Occupation time fluctuations of particle systems

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joint work with T. Bojdecki and L. Gorostiza

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Particle system Problem

Particle system in \mathbb{R}^d

- initial distribution of particles given by a Poisson random measure with intensity measure ν (or ν_T)
- particle motion standard, spherically symmetric α-stable Lévy process
- intensity of branching V > 0,

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Particle system Problem

Particle system in \mathbb{R}^d

 (1 + β) branching mechanism, 0 < β ≤ 1, with probability generating function

$$g(s) = s + \frac{1}{1+\beta}(1-s)^{1+\beta}, \ 0 < s < 1.$$

branching is critical – mean value = 1

- if $\beta = 1 \text{binary branching}$
- if $\beta < 1$ infinite variance branching
- particles move and branch independently

Empirical process

 N_t : $N_t(A)$ = number of particles in the set A at time t.

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Model

Particle syste Problem

Problem

$$\int_0^{Tt} N_s \qquad ds, \qquad t \ge 0$$

Anna Talarczyk Occupation time fluctuations

Particle system Problem

Problem

Rescaled occupation time fluctuations

$$X_T(t) = rac{1}{F_T}\int_0^{T_t} (N_s - EN_s) ds, \qquad t \geq 0$$

where F_T is a proper norming.

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Particle system Problem

Problem

Rescaled occupation time fluctuations

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where F_T is a proper norming.

Problem

- find a suitable norming *F_T*, such that *X_T* converges in law as *T* → ∞ to a nontrivial limit
- identification of the limit
- properties of the limit

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Two cases

- initial intensity measure $= \nu$ does not depend on *T* (only speed up the time)
- ② initial intensity measure = $\nu_T = H_T \nu$, $H_T \rightarrow \infty$ (speed up the time + high density)

We consider the following initial intensity measures

- $\nu = \lambda$ (Lebesgue measure)
- $H_T\mu$, where μ finite measure, $H_T \rightarrow \infty$

• $H_T \mu_\gamma$ with

$$\mu_{\gamma}(dx) = rac{dx}{1+|x|^{\gamma}}, \quad \gamma \geq 0$$

either $H_T \equiv 1$ or $H_T \rightarrow \infty$ sufficiently quickly

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Particle system Problem

History and related problems

- Deuschel, Wang (1994), $\alpha = 2$ (Brownian particles), Lebesgue measure, without branching
- Dawson, Gorostiza, Wakolbinger (2001), convergence of random variable $X_T(1)$, $\beta = 1$, general β only for "large" dimensions
- Iscoe (1986), for superprocesses, Lebesgue, convergence of X_T(1).
- Birkner, Zähle (2007), for branching random walks on a lattice.

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Particle system Problem

Notation

- S(ℝ^d) the Schwartz space of smooth quickly decreasing functions
 S'(ℝ^d) the space of tempered distributions on ℝ^d,
 - It is convenient to consider X_T as a process in $\mathcal{S}'(\mathbb{R}^d)$
- p_t the transition density of the α -stable process, $(\hat{p}_t(x) = e^{-t|x|^{\alpha}})$
 - T_t transition semigroup $T_t f = p_t * f$,
- \Rightarrow_{c} convergence in law in $C([0, \tau], S'(\mathbb{R}^{d}))$, for any $\tau > 0$ \Rightarrow_{f} - convergence of finite dimensional distributions

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Particle system Problem

Expected value of the occupation time

By the Poisson property of the initial distribution: $E \langle N_t, \varphi \rangle = \int_{\mathbb{R}^d} \mathcal{T}_t \varphi(x) \nu(dx)$

Proposition 1

Initial intensity measure $\nu = \mu_{\gamma} = \frac{dx}{1+|x|^{\gamma}}$,

$$E \int_0^T \langle N_t, \varphi \rangle \, dt \sim \begin{cases} T^{1-\frac{\gamma}{\alpha}} & \text{if} \quad \gamma < d, \gamma < \alpha \\ \log T & \gamma = \alpha < d \\ T^{1-\frac{d}{\alpha}} \log T & \gamma = d < \alpha \\ (\log T)^2 & \gamma = d = \alpha \\ T^{1-\frac{d}{\alpha}} & \gamma > d, d < \alpha \\ \log T & \gamma > d, d = \alpha \\ 1 & \gamma > \alpha, d > \alpha \end{cases}$$

