ON ASYMPTOTICS OF EXCHANGEABLE

COALESCENTS WITH MULTIPLE COLLISIONS

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CIRM, Marseille Luminy, May 26, 2009

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Coalescent with multiple collisions (Pitman 1999, Sagitov 1999)

Fix $n \in \mathbb{N}$. Let Λ be a finite measure on [0, 1].

Definition. (Coalescent with multiple collisions, Λ -coalescent)

The n-coalescent with multiple collisions is a Markov process $\Pi_n = (\Pi_n(t))_{t\geq 0}$ with state space \mathcal{P}_n , the set of partitions of $\{1, \ldots, n\}$, and generator $Q = (q_{\pi\pi'})_{\pi,\pi'\in\mathcal{P}_n}$ with rates

$$q_{\pi\pi'} = \int_{[0,1]} x^{k-2} (1-x)^{b-k} \Lambda(dx) =: \lambda_{b,k}$$
 for $\pi \prec_k \pi', k \ge 2$, where

 $b:=|\pi|:=\text{number of blocks of }\pi$

 $\pi \prec_k \pi' :\iff$ Exactly k blocks of π merge to form a single block of π'

Examples.

 $\Lambda = \delta_0$ (Dirac measure in 0). \Rightarrow Kingman coalescent (1982)

 $\Lambda = U_{[0,1]}$ (uniform on [0,1]). \Rightarrow Bolthausen-Sznitman coalescent (1998)

 $\Lambda = \delta_1$ (Dirac measure in 1). \Rightarrow star-shaped coalescent

Functionals of coalescent processes



 $\tau_n := \inf\{t \ge 0 : |\Pi_n(t)| = 1\} =$ time back to MRCA = absorption time



Graphical representation of three realisations of a 6-coalescent with multiple collisions

Example: Bolthausen-Sznitman coalescent

Proposition 1. (Asymptotics of X_n , Iksanov and M. 2007, Drmota, Iksanov, Rösler and M. 2007, 2009)

$$Y_n := \frac{(\log n)^2}{n} X_n - \log n - \log \log n \quad \stackrel{d}{\to} Y$$

where Y is 1-stable with characteristic function $E(e^{i\lambda Y}) = e^{i\lambda \log |\lambda| - \frac{\pi}{2}|\lambda|}$, $\lambda \in \mathbb{R}$.

Proposition 2. (Asymptotics of τ_n , Goldschmidt and Martin 2005, Freund and M. 2007)

$$\tau_n - \log \log n \xrightarrow{d} \tau$$

where au is Gumbel distributed.

Remark. Analytic proofs are based on singularity analysis. Probabilistic proofs use relations to random recursive trees and/or to random walks with a barrier.

Summary: Number of collisions of beta(a,1)-coalescent (2007/2008)

coalescent	parameter a	number $\mathbf{X_n}$ of collisions
Kingman	$a \rightarrow 0$	$X_n = n - 1$
	0 < a < 1	$\frac{X_n - n(\alpha - 1)}{(\alpha - 1)n^{1/\alpha}} \xrightarrow{d} Y_\alpha \text{(α-stable, α := 2 - a$)}$
Bolthausen-Sznitman	a = 1	$\frac{(\log n)^2}{n} X_n - \log(n \log n) \xrightarrow{d} Y (1-\text{stable})$
	1 < a < 2	$\frac{X_n}{\Gamma(a)n^{2-a}} \xrightarrow{d} \int_0^\infty e^{-U_t} dt$ ($U = $ subord.)
	a=2	$\frac{X_n - (2m_1)^{-1} \log^2 n}{(\frac{m_2}{3m_1^2} \log^3 n)^{1/2}} \xrightarrow{d} N \text{(standard normal)}$
	$2 < a < \infty$	$\frac{X_n - \mu^{-1} \log n}{(\sigma^2 \mu^{-3} \log n)^{1/2}} \xrightarrow{d} N \text{(standard normal)}$
star-shaped	$a \to \infty$	$X_n = 1$

Intuition: High mass of Λ near $0 \cong \Pi_n$ has many small jumps \cong increase of X_n

Assumptions

Assume that $\nu(dx) := x^{-2} \Lambda(dx)$ is a probability measure on (0,1) such that

1. the support of ν is not contained in $\{1-\delta\gamma^n:n=0,1,\ldots\}$ for some $\delta>0$ and $\gamma\in(0,1),$

2.
$$\int_{(0,1)} |\log x| \nu(dx) < \infty$$
.

