

Hyperbolic equations and systems with discontinuous coefficients or source terms

T. Gallouët¹

¹University of Marseille

Equadiff 11, 2005

Notations:

$$(\cdot)_t = \frac{\partial(\cdot)}{\partial t},$$

$$(\cdot)_x = \frac{\partial(\cdot)}{\partial x},$$

$$t \in \mathbb{R}_+.$$

$x \in \mathbb{R}$, but extensions to $x \in \mathbb{R}^d$, $d = 2$ or 3 are possible.

Hyperbolic equation with a discontinuous coefficient

Two phase flow in an heterogeneous porous medium:

unknown: $u : \mathbb{R} \times \mathbb{R}_+ \rightarrow \mathbb{R}$:

$$u_t(x, t) + (kg(u))_x = 0,$$

$$k(x) = k_l, \text{ for } x < 0,$$

$$k(x) = k_r, \text{ for } x > 0,$$

$$k_l, k_r > 0, k_l \neq k_r,$$

$g : [0, 1] \rightarrow \mathbb{R}$, Lipschitz continuous and such that

$$g(0) = g(1) = 0. \text{ Example: } g(u) = u(1 - u)$$

This hyperbolic equation with a discontinuous coefficient can be viewed as a conservative 2×2 hyperbolic system, adding k has an unknown and the equation $k_t = 0$.

Hyperbolic equation with a discontinuous coefficient

$$\begin{aligned}u_t(x, t) + (kg(u))_x &= 0, \\k_t &= 0.\end{aligned}$$

$$W = \begin{bmatrix} u \\ k \end{bmatrix} \text{ and } F(W) = \begin{bmatrix} kg(u) \\ 0 \end{bmatrix},$$

$$W_t + (F(W))_x = 0,$$

or equivalently (for regular solutions), with $A(W) = DF(W)$:

$$W_t + A(W)W_x = 0.$$

Hyperbolic system with a source term

Saint Venant Equations with topography (nonflat bottom)

unknowns: $h, u : \mathbb{R} \times \mathbb{R}_+ \rightarrow \mathbb{R}$ (with $h > 0$):

$$\begin{aligned}h_t + (hu)_x &= 0, \\(hu)_t + (hu^2 + \tfrac{1}{2}gh^2)_x &= -ghz_x,\end{aligned}$$

g is a given constant and z is a given function of x .

This 2×2 conservative hyperbolic system with a source term can be viewed as a nonconservative 3×3 hyperbolic system, adding z has an unknown and the equation $z_t = 0$.

Hyperbolic system with a source term

$$\begin{aligned}h_t + (hu)_x &= 0, \\(hu)_t + (hu^2 + \tfrac{1}{2}gh^2)_x + ghz_x &= 0, \\z_t &= 0.\end{aligned}$$

$$W = \begin{bmatrix} u \\ hu \\ z \end{bmatrix}, F(W) = \begin{bmatrix} hu \\ \frac{(hu)^2}{h} + \frac{1}{2}gh^2 \\ 0 \end{bmatrix} \text{ and}$$

$$B(W) = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & gh \\ 0 & 0 & 0 \end{bmatrix},$$

$$W_t + (F(W))_x + B(W)W_x = 0,$$

or equivalently (for regular solutions):

$$W_t + A(W)W_x = 0, \quad (\text{with } A(W) = DF(W) + B(W)).$$

Common feature : Resonance

Unknown: $W : \mathbb{R} \times \mathbb{R}_+ \rightarrow \mathbb{R}^p$, $p > 1$

Equation: $W_t + A(W)W_x = 0$

For all admissible $W \in \mathbb{R}^p$, the $p \times p$ matrix $A(W)$ has only real eigenvalues, but $A(W)$ is, for some W , not diagonalizable.

$A(W)$ strictly hyperbolic: p real distinct eigenvalues and then diagonalizable (in \mathbb{R}).

