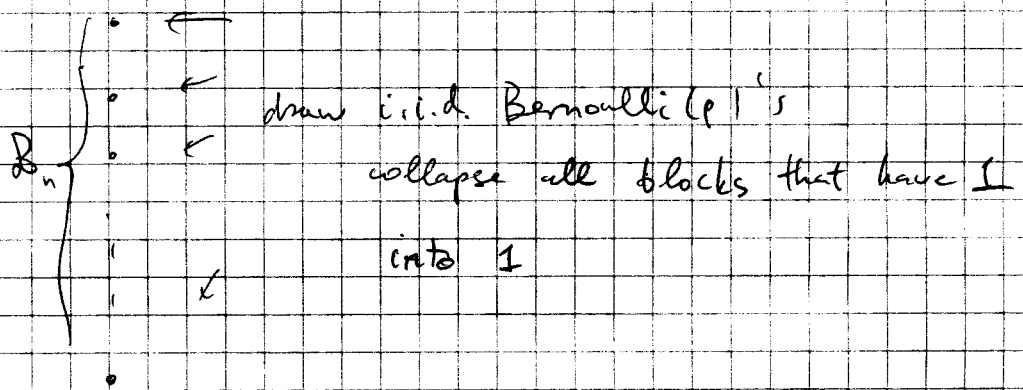


Speed of CDI for \boxminus -coalescents

CIRM 2003

Assume we are given n blocks B_n , and perform coalescence operator

$p \in (0, 1)$ coal_p
 $1 \in \mathcal{K} \cap \mathcal{Q}$
 think "small"



interested in
$$\frac{\# \text{ blocks } (\text{coal}_p(B_n))}{\# \text{ blocks } (B_n)} = \frac{n - Y_n + \mathbb{1}(Y_n > 0)}{n}$$

$$= \frac{n - (Y_n - 1)^+}{n}$$
 where $Y_n = \text{Bin}(n, p)$

$$\log \left(\frac{n - Y_n + \mathbb{1}(Y_n > 0)}{n} \right) = - \frac{Y_n + \mathbb{1}(Y_n > 0)}{n} + \mathcal{O} \left(\left(\frac{Y_n - 1}{n} \right)^2 \right)$$

if $\frac{Y_n}{n} \leq 1$ will be true with overwhelming probab. if $p \leq 1 - \epsilon$

generalisation:

$$\vec{p} \in \Delta = \{ (x_1, x_2, \dots) : x_i \geq x_{i+1}, \dots \geq 0, \sum x_i \leq 1 \}$$

$\text{coal}_{\vec{p}}(B_n) =$ draw indep. color according to \vec{p} for each block, simultaneously collapse all blocks with same color into 1 $\mathbb{P}(\text{color } i) = p_i$

$$\# \text{ blocks } (\text{coal}_{\vec{p}}(B_n)) = \frac{n - \sum_{i=1}^{\infty} Y_n^i + \mathbb{1}(Y_n^i > 0)}{n}$$

where $Y_n^i = \# \text{ blocks with color } i$

again
$$\log \frac{\# \text{ blocks } (\text{coal}_{\vec{p}}(B_n))}{n} = - \frac{\sum Y_n^i - \mathbb{1}(Y_n^i > 0)}{n} - \frac{1}{2} \left(\frac{\sum (Y_n^i - 1)^+}{n} \right)^2 - \frac{1}{3} \left(\frac{\sum (Y_n^i - 1)^+}{n} \right)^3 + \dots$$

Moreover if $\sum p_i \leq 1 - \epsilon < 1$, then uniformly over such \vec{p}

$$\log \frac{\# \text{blocks}(\text{coal}_{\vec{p}}(B_n))}{n} = - \frac{\sum_i Y_n^i - \mathbb{1}(Y_n^i > 0)}{n} + O_{\epsilon} \left(\frac{(\sum_i Y_n^i - 1)^2}{n} \right)$$

$$\mathbb{E} \left(\frac{\sum_{i=1}^{\infty} n p_i - 1 + (1-p_i)^n}{n} \right) + O_{\epsilon} \left((\sum p_i)^2 \right)$$

Since $(1-p_i)^n = e^{-n p_i} + O(n p_i^2)$

$$= - \frac{\sum_{i=1}^{\infty} e^{-n p_i} - 1 + n p_i}{n} + O_{\epsilon} \left((\sum p_i)^2 + \sum p_i^2 \right)$$