The more general case of a finite measure ν can be reduced to the case of a probability measure by a linear time change of the coalescent.

The annihilator

There is a simpler process $A_n := (A_n(t))_{t \ge 0}$ with state space $\{0, \ldots, n\}$ moving from m to m - k with rate $\binom{m}{k} \lambda_{m,k}$, $1 \le k \le m \le n$.

Remark. Total rates are
$$\sum_{k=1}^{m} {m \choose k} \lambda_{m,k} = \int_{[0,1]} (1 - (1 - x)^m) \nu(dx), m \in \{0, \dots, n\}.$$

Interpretation.

When there are m particles, any k-tuple of them collides and annihilates at rate $\lambda_{m,k}$.

Pitman's Poisson process construction.

At the generic time s_j of a unit Poisson process, a random variable $1 - \eta_j$ is sampled from ν , and each of the remaining particles is marked 'head' with probability $1 - \eta_j$ or 'tail' with probability η_j . The particles marked 'head' are removed.

A coupling

Define the coalescent Π_n and the annihilator A_n using the same unit Poisson process and the same sample $1 - \eta_1, 1 - \eta_2, \ldots$ from ν .

Call the initial *n* particles primary and their followers resulting from mergers secondary.

 $K_n :=$ number of transitions of A_n as the process proceeds from n to 0

 $K_{n,1} :=$ number of jumps of A_n of size 1

 $\sigma_n := \inf\{t \ge 0 : A_n(t) = 0\}$ (absorption time, first time when there are no primary particles anymore)

 $K_{n,0} :=$ number of epochs $s_j < \sigma_n$ when none of the primary particles are marked 'head' $U_n := |\Pi_n(\sigma_n)|$ (remaining secondary particles) Lemma 1. The following stochastic order relations hold:

$$K_n - K_{n,1} \leq_d X_n \leq_d K_n + K_{n,0} + X_{U_n}$$

where (on the r.h.s.) $(K_n, K_{n,0}, U_n)$ is independent of (X_1, X_2, \ldots) .

Proof. Coupling. If ≥ 2 primary particles collide, then both processes A_n and Π_n jump. The number of jumps of Π_n up to time σ_n does not exceed the number $K_n + K_{n,0}$, and after time σ_n the coalescent evolves with U_n particles.

Lemma 2. The sequence of distributions of X_{U_n} , $n \in \mathbb{N}$, is tight.

Sketch of proof. Use $X_{U_n} < U_n$ and study the number Q_j of secondary particles at time s_j . Given all the η_k 's, $(Q_j)_j$ is a Markov chain. Introduce $N_n := \inf\{k \ge 1 : \eta_1 \cdots \eta_k \le 1/n\}$ and use renewal theory to show that $(Q_{N_n})_n$ is tight. Now replace the fixed drop level 1/nby the appropriate random drop level associated with the last primary particle disappearing at time σ_n in order to show that $|Q_{N_n} - U_n|$ is stochastically bounded (technical ...)

Number of collisions X_n : Asymptotics

Theorem 1. The following two assertions are equivalent.

- (i) There exist constants $a_n > 0$ and $b_n \in \mathbb{R}$ such that $(X_n a_n)/b_n$ weakly converges to some non-degenerate and proper distribution.
- (ii) The distribution of $(-\log \eta)$ either belongs to the domain of attraction of a stable law or the function $x \mapsto P(-\log \eta > x)$ slowly varies at ∞ .

There are five different regimes (A) - (E) of convergence.

Possible limiting laws are normal, α -stable with $\alpha \in [1, 2)$ and (scaled) Mittag-Leffler.

