# $+|x(t)|^{2} - ||x_{t}||^{2} + \int_{t-h}^{t} \frac{d}{dt} ||x^{2}(\cdot)||^{[s,t]} ds$ $\leq (-2a(t) + |b(t)| + \alpha(t) + 1)|x(t)|^{2} + 2|x(t)||f(t, x_{t})| - ||x_{t}||^{2}$ $- \frac{1}{h} (1 - |b(t)|h^{2}) \int_{t-h}^{t} \alpha(s)|x(s)|^{2} ds + \int_{t-h}^{t} \frac{d}{dt} ||x^{2}(\cdot)||^{[s,t]} ds.$

For each fixed s, if  $||x^2(\cdot)||^{[s,t]} = |x(\theta)|^2$  with  $s \le \theta < t$  and  $|x(\tau)| < |x(\theta)|$  for all  $\theta < \tau \le t$ , then  $\frac{d}{dt}||x^2(\cdot)||^{[s,t]} = 0$  (see Hale [4], p. 127). Now suppose that  $||x^2(\cdot)||^{[s,t]} = |x(t)|^2$ . Then  $|x(t)| \ge |x(\tau)|$  for  $\tau \in [s,t]$  and

$$\frac{d}{dt} \|x^{2}(\cdot)\|^{[s,t]} \leq -2a(t)|x(t)|^{2} + |b(t)||x(t)|^{2} 
+ |b(t)|h \int_{t-h}^{t} |C(s,x(s))|^{2} ds + 2|x(t)||f(t,x_{t})|| 
\leq |b(t)|h \int_{t-h}^{t} |C(s,x(s))|^{2} ds + 2|x(t)||f(t,x_{t})|.$$

Thus,

$$V'_{(E)}(t,x_{t}) \leq (-2a(t) + |b(t)| + \alpha(t) + 1)|x(t)|^{2} + 2(1+h)|x(t)| |f(t,x_{t})|$$

$$- ||x_{t}||^{2} - \frac{1}{h}[1 - |b(t)|(1+h)h^{2}] \int_{t-h}^{t} \alpha(s)|x(s)|^{2} ds$$

$$\leq (-2a(t) + |b(t)| + \alpha(t) + 1)|x(t)|^{2} + 2(1+h)p||x_{t}||^{2} - ||x_{t}||^{2}$$

$$\leq -\beta(t)|x(t)|^{2}.$$
(11)

It follows from (10) that

$$(1+h)|\phi(0)|^2 \le V(t,\phi) \le 2(1+L)[|\phi(0)| + |\phi|_2]^2 + h\|\phi\|^2. \tag{12}$$

Let  $W(r) = r^2$ ,  $W_2(r) = 2(1 + L)r^2$ ,  $W_1(r) = (1 + h)r^2$ , and  $W_3(r) = hr^2$ . Then all conditions of Corollary 7 are satisfied and the zero solution of (E) is UAS.

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## ANISOTROPIC EQUATIONS IN $L^1$

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**1.** Introduction. Existence results for weak solutions of equations involving the p-Laplacian with right hand side measure have been recently obtained in [1]. G. Stampacchia [5] first considered the problem in the linear case; he obtained an existence result through a duality method: if  $T \in W^{-1,r}(\Omega)$ , r > N, then the solution of the Dirichlet problem

$$w \in H_0^1(\Omega) : -\operatorname{div}(A(x)Dw) = T \tag{1.1}$$

lies in  $C^0(\Omega)$  and the mapping  $T \to w$  is linear and continuous from  $W^{-1,r}(\Omega)$  to  $C^0(\Omega)$ . Therefore the adjoint operator maps  $M_b(\Omega)$  into  $W_0^{1,q}(\Omega)$ , for any  $q < \frac{N}{N-1}$ .

In this paper we consider anisotropic equations with right hand side measures. We confine ourselves to the model case

$$\begin{cases}
-\operatorname{div}(j(Du)) = f & \text{in } \Omega \\
u = 0 & \text{on } \partial\Omega,
\end{cases}$$
(1.2)

where  $j(\xi)$  is the vector field whose components are  $|\xi_i|^{p_i-2}\xi_i$  ( $i=1,\ldots,N;\ p_i>1$ ). We shall prove the existence of solutions of (1.2). More precisely we obtain the existence of a solution in the Sobolev space

$$W_0^{1,q_i}(\Omega)\big\{v\in W_0^{1,1}(\Omega): \frac{\partial v}{\partial x_i}\in L^{q_i}(\Omega),\quad \forall i=1,\ldots,N\big\},$$

where  $q_i$  (for i = 1, ..., N) is any real number greater than 1 and such that

$$1 < q_i < \frac{N(\bar{p}-1)}{\bar{p}(N-1)} p_i,$$

and

$$\frac{1}{\bar{p}} = \frac{1}{N} \sum_{i=1}^{N} \frac{1}{p_i}, \quad p_i > 1. \tag{1.3}$$