$A(W)$ genuine-hyperbolic: real eigenvalues and diagonalizable.

$A(W)$ resonant-hyperbolic: real eigenvalues and not diagonalizable.

First example:

$p = 2$, $W = \begin{bmatrix} u \\ k \end{bmatrix}$, $A(W) = \begin{bmatrix} kg'(u) & g(u) \\ 0 & 0 \end{bmatrix}$ which is not diagonalizable if $g'(u) = 0$ and $g(u) \neq 0$ (if $g(u) = u(1 - u)$, this is the case for $u = \frac{1}{2}$).

Common feature : Resonance

Second example:

$$p = 3, W = \begin{bmatrix} u \\ hu \\ z \end{bmatrix}, A(W) = \begin{bmatrix} 0 & 1 & 0 \\ -u^2 + gh & 2u & gh \\ 0 & 0 & 0 \end{bmatrix}.$$

Eigenvalues of $A(W)$ are $u \pm c$ and 0 , with $c = \sqrt{gh}$.

$A(W)$ is not diagonalizable if $u - c = 0$ or $u + c = 0$ (and $h > 0$).

Linear resonant problem, ill posed

The Cauchy problem for a linear resonant problem is ill-posed in L^∞ (or in $L^1, L^2 \dots$, but well posed in C^∞).

Riemann problem for a typical example:

$$\begin{bmatrix} u \\ v \end{bmatrix}_t + \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} u \\ v \end{bmatrix}_x = 0,$$

$$\begin{bmatrix} u(x, 0) \\ v(x, 0) \end{bmatrix} = \begin{bmatrix} u_l \\ v_l \end{bmatrix}, \text{ if } x < 0, \text{ and } \begin{bmatrix} u_r \\ v_r \end{bmatrix}, \text{ if } x > 0,$$

The solution is , for all $t > 0$, $v(\cdot, t) = v(\cdot, 0)$ and

$$u(\cdot, t) = u_l 1_{\mathbb{R}_-} + u_r 1_{\mathbb{R}_+} + t(v_l - v_r)\delta_0.$$

Nonlinear resonant problem, sometimes ill posed

Academic simple example:

$$\begin{aligned}u_t + (au)_x &= 0, \\a_t &= 0,\end{aligned}$$

$$\begin{bmatrix} u(x, 0) \\ a(x, 0) \end{bmatrix} = \begin{bmatrix} u_l \\ a_l \end{bmatrix}, \text{ if } x < 0, \text{ and } \begin{bmatrix} u_r \\ a_r \end{bmatrix}, \text{ if } x > 0,$$

has no weak solution in L^∞ if $a_l > 0$, $a_r < 0$ and $a_l u_l \neq a_r u_r$
(and has infinitely many solutions if $a_l < 0$ and $a_r > 0$).

Nonlinear resonant problem, sometimes ill posed

$$\begin{aligned}u_t + (au)_x &= 0, \\ a_t &= 0,\end{aligned}$$

is equivalent (for regular solution) to $W_t + A(W)W_x = 0$, with $W = \begin{bmatrix} u \\ a \end{bmatrix}$ and $A(W) = \begin{bmatrix} a & u \\ 0 & 0 \end{bmatrix}$. Resonance occurs when $a = 0$ and $u \neq 0$.

The Riemann problem for the nonlinear system is ill posed in L^∞ provided that 0 is between a_r and a_l (except if $u_r = u_l = 0$).

This nonlinear case is “worse” than the linear case.

Nonlinear resonant problem, sometimes well posed

First example

$$\begin{aligned}u_t + (kg(u))_x &= 0, \\k_t &= 0,\end{aligned}$$

$g : [0, 1] \rightarrow \mathbb{R}$, Lipschitz continuous and such that $g(0) = g(1) = 0$. The Cauchy problem is well posed in the following sense:

$k(x) = k_l$ if $x < 0$, $k(x) = k_r$ if $x > 0$, $u(\cdot, 0) = u_0 \in L^\infty$, $0 \leq u_0 \leq 1$. Then, the Cauchy problem has a unique entropy weak solution.