Convergence Regimes

Logarithmic moments: $\mu := E(-\log \eta)$ and $\sigma^2 := Var(\log \eta)$

(A) If $\sigma^2 < \infty$ then, with $b_n := \mu^{-1}$ and $a_n := (\mu^{-3}\sigma^2 \log n)^{1/2}$, the limiting law of $(X_n - a_n)/b_n$ is standard normal.

(B) If $\sigma^2=\infty$ and

$$\int_{(x,1)} (\log y)^2 \nu(dy) \sim L(-\log x) \quad \text{as } x \to 0$$

for some L slowly varying at ∞ , then, with $b_n := \mu^{-1} \log n$ and $a_n := \mu^{-3/2} c_{[\log n]}$, where c_n is any sequence satisfying $\lim_{n\to\infty} nL(c_n)/c_n^2 = 1$, the limiting law of $(X_n - b_n)/a_n$ is standard normal. (C) Assume that, for some function L slowly varying at ∞ , the relation

$$P(\eta \le x) \sim (-\log x)^{-\alpha} L(-\log x) \quad \text{as } x \to 0 \tag{(*)}$$

holds with $\alpha \in [1,2)$ and that $\mu < \infty$ if $\alpha = 1$. Then, with $b_n := \mu^{-1} \log n$ and $a_n := \mu^{-(\alpha+1)/\alpha} c_{[\log n]}$, where c_n is any sequence satisfying $\lim_{n\to\infty} nL(c_n)/c_n^{\alpha} = 1$, the limiting law of $(X_n - a_n)/b_n$ is α -stable with characteristic function

$$t \mapsto \exp\left\{-|t|^{\alpha}\Gamma(1-\alpha)\left(\cos\left(\frac{\pi\alpha}{2}\right)+i\sin\left(\frac{\pi\alpha}{2}\right)\operatorname{sgn}(t)\right)\right\}, \quad t \in \mathbb{R}.$$

The remaining two regimes (D) and (E) cover the case when $\mu = \infty$.

(D) Assume that $\mu = \infty$ and that (*) in (C) holds with $\alpha = 1$. Let c be any positive function satisfying $\lim_{x \to \infty} \frac{xL(c(x))}{c(x)} = 1$, and set $\psi(x) := x \int_{e^{-c(x)}}^{1} \frac{P(\eta \le y)}{y} dy$. Let b be any positive function satisfying $b(\psi(x)) \sim \psi(b(x)) \sim x, x \to \infty$. Then, with $b_n := b(\log n)$ and $a_n := b(\log n)c(b(\log n))/\log n$, the limiting law of $(X_n - a_n)/b_n$ is 1-stable with characteristic function

$$t \mapsto \exp\left(-|t|\left(\frac{\pi}{2}-i\log|t|\operatorname{sgn}(t)\right)\right), \quad t \in \mathbb{R}$$

The last regime:

(E) If (*) in (C) holds with $\alpha \in [0, 1)$ then, with $a_n := \frac{\log^{\alpha} n}{L(\log n)}$, the limiting law of $\frac{X_n}{a_n}$ is scaled Mittag-Leffler θ_{α} (exponential if $\alpha = 0$) with moments

$$\int_{[0,\infty)} x^n \,\theta_\alpha(dx) = \frac{n!}{\Gamma^n(1-\alpha)\Gamma(1+n\alpha)}, \qquad n \in \mathbb{N}$$

Basic ideas of the proof

- Use the coupling $K_n K_{n,1} \leq_d X_n \leq_d K_n + K_{n,0} + X_{U_n}$ and the tightness of X_{U_n} to show that $(X_n b_n)/a_n$ weakly converges if and only if $(K_n b_n)/a_n$ weakly converges to the same distribution.
- Interpret K_n as the number of non-empty boxes in an occupancy scheme, where n balls are dropped in box 1 with probability $\xi_1 := 1 - \eta_1$, the remaining balls are dropped in box 2 with probability $\xi_2 := 1 - \eta_2$ and so on (Bernoulli sieve).