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Our existence result can be seen as a "nonlinear version" of the Stampacchia method, since it is known that the Dirichlet problem

$$\begin{cases} -\operatorname{div}(j(Dw)) = -\operatorname{div}(j(F)) \\ u \in W_0^{1,p_i}(\Omega) \end{cases}$$
 (1.4)

has an  $L^{\infty}$  solution if each component  $F_i$  of the vector field F belongs to  $L^{r_i}(\Omega)$ ,  $r_i > N \frac{p_i}{\bar{p}}$  (see [6], [3], [4]).

**2. Existence theorems.** Let  $\Omega$  be an open bounded set of  $\mathbb{R}^N$   $(N \ge 1)$ ,  $p_i > 1$  (i = 1, ..., N) and  $\bar{p} < N$ . The aim of this paper is to obtain weak solutions of the (model) problem

 $\begin{cases} -\sum_{i=1}^{N} \frac{\partial}{\partial x_i} (|\frac{\partial u}{\partial x_i}|^{p_i - 2} \frac{\partial u}{\partial x_i}) = \mu & \text{in } \Omega \\ u = 0 & \text{on } \partial \Omega \end{cases}$  (2.1)

in the sense of distributions, when  $\mu$  is a bounded Randon measure on  $\Omega$  (i.e.,  $\mu \in M_b(\Omega)$ ).

We shall prove the following existence theorem.

**Theorem 1.** Let  $\mu \in M_b(\Omega)$ . Then there exists  $u \in W_0^{1,1}(\Omega)$  (such that  $\frac{\partial u}{\partial x_i} \in L^{q_i}(\Omega)$ ,  $q_i < \frac{N(\bar{p}-1)}{\bar{p}(N-1)}p_i$ ) solution of (2.1)

**Proof.** As in [1] we consider a sequence of smooth functions  $f_k$  converging to  $\mu$  in  $M_b(\Omega)$  weak, such that  $||f_k||_{L^1(\Omega)} \le ||\mu||_{M_b(\Omega)}$ , and the solutions  $u_k$  of the equations

$$u_k \in W_0^{1,p_i}(\Omega) : -\sum_{i=1}^N \frac{\partial}{\partial x_i} \left( \left| \frac{\partial u_k}{\partial x_i} \right|^{p_i - 2} \frac{\partial u}{\partial x_i} \right) = f_k.$$
 (2.2)

We define

$$\varphi(s) = \begin{cases} 0 & 0 \le s \le n \\ s - n & n < s < n + 1 \\ 1 & s \ge n + 1 \\ -\varphi(-s) & s < 0 \end{cases}$$

and we use  $\varphi(u_k)$  as test function in (2.2). Let  $B_{n,k} = \{x \in \Omega : n \le |u_k(x)| < n+1\}$ . Then

$$\int_{B_{n,k}} \sum_{i=1}^{N} \left| \frac{\partial u_k}{\partial x_i} \right|^{p_i} \le c_1,$$

which implies

$$\int_{R_{-}} \left| \frac{\partial u_k}{\partial x_i} \right|^{p_i} \le c_1, \quad i = 1, \dots, N.$$
 (2.3)

Let  $q_i = \theta p_i$ ,  $\lambda = \frac{(1-\theta)\bar{q}^*}{\theta}$  and  $\theta \in (0, \frac{N(\bar{p}-1)}{\bar{p}(N-1)})$ . Then, using (2.3),  $\lambda > 1$  and Holder inequality,

$$\int_{\Omega} \left| \frac{\partial u_{k}}{\partial x_{i}} \right|^{q_{i}} \leq \left[ \int_{\Omega} \left| \frac{\partial u_{k}}{\partial x_{i}} \right|^{p_{i}} (1 + |u_{k}|)^{-\lambda} \right]^{\frac{q_{i}}{p_{i}}} \left[ \int_{\Omega} (1 + |u_{k}|)^{\frac{\lambda q_{i}}{p_{i} - q_{i}}} \right]^{1 - \frac{q_{i}}{p_{i}}} \right] dt \\
\leq \left[ \sum_{n=0}^{\infty} \frac{1}{(1 + n)^{\lambda}} \int_{B_{n,k}} \left| \frac{\partial u_{k}}{\partial x_{i}} \right|^{p_{i}} \left[ \int_{\Omega} (1 + |u_{k}|)^{\frac{\lambda q_{i}}{p_{i} - q_{i}}} \right]^{1 - \frac{q_{i}}{p_{i}}}. \tag{2.4}$$