Karlsen-Risebro-Towers

Seguin-Vovelle, Bachmann-Vovelle

Bachmann: $k(x)g(u) \rightsquigarrow g(x, u)$, $g(x, 0) = g(x, 1) = 0$.

Nonlinear resonant problem, sometimes well posed

Second example

Saint Venant Equations with topography. The Riemann problem:

$$\begin{aligned}h_t + (hu)_x &= 0, \\(hu)_t + (hu^2 + \frac{1}{2}gh^2)_x + ghz_x &= 0, \\z_t &= 0\end{aligned}$$

$$\begin{bmatrix} h \\ hu \\ z \end{bmatrix} (x, 0) = \begin{bmatrix} h_l \\ (hu)_l \\ z_l \end{bmatrix}, \text{ if } x < 0, \text{ and } \begin{bmatrix} h_r \\ (hu)_r \\ z_r \end{bmatrix}, \text{ if } x > 0,$$

has one (sometimes three...) solution (composed of constant states and waves), satisfying a classical entropy condition, assuming continuity of (the Riemann invariants) hu and ψ at the contact discontinuity (at $x = 0$) with $\psi = \frac{1}{2}u^2 + g(h + z)$.

Chinnayya-LeRoux-Seguin, Goatin-LeFloch

Nonlinear resonant problem, sometimes well posed

Third example

Isentropic Euler Equations with an EOS taking into account a simple model of “phase transition”, that is:

$$\begin{aligned}p &= a_1 \rho, \text{ if } 0 < \rho < \rho_1, \\p &= a_1 \rho_1, \text{ if } \rho_1 \leq \rho \leq \rho_2, \\p &= a_2 \rho, \text{ if } \rho_2 < \rho,\end{aligned}$$

with ρ_1, ρ_2, a_1, a_2 given constants, $0 < \rho_1 < \rho_2$, $0 < a_1 < a_2$.

$$\begin{aligned}\rho_t + (\rho u)_x &= 0, \\(\rho u)_t + (\rho u^2 + p)_x &= 0.\end{aligned}$$

For $\rho_1 \leq \rho \leq \rho_2$ and any u , the system is resonant (with u as eigenvalue, and the 2 genuinely nonlinear fields lead to a linear degenerate field). But the Riemann problem is well posed.

Recent result of [Seguin](#).

Discretization by Finite Volume Schemes

$$W_t + (F(W))_x + B(W)W_x = 0,$$

$$W(\cdot, 0) = W_0.$$

Time step: k , $t_n = nk$

Space step: h , $x_{i+\frac{1}{2}} = (i + \frac{1}{2})h$

Approximate solution for $x \in (x_{i-\frac{1}{2}}, x_{i+\frac{1}{2}})$ and $t = t_n$: W_i^n

$$W_i^0 = \frac{1}{h} \int_{x_{i-\frac{1}{2}}}^{x_{i+\frac{1}{2}}} W_0(x) dx, \quad i \in \mathbb{Z}.$$

$$W_t + (F(W))_x + B(W)W_x = 0,$$

$$W(x, 0) = \begin{cases} W_l & \text{if } x < 0, \\ W_r & \text{if } x > 0. \end{cases}$$

Let W be the (or a) self similar solution, that is:

$$W(x, t) = R\left(\frac{x}{t}, W_l, W_r\right)$$

and set

$$W^{*,\pm}(W_l, W_r) = R(0^\pm, W_l, W_r).$$

Godunov scheme for nonconservative system

For $i \in \mathbb{Z}$, $n \geq 0$:

$$\frac{W_i^{n+1} - W_i^n}{k} + F_{i+\frac{1}{2}}^{n,-} - F_{i-\frac{1}{2}}^{n,+} + B(W_i^n)(W_{i+\frac{1}{2}}^{n,-} - W_{i-\frac{1}{2}}^{n,+}) = 0,$$

with $F_{i+\frac{1}{2}}^{n,\pm} = F(W_{i+\frac{1}{2}}^{n,\pm})$,

$W_{i+\frac{1}{2}}^{n,\pm} = W^{*,\pm}(W_i^n, W_{i+1}^n)$.