⊖-coalescent

⊖ a probab. measure on Δ , $\mathbb{E}(\sum \vec{0}) = 0$ (assume in this talk)

B_s = configuration of blocks at time s , $s > 0$

\vec{p} arrives at rate $\frac{\mathbb{E}(d_{\vec{p}})}{\sum p_i^2}$, "upon" each arrival perform $\text{coal}_{\vec{p}}(B_{\text{current time}})$

So if (t, \vec{p}) is an atom of PPP then $B_t = \text{coal}_{\vec{p}}(B_{t-})$

Notes: 1) $\frac{\mathbb{E}(d_{\vec{p}})}{\sum p_i^2}$ may not be finite, but the construction still works (recall Lévy processes)

2) can start with $B_{0+} =$ infinitely many blocks, standard version
 due to consistency: any b -tuple of blocks behaves as \ominus -coalescent started from b blocks

3) Any pair of blocks coalesces at rate 1

$$\lambda_2 = \int P(\text{1 and 2 colored by same color}) \frac{\mathbb{E}(d_{\vec{p}})}{\sum p_i^2} = 1 \quad (\mathbb{E} \vec{p} \text{ is p.m.})$$

4) if $\mathbb{E}(\vec{p} : p_2 > 0) = 0$ $\xrightarrow{\sum_i p_i^2}$ λ -coalescent $\text{coal}_{\vec{p}}(B_{\cdot})$

(3)

Suppose for some $\epsilon > 0$ $\int_{\{\vec{p} \in \Delta : \sum p_i > 1 - \epsilon\}} \frac{\Theta(d\vec{p})}{\sum p_i^2} < \infty$ then if (C_ϵ) for some $\epsilon > 0$.

$T_\epsilon := \min\{\text{arrival time of } \vec{p} \text{ s.t. } \sum p_i > 1 - \epsilon\}$ hence $P(T_\epsilon = 0) = 0$

continuous dependent on n .

Fix large n , interested in the evolution of

(# blocks $(B_s^{(n)})$, $s \geq 0$) and

(# blocks $(B_s^{(\infty)})$, $s \geq 0$) — standard case

Def Say that the standard coalescent B CDI if

$$P(\# \text{ blocks } (B_t) < \infty, \forall t \geq 0) = 1$$

Falt Schweisberg

If $P(\vec{p} : \sum_{i=1}^n p_i = 1, \exists \text{ finite}) = 0$, then either

standard Θ -coal CDI or $P(\# \text{ blocks } (B_t) = \infty, \forall t \geq 0) = 1$

From now on assume

$$P(\vec{p} : \sum_{i=1}^n p_i = 1, \exists n < \infty) = 0$$

and moreover $P(\vec{p} : \sum p_i \leq 1 - \epsilon) = 1$ for some $\epsilon > 0$.

Consider $(B_s^{(n_0)}, s \geq 0)$, n_0 large finite

$$N^{(n_0)}(s) = \# \text{ blocks } (B_s^{(n_0)}), s \geq 0$$

$$\int_{N^{(n_0)}(t)=n} \mathbb{E}(d \log N^{(n_0)}(t) | \mathcal{F}_t) = \int_{\Delta} \left(\frac{\sum_{i=1}^{\infty} (e^{-np_i} - 1 + np_i)}{n} \frac{\Theta(d\vec{p})}{\sum p_i^2} + \int_{\Delta} \left[\Theta_\epsilon(\sum p_i^2) \frac{\Theta(d\vec{p})}{\sum p_i^2} + \Theta_\epsilon\left(\left(\sum_{i=1}^{\infty} p_i\right)^2\right) \frac{\Theta(d\vec{p})}{\sum p_i^2} \right] dt \right)$$

uniform in n

not used

(C_ε)

④

So provided

$$(R) \quad \int \frac{\left(\sum_{i=1}^{\infty} p_i\right)^2}{\sum_{i=1}^{\infty} p_i^2} \Theta(d\vec{p}) < \infty$$

can write

$$\mathbb{E}(d \log N^{(n)}(t) | \mathcal{F}_t) = \left(- \frac{\Psi(N^{(n)}(t))}{N^{(n)}(t)} + \lambda(t) \right) dt$$