(Gnedin, Iksanov, Negadajlov, Rösler 2008)

Example 1. Number of collisions for beta coalescents

An application of Regime (A) to beta coalescents with

$$\nu(dx) = \frac{\Lambda(dx)}{x^2} = cx^{a-3}(1-x)^{b-1}dx$$

(a > 2 and b, c > 0) shows that

$$\frac{X_n - \mu^{-1} \log n}{(\sigma^2 \mu^{-3} \log n)^{\frac{1}{2}}} \stackrel{d}{\to} N(0, 1)$$

where

$$\mu = \Psi(a - 2 + b) - \Psi(b),$$

$$\sigma^{2} = \Psi'(b) - \Psi'(a - 2 + b)$$

and Ψ is the logarithmic derivative of the gamma function.

Example 2

Suppose that
$$P(\eta \le x) = \frac{1}{(1 - \log x)^{\alpha}}$$
, $x \in (0, 1)$, for some $\alpha > 0$.

Then all five regimes occur.

Parameter α	$\mu := \mathcal{E}(-\log \eta)$	$\sigma^2 := \operatorname{Var}(\log \eta)$	Regime
$0 < \alpha < 1$	∞	∞	(E) Mittag-Leffler
1	∞	∞	(D) 1-stable
$\boxed{1 < \alpha < 2}$	$\frac{1}{\alpha - 1}$	∞	(C) α -stable
2	1	∞	(B) normal
$\boxed{2 < \alpha < \infty}$	$\frac{1}{\alpha - 1}$	$\frac{\alpha}{(\alpha-1)^2(\alpha-2)}$	(A) normal

Absorption times

Lemma. The following relation holds:

$$\tau_n \stackrel{d}{=} \sigma_n + \tau'_{U_n}$$

where, on the right hand side, $\tau'_j \stackrel{d}{=} \tau_j$, $j \in \mathbb{N}$, and (σ_n, U_n) and $(\tau'_j)_j$ are independent.

Proof. Coupling. When all the primary particles disappear at time σ_n , there are U_n secondary particles left.

Absorption times: Asymptotics

Theorem 2. Assume that condition (ii) of Theorem 1 holds. Then, with X_n replaced by τ_n , condition (i) of Theorem 1 holds.

There are again five different regimes of convergence.

Sketch of proof. By the previous lemma and Lemma 2, σ_n and τ_n have the same asymptotics. σ_n has the same limiting law as the first passage time through level $\log n$ for a compound Poisson process with the generic jump $-\log \eta$ and intensity 1.

Example: Absorption time for beta coalescent

An application of Regime (A) to beta coalescents with

$$\nu(dx) = \frac{\Lambda(dx)}{x^2} = cx^{a-3}(1-x)^{b-1}dx$$

(a > 2 and b, c > 0) shows that

$$\frac{C\tau_n - \mu^{-1}\log n}{((\sigma^2 + \mu^2)\mu^{-3}\log n)^{\frac{1}{2}}} \stackrel{d}{\to} N(0,1)$$

with μ and σ^2 as before and $C := \frac{a-1+b}{a-1} \frac{a-2+b}{a-2}$.

Remarks.

 \circ The factor C comes from the linear time change (ν is in general not a probability measure).

• Devroye (1999) obtains the same scaling for the depth of a random node in random trees.

Summary: Absorption time of beta(a,1)-coalescent

coalescent	parameter a	absorption time $ au_{\mathbf{n}}$
Kingman	$a \rightarrow 0$	$ au_n \stackrel{\text{a.s.}}{\to} au, \mathbf{E}(au) = 2$
	0 < a < 1	$\tau_n \stackrel{\text{a.s.}}{\to} \tau, E(\tau) < \infty$
Bolthausen-Sznitman	a = 1	$ au_n - \log \log n \xrightarrow{d} au$ (Gumbel)
	1 < a < 2	Conjecture: asymptotically normal?
	a=2	Conjecture: asymptotically normal?
	$2 < a < \infty$	$\frac{C\tau_n - \mu^{-1}\log n}{((\sigma^2 + \mu^2)\mu^{-3}\log n)^{1/2}} \xrightarrow{d} N \text{(standard normal)}$
star-shaped	$a \to \infty$	$\tau_n \stackrel{d}{=} \operatorname{Exp}(1)$

External branch length

 $Z_n :=$ length of an external branch, chosen uniformly at random from the *n* external branches of the coalescent Π_n

Let T_n be the time of the first jump of Π_n and let $I_n := |\Pi_n(T_n)|$ the number of blocks after that first jump.