CFL condition: $k \leq Ch$.

Godunov scheme, particular cases

$$\frac{W_i^{n+1} - W_i^n}{k} + F_{i+\frac{1}{2}}^{n,-} - F_{i-\frac{1}{2}}^{n,+} + B(W_i^n)(W_{i+\frac{1}{2}}^{n,-} - W_{i-\frac{1}{2}}^{n,+}) = 0,$$

- 1 conservative case: $F_{i+\frac{1}{2}}^{n,+} = F_{i+\frac{1}{2}}^{n,-}$, even if $W_{i+\frac{1}{2}}^{n,+} \neq W_{i+\frac{1}{2}}^{n,-}$
(thanks to Rankine-Hugoniot condition on the Riemann problem).
- 2 In the case of Saint Venant equations with topography, one has $z_{i+\frac{1}{2}}^{n,-} = z_{i-\frac{1}{2}}^{n,+} = z_i^n$ and then $B(W_i^n)(W_{i+\frac{1}{2}}^{n,-} - W_{i-\frac{1}{2}}^{n,+}) = 0$.
The non conservativity of the equation (that is the source term) appears only in the fact that, generally, $F_{i+\frac{1}{2}}^{n,+} \neq F_{i+\frac{1}{2}}^{n,-}$.

Linearized system

- Initial system: $W_t + A(W)W_x = 0$,
 $A(W) = DF(W) + B(W)$.
- Linearized system: $W_t + A(\overline{W})W_x = 0$, with some fixed $\overline{W} \in \mathbb{R}^p$.
- Initial system with a change of unknown: $Y = \phi(W)$. ϕ is not necessarily invertible, but one assumes that there exists C, G, D , such that $A(W) = C(Y)$, $F(W) = G(Y)$, $B(W)W_x = D(W)Y_x$. Then $W_t + A(W)W_x$ leads to

$$Y_t + C(W)Y_x = 0.$$

- Linearized system with a change of unknown: $Y = \phi(W)$,

$$Y_t + C(\overline{W})Y_x = 0,$$

with some fixed $\overline{W} \in \mathbb{R}^p$.

linearized Riemann problem

$$Y_t + C(\overline{W}) Y_x = 0,$$

$$Y(x, 0) = \begin{cases} Y_l = \phi(W_l) & \text{if } x < 0, \\ Y_r = \phi(W_r) & \text{if } x > 0. \end{cases}$$

For instance, $\overline{W} = \frac{W_r + W_l}{2}$ or another mean value between W_l and W_r .

Let Y be the self similar solution of this problem (when it exists...), that is:

$$Y(x, t) = LR\left(\frac{x}{t}, Y_l, Y_r\right)$$

and set $Y^{*,\pm}(W_l, W_r) = LR(0^\pm, Y_l, Y_r)$.

For $i \in \mathbb{Z}$, $n \geq 0$:

$$\frac{W_i^{n+1} - W_i^n}{k} + F_{i+\frac{1}{2}}^{n,-} - F_{i-\frac{1}{2}}^{n,+} + D(W_i^n)(Y_{i+\frac{1}{2}}^{n,-} - Y_{i-\frac{1}{2}}^{n,+}) = 0,$$

with:

$$F_{i+\frac{1}{2}}^{n,\pm} = G(Y_{i+\frac{1}{2}}^{n,\pm}) (F(W) = G(\phi(W)), B(W)W_x = D(W)Y_x),$$
$$Y_{i+\frac{1}{2}}^{n,\pm} = Y^{*,\pm}(W_i^n, W_{i+1}^n)$$

CFL condition: $k \leq Ch$.