Note: Always holds for

→ can write with ~~...~~

where

$$\Psi(q) = \int \sum_{i=1}^{\infty} (e^{-2p_i} - 1 + 2p_i) \frac{d\Theta(\vec{p})}{\sum_{i=1}^{\infty} p_i^2} \quad \left| \text{bounded by constant} \right.$$

$$(R) \Rightarrow (C_\varepsilon) \Delta$$

- Note
- 1) (R) is true for all \rightarrow -coalescents ~~...~~
 - 2) (R) is not true for all \ominus -coalescents

exple $\Theta\left(\underbrace{\left(\frac{1}{2^n}, \frac{1}{2^n}, \dots, \frac{1}{2^n}, 0, 0, \dots\right)}_{2^{n-1}, \text{ or } 2^n - 2, \dots}\right) = \frac{1}{2^n}, n \geq 1$

$$\sum p_i^2 = \sum \frac{1}{2^{2n}} = \frac{1}{2^n} \quad \Theta(\vec{p}) = \frac{1}{2^n}$$

$$\sum p_i \approx 1$$

3) probabilistic interpretation of (R):

- if in the PPP graphical construction using $\text{Coal}_p(\mathbb{R}_+^n)$ one becomes "color blind" - all blocks are colored by any p_1, p_2, \dots are merged together in one block, the corresponding λ -coalescent "exists" in the sense that $\# \text{blocks}(\text{collapse}(\mathbb{R}))_{0+} = \infty$ iff (R), and if not (R) $\# \text{blocks}(\text{collapse}(\mathbb{R})) = 1$ ~~...~~

- 4) could have CDI + (R)
- CDI + ∇ (R)
- non CDI + (R)
- non CDI + ∇ (R)

(5)

no precise (suff. - nec. criterion for CDI) but
 Theorem of Schweierberg says that if (C_2) then

CDI $\Leftrightarrow \int \frac{dq}{\psi(q)} < \infty$ (in fact his criterion in terms of quantities that behave asympt. the same as ψ , also summability)

And also $\int \frac{dq}{\psi(q)} < \infty \Rightarrow$ CDI

but there are exmples of CDI + $\int_{(0, \infty)} \frac{dq}{\psi(q)} = \infty$, of course not satisfying (C_2) , $\forall \epsilon > 0$

Theorem If (R) then the standard \mathbb{P} coalescent has speed of coming down from infinity

$\lim_{t \rightarrow \infty} \frac{N^{\mathbb{P}}(t)}{v_{\mathbb{P}}(t)} = 1$ a.s.

where $v_{\mathbb{P}}(t)$ solves $\int_{v_{\mathbb{P}}(t)}^{\infty} \frac{dq}{\psi(q)} < \infty$

Why? Similar to the work Berestycki x 2, 1 for Δ -coalescents

use $\log \frac{N^{\mathbb{P}}(t)}{v_{\mathbb{P}}(t)} - \log \frac{N^{\mathbb{P}}(z)}{v_{\mathbb{P}}(z)} + \int_z^t \left[\frac{\psi(N(u))}{N(u)} - \frac{\psi(v_{\mathbb{P}}(u))}{v_{\mathbb{P}}(u)} + H(u) \right] du$
 a MG $O(1)$

$\text{var}(\log N(t) | \mathcal{F}_z) = \int_z^t H_1(u) du$
 again $\leq C$ uniformly in $N(t) \geq m_0$

and $\lim_{q \rightarrow \infty} \frac{\psi(q)}{q} \uparrow$ funct. —

X_z chosen so that $v_{X_z}^{\mathbb{P}}(z) = N(z) = v(X_z + z)$

(6)

or equivalently (R) will guarantee that

$\log \frac{N^{(n)}(t)}{v_n(t)}$ stays close to 0 for some order 1 amount of time uniformly over n

where $v_n(t) = v_{\Theta}(t_n + t)$ and $v_{\Theta}(t_n) = n$

Why $\frac{d}{dt} \log v_n(t) = (\log v_n(t))' = \frac{1}{v_n(t)} v_n'(t) = - \frac{\psi(v_n(t))}{v_n(t)}$

$v_{\Theta}(t_n + t)$ s.t. $\int_{v_{\Theta}(t_n + t)}^{\infty} \frac{dq}{\psi(q)} = t_n + t$

$-\frac{1}{\psi(v_{\Theta}(t_n + t))} v_{\Theta}'(t_n + t) = 1$

$\int \frac{dq}{\psi(q)} < \infty$	CDI	(R)	
0	0	0	✓
0	0	1	✓ \rightarrow bad
0	1	0	✓
0	1	1	✗
1	0	0	✗
1	0	1	✗
1	1	0	✓
1	1	1	✓