 Z_n satisfies the distributional recursion

$$Z_1 = 0$$
 and $Z_n = T_n + B_n Z_{I_n}, \quad n \in \{2, 3, \ldots\},$

where T_n is independent of $B_n Z_{I_n}$, B_2, B_3, \ldots are Bernoulli random variables with $P(B_n = 1|I_n) = (I_n - 1)/n$, and, conditional on I_n , B_n and Z_{I_n} are independent.

Summary: External branch length of beta(a,1)-coalescent

coalescent	parameter a	external branch length $\mathbf{Z_n}$	
Kingman	$a \rightarrow 0$	$nZ_n \xrightarrow{d} Z$ having density $x \mapsto 8/(2+x)^3$	
	0 < a < 1	Conj.: $n^{1-a}Z_n \xrightarrow{d} Z = ?$ (work in progress?)	
Bolthausen-Sznitman	a = 1	$(\log n)Z_n \xrightarrow{d} Z \stackrel{d}{=} \operatorname{Exp}(1)$	
	$1 < a < \infty$	$Z_n \xrightarrow{d} Z \stackrel{d}{=} \operatorname{Exp}(a/(a-1))$	
star-shaped	$a \to \infty$	$Z_n \stackrel{d}{=} \operatorname{Exp}(1)$	

Extensions to Xi-coalescents without proper frequencies

Some results can be easily extended to Ξ -coalescents (with simultaneous multiple collisions). For example, $P(Z_n > t)$ is the probability that a randomly chosen individual is still a singleton at time t. By exchangeability, we can call this individual '1', and it follows that

 $P(Z_n > 1) = P(\{1\} \text{ is a singleton block of } \Pi_n(t)).$

If the coalescent does no have proper frequencies, i.e., if the measure Ξ on the simplex

$$\begin{split} \Delta &:= \{x = (x_1, x_2, \ldots) : x_1 \ge x_2 \ge \cdots \ge 0, |x| := \sum_i x_i \le 1\}\\ \text{satisfies} \boxed{\alpha := \int_{\Delta} \frac{|x|}{(x, x)} \Xi(dx) < \infty}, \text{ where } (x, x) := \sum_i x_i^2, \text{ then,}\\ \lim_{n \to \infty} P(Z_n > t) \ = \ P(\{1\} \text{ is a block of } \Pi_{\infty}(t)) \ = \ \exp(-\alpha t), \end{split}$$

using an argument of Pitman (1999), Eq. (37).

Examples. Dirac-coalescents, Poisson-Dirichlet coalescent

Discussion: A more general approach

Suppose that $\alpha < \infty$. For $t \ge 0$ let S_t denote the frequency of singletons of $\Pi_{\infty}(t)$ and set $X_t := -\log S_t$. Then $X = (X_t)_{t>0}$ is a subordinator with Laplace exponent

$$\Phi(a) = \int_{\Delta} (1 - (1 - |x|)^a) \frac{\Xi(dx)}{(x, x)}, \qquad a \ge 0.$$

We can still define the annihilator using Schweinberg's Poisson construction or Kingman's paint-box construction (drop all particles marked head). The annihilator has total rates $\Phi(m)$, $m \in \{0, \ldots, n\}$, and the inequality $X_n \ge_d K_n - K_{n,1}$ holds true. The other bound in general does not make sense as a variable analogous to $K_{n,0}$ might be ∞ . It seems plausible that, under additional conditions, X_n and $K_n - K_{n,1}$ still have the same limiting law.

Remark. This approach is for example also useful to analyse the asymptotics of the number of types M_n for Ξ -coalescents without proper frequencies and with mutation (rate r > 0): $\boxed{\frac{M_n}{n} \stackrel{d}{\rightarrow} \int_0^\infty r e^{-rt} S_t \, dt}$ (exponential integral of a subordinator with drift)

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