- 1 Modification for conservativity (null eigenvalue)
- 2 Discontinuity of the numerical flux when there is 0 as eigenvalue. $F(W^*(W_l, W_r))$ is a discontinuous fonction of W_l and W_r .
- 3 Choice of Y .
 $Y = (kg(u), k)^t$ for porous media,
 $Y = (2c, u, z)^t$ or $Y = (q, \psi, z)^t$ for Saint Venant with topography.

Remark on resonance, porous media

$$\begin{aligned}u_t(x, t) + (kg(u))_x &= 0, \\k_t &= 0.\end{aligned}$$

With the choice $Y = (kg(u), k)^t$ the linearized system is $Y_t + C(\overline{W})Y_x = 0$, $\overline{W} = (\overline{u}, \overline{k})^t$,

$$C(\overline{W}) = \begin{bmatrix} \overline{k}g'(\overline{u}) & 0 \\ 0 & 0 \end{bmatrix},$$

which is never resonant. . .

Eigenvalues: $\lambda_1 = \overline{k}g'(\overline{u})$, $\lambda_2 = 0$

Eigenvectors: $e_1 = (1, 0)^t$, $e_2 = (0, 1)^t$

Remark on resonance, Saint Venant with topography

$$\begin{aligned}h_t + (hu)_x &= 0, \\(hu)_t + (hu^2 + \tfrac{1}{2}gh^2)_x + ghz_x &= 0, \\z_t &= 0\end{aligned}$$

With the choice $Y = (q, \psi, z)^t$, the linearized system is

$$Y_t + C(\overline{W})Y_x = 0, \quad \overline{W} = (\overline{h}, \overline{h}\overline{u}, \overline{z})^t, \quad \overline{c} = \sqrt{g\overline{h}},$$

$$C(\overline{W}) = \begin{bmatrix} \overline{u} & \overline{h} & 0 \\ g & \overline{u} & 0 \\ 0 & 0 & 0 \end{bmatrix},$$

which is never resonant (for $\overline{h} > 0$)...

Eigenvalues: $\lambda_1 = 0$, $\lambda_2 = \overline{u} + \overline{c}$, $\lambda_3 = \overline{u} - \overline{c}$

Eigenvectors: $e_1 = (0, 0, 1)^t$, $e_2 = (\overline{h}, \overline{c}, 0)^t$, $e_3 = (-\overline{h}, \overline{c}, 0)^t$

Remark on resonance, Saint Venant with topography

With the choice $Y = (2c, u, z)^t$, the linearized system is

$$Y_t + C(\overline{W})Y_x = 0, \quad \overline{W} = (\overline{h}, \overline{h}\overline{u}, \overline{z})^t, \quad \overline{c} = \sqrt{g\overline{h}},$$

$$C(\overline{W}) = \begin{bmatrix} \overline{u} & \overline{c} & 0 \\ \overline{c} & \overline{u} & g \\ 0 & 0 & 0 \end{bmatrix},$$

which is resonant if $\overline{u} + \overline{c} = 0$ or $\overline{u} - \overline{c} = 0$.

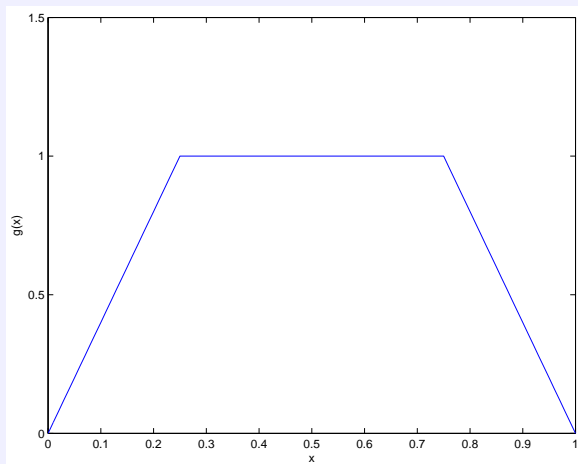
Eigenvalues: $\lambda_1 = 0$, $\lambda_2 = \overline{u} + \overline{c}$, $\lambda_3 = \overline{u} - \overline{c}$