The use of the nonisotropic Sobolev inequality (see [7]) leads to

$$||u_{k}||_{L^{\bar{q}^{*}}} \leq c_{3} \prod_{i=1}^{N} \left\| \frac{\partial u_{k}}{\partial x_{i}} \right\|_{L^{q_{i}}}^{1/N} \leq c_{4} \prod_{i=1}^{N} \left[ \int_{\Omega} (1 + |u_{k}|)^{\frac{\lambda q_{i}}{p_{i} - q_{i}}} \right]^{(\frac{1}{q_{i}} - \frac{1}{p_{i}})^{\frac{1}{N}}}$$

$$= c_{4} \prod_{i=1}^{N} \left[ \int_{\Omega} (1 + |u_{k}|)^{\frac{\lambda \theta}{1 - \theta}} \right]^{(\frac{1}{\theta} - 1)^{\frac{1}{Np_{i}}}} = c_{4} \left[ \int_{\Omega} (1 + |u_{k}|)^{\bar{q}^{*}} \right]^{\frac{1 - \theta}{\theta} \frac{1}{\bar{p}}}.$$

$$(2.5)$$

By estimate (2.4),  $\frac{\partial u_k}{\partial x_i}$  is bounded in  $L^{q_i}(\Omega)$ . Then we can assume (for some u and for some subsequence still denoted  $u_k$ ) that

$$\frac{\partial u_k}{\partial x_i} \to \frac{\partial u}{\partial x_i} \quad \text{weakly in } L^{q_i}(\Omega), \tag{2.6}$$

$$u_k \to u$$
 strongly in  $L^{\bar{q}}(\Omega)$ . (2.7)

This is not sufficient to pass to the limit, but the monotonicity properties of the differential operator allow us to use the techniques of [1] in order to deduce that we have also

$$\frac{\partial u_k}{\partial x_i} \to \frac{\partial u}{\partial x_i} \quad \text{in } L^{r_i}(\Omega), \quad r_i < q_i.$$
 (2.8)

Indeed, by (2.2), we have, for any  $\eta > 0$ ,

$$\sum_{i=1}^{N} \int_{\Omega} \left[ \left| \frac{\partial u_k}{\partial x_i} \right|^{p_i - 2} \frac{\partial u_k}{\partial x_i} - \left| \frac{\partial u_h}{\partial x_i} \right|^{p_i - 2} \frac{\partial u_h}{\partial x_i} \right] \frac{\partial}{\partial x_i} T_{\eta}(u_k - u_h) = \int (f_k - f_h) T_{\eta}(u_k - u_h),$$

where  $T_n$  is the truncation at levels  $\pm \eta$  ( $\eta > 0$ ). Then, for  $i = 1, \dots, N$ ,

$$\int_{|u_k-u_h|\leq \eta} \left[ \left| \frac{\partial u_k}{\partial x_i} \right|^{p_i-2} \frac{\partial u_k}{\partial x_i} - \left| \frac{\partial u_h}{\partial x_i} \right|^{p_i-2} \frac{\partial u_h}{\partial x_i} \right] \frac{\partial}{\partial x_i} (u_k - u_h) \leq 2\eta \|\mu\|.$$

If, for a fixed i,  $p_i$  is greater than or equal to 2, we deduce that

$$\int_{|u_k-u_h|<\eta} \left| \frac{\partial (u_k-u_h)}{\partial x_i} \right|^{p_i} \le c_5 \eta$$

and then

$$\int_{\Omega} \left| \frac{\partial (u_k - u_h)}{\partial x_i} \right|^{p_i} \le c_6 \left[ \int_{|u_k - u_h| \le \eta} \left| \frac{\partial (u_k - u_h)}{\partial x_i} \right|^{p_i} \right]^{\frac{r_i}{p_i}}$$

$$+ c_7 \text{mis} \left\{ x \in \Omega : |u_k(x) - u_h(x)| > \eta \right\}^{1 - \frac{r_i}{q_i}}$$

$$\le c_8 \eta^{\frac{r_i}{p_i}} + c_7 \text{mis} \left\{ x \in \Omega : |u_k(x) - u_h(x)| > \eta \right\}^{1 - \frac{r_i}{q_i}}.$$

Recall that  $u_k$  is a Cauchy sequence is measure. The arbitrariness of  $\eta > 0$  leads to the conclusion that  $\frac{\partial u_k}{\partial x_i}$  is a Cauchy sequence in  $L^{r_i}(\Omega)$ . A slight modification is needed to obtain (2.8) for  $p_i < 2$ . Then by (2.8), we have

$$\left|\frac{\partial u_k}{\partial x_i}\right|^{p_i-2}\frac{\partial u_k}{\partial x_i} \to \left|\frac{\partial u}{\partial x_i}\right|^{p_i-2}\frac{\partial u}{\partial x_i} \quad \text{in } L^1(\Omega)$$

and therefore u is a solution of equation (2.1).