Eigenvectors:

- If $\overline{u} \pm \overline{c} \neq 0$,
 $e_1 = (\overline{c}g, -\overline{u}g, \overline{u}^2 - \overline{c}^2)^t$, $e_2 = (1, 1, 0)^t$, $e_3 = (1, -1, 0)^t$
- If $\overline{u} = \overline{c}$, $e_1 = (1, -1, 0)^t$, $e_2 = (1, 1, 0)^t$
- If $\overline{u} = -\overline{c}$, $e_1 = (1, 1, 0)^t$, $e_3 = (1, -1, 0)^t$

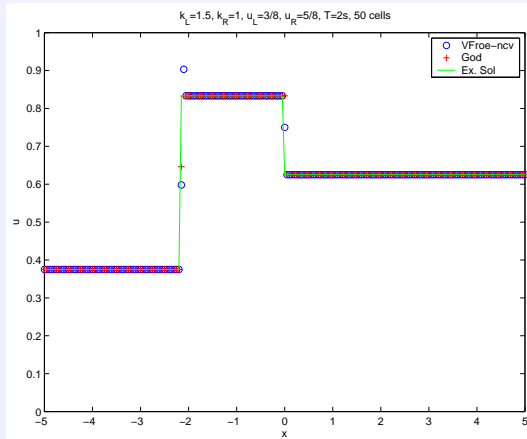
- 1 Perfect with Godunov. In particular, for Saint Venant with topography, one has preservation of **all** steady state solutions : $q = hu$ and $\psi = \frac{u^2}{2} + g(h + z)$ constant (and not only those with $u = 0$ which are called "lake at rest")
- 2 Perfect also with VFRoe-ncv, sometimes with one "incorrect point". For Saint Venant with topography: Preservation of steady state solution with $u = 0$ (for any choice of the variable of linearization Y) and preservation of **all** steady state solutions for $Y = (q, \psi, z)^t$

numerical results, two phase flow in porous media



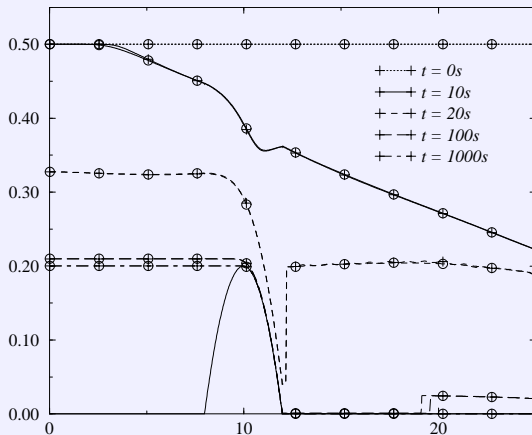
Resonance occurs for all (k, u) with $u \in (\frac{1}{4}, \frac{3}{4})$.

numerical results, two phase flow in porous media



Resonance occurs for all (k, u) with $u \in (\frac{1}{4}, \frac{3}{4})$.

numerical results, Saint Venant with topography



A case with $h = 0$, steady state, sonic point. . . VFRoe-ncv with $Y = (2c, u, z)^t$ (best choice for the problem of vanishing h).

- ① Porous medium: Towers, Seguin-Vovelle, Bachmann
- ② Saint Venant: Kurganov-Levy, Simeoni-Perthame, Chinnayya-LeRoux-Seguin, Goutal-Karni Hérard-Gallouët-Seguin. . .