Theorem 1 allows us to find a solution when, for instance, the data are  $L^1$ -functions. Anyway, to obtain  $q_i = \frac{N(\bar{p}-1)}{\bar{p}(N-1)} p_i$ , we have to make stronger assumptions on the right hand side. This is stated in the following.

**Theorem 2.** Let f be a measurable function such that

$$\int_{\Omega} |f| \log(1+|f|) < \infty.$$

Then there exists a solution u of the Dirichlet problem

$$u \in W_0^{1,q_i}(\Omega) : -\sum_{i=1}^N \frac{\partial}{\partial x_i} \left( \left| \frac{\partial u}{\partial x_i} \right|^{p_i-2} \frac{\partial u}{\partial x_i} \right) = f,$$

where  $q_i = \frac{N(\bar{p}-1)}{\bar{p}(N-1)}p_i$ .

**Proof.** We modify the previous proof with the help of techniques used in [2]. Using  $\log(1+|u_k|)\operatorname{sgn}(u_k)$  as test function in (2.2), the inequality (2.3) becomes

$$\int_{\Omega} \left| \frac{\partial u_k}{\partial x_i} \right|^{p_i} \cdot \frac{1}{1 + |u|} \le \int_{\Omega} f \log(1 + |u_k|)$$

$$\le \int_{\Omega} |f| \log(1 + |f|) + \int (1 + |u_k|) \le c_1 + \int_{\Omega} (1 + |u_k|).$$

This implies the change in inequality (2.5),

$$||u_k||_{\tilde{q}^*} \leq c_2 + c_2 \left[ \int_{\Omega} (1+|u_k|)^{\frac{\theta}{1-\theta}} \right]^{\frac{1-\theta}{\tilde{q}}},$$

and so the sequence  $u_k$  is bounded in  $W_0^{1,q_i}(\Omega)$ . The a priori estimate and the strong  $L^1$  convergence of  $\frac{\partial u_k}{\partial x_i}$  allow us to pass to the limit.

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# A FINITE RANGE OPERATOR WITH A QUASI-PERIODIC POTENTIAL

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**Abstract.** We consider the operator  $H = \varepsilon h + \cos(\alpha \cdot j + \vartheta)$  for  $j \in \mathbb{Z}^{\nu}$ , where h is self-adjoint, translation invariant and finite range. The vector  $\alpha = (\alpha_1, \alpha_2, \dots, \alpha_{\nu})$  is assumed to have the diophantine property  $|j \cdot \alpha \mod 2\pi| \ge C/|j|^2$ , where  $j \ne 0$  and C is some constant. For  $\varepsilon$  sufficiently small we prove that H has pure point spectrum for almost every  $\vartheta$ . Moreover, every polynomially bounded eigenfunction of H decays exponentially fast. Finally, we will show how this operator comes up in the study of electrons in a transverse magnetic field subject to a two dimensional periodic potential.

### 1. Introduction. We consider the equation

$$H\psi \equiv (\varepsilon h + V)\psi = E\psi$$

on the lattice  $\mathbb{Z}^{\nu}$ . We assume that h is self-adjoint, translation invariant and finite range; i.e., the matrix elements of h satisfy,  $h_{ij}=0$  when  $|i-j|\geq R$ , where  $i,j\in\mathbb{Z}^{\nu}$  and R is some positive constant. The potential V is given by matrix elements

$$V_{ij}(\vartheta) = \delta_{ij} v_j(\vartheta) = \delta_{ij} \cos(\alpha \cdot j + \vartheta)$$

and the vector  $\alpha = (\alpha_1, \alpha_2, \dots, \alpha_{\nu})$  is assumed to have the Diophantine property

$$|j \cdot \alpha \bmod 2\pi| \ge C/|j|^2,$$

where  $i \neq 0$  and C is some constant.

We state our main result:

**Theorem.** Consider the operator  $H = \varepsilon h + \cos(\alpha \cdot j + \vartheta)$   $(j \in \mathbb{Z}^{\nu})$  subject to the conditions listed above. For  $\varepsilon$  sufficiently small H has pure point spectrum for almost every  $\vartheta$ . Moreover, every polynomially bounded eigenfunction of H decays exponentially fast.

Our motivation for considering operators of this type comes from the study of electrons in a transverse magnetic field subject to a two dimensional periodic potential, [26], [27]. In Section 4, we will examine the Schrödinger equation in a Magnetic field. We will make a series expansion of the solution to the differential equation and show that the recursion formula for the coefficients in the expansion leads us to the finite difference equation treated in this paper. In recent years many problems of this type have been studied. We now mention a few of the latest results.

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