Lebesgue measure ($\gamma = 0$) General case ($gamma \ge 0$) Long range dependence

stable random measure

Let *M* be independently scattered $(1 + \beta)$ -stable random measure on \mathbb{R}^{d+1} with control measure λ_{d+1} (Lebesgue) totally skewed to the right, i.e. for any $A \in \mathcal{B}(\mathbb{R}^{d+1})$ such that $0 < \lambda_{d+1}(A) < \infty$, M(A) is $(1 + \beta)$ -stable, with characteristic function

$$\exp\left\{-\lambda_{d+1}(A)|z|^{1+\beta}\left(1-i(\operatorname{sgn} z)\tan\frac{\pi}{2}(1+\beta)\right)\right\},$$

$$z \in \mathbb{R},$$

M is σ -additive and $M(A_j)$, j = 1, 2, ... are independent if A_j are disjoint.

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Model	Lebesgue measure ($\gamma = 0$)
Results	General case ($gamma \ge 0$)
Idea of the proof	

Theorem 2

Assume that the initial intensity measure is $\nu = \lambda$ ($\gamma = 0$), then (i) $\frac{\alpha}{\beta} < d < \frac{\alpha(1+\beta)}{\beta}$, $F_T^{1+\beta} = T^{(2+\beta-\frac{d}{\alpha}\beta)}$, then $X_T \rightleftharpoons K\lambda\xi$, where $\xi_t = \int_{\mathbb{R}^{d+1}} \left(\mathbf{1}_{[0,t]}(r) \int_r^t p_{u-r}(x) du \right) M(drdx), t \ge 0,$

K is a constant.

 ξ is well defined for $d < \frac{\alpha(1+\beta)}{\beta}$, since

$$\int_{\mathbb{R}^d} \left(\int_0^\tau p_u(x) du \right)^{1+\beta} dx < \infty, \quad \text{for any } \tau > 0.$$

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Model	Lebesgue measure ($\gamma=0$)
Results	General case ($gamma \ge 0$)
Idea of the proof	

$$\xi_t = \int_{\mathbb{R}^{d+1}} \left(\mathbf{1}_{[0,t]}(r) \int_r^t p_{u-r}(x) du \right) M(drdx), t \ge 0.$$

Remark: If $\beta = 1$ (binary branching), then ξ is a centered Gaussian process with covariance function

$$s^{h} + t^{h} - \frac{1}{2} \left[(s+t)^{h} + |s-t|^{h} \right],$$

 $h = 3 - d/\alpha$ (0 < h < 2) (sub-fractional Brownian motion)

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Model	Lebesgue measure ($\gamma=$ 0)
Results	General case ($gamma \ge 0$
dea of the proof	Long range dependence

Theorem

(ii)
$$d = \frac{\alpha(1+\beta)}{\beta}$$
, $F_T^{1+\beta} = T \log T$, then $X_T \Rightarrow K\lambda \tilde{\xi}$, where $\tilde{\xi}$ is one-dimensional $(1 + \beta)$ -stable process, totally skewed to the right with stationary independent increments.

$$e^{iz ilde{\xi}_t} = \exp\left\{-t\left|z
ight|^{(1+eta)}\left(1-i(sgnz) anrac{\pi}{2}(1+eta)
ight)
ight\}.$$

ModelLebesgue measure ($\gamma = 0$)ResultsGeneral case (gamma ≥ 0)dea of the proofLong range dependence

(iii) $d > \frac{\alpha(1+\beta)}{\beta}$, $F_T^{1+\beta} = T$, then $X_T \Rightarrow X$, where X is a $(1+\beta)$ -stable process with values in $\mathcal{S}'(\mathbb{R}^d)$, it has stationary independent increments

$$\begin{split} & E \exp\{i\langle X(t),\varphi\rangle\} = \exp\left\{-\mathcal{K}t \int_{\mathbb{R}^d} \left[c_\beta \varphi(x) G\varphi(x) + \frac{V}{2} |G\varphi(x)|^{1+\beta} \left(1 - i(\operatorname{sgn} G\varphi(x)) \tan \frac{\pi}{2}(1+\beta)\right)\right] dx\right\}, \\ & \varphi \in \mathcal{S}(\mathbb{R}^d), t \ge 0, \\ & \text{where} \quad G\varphi(x) = \int_0^\infty \mathcal{T}_s \varphi(x) ds = C_{\alpha,d} \int_{\mathbb{R}^d} \frac{1}{|x-y|^{d-\alpha}} \varphi(y) dy, \\ & c_\beta = \begin{cases} 0 & \text{if } 0 < \beta < 1 \\ 1 & \text{if } \beta = 1. \end{cases} \end{split}$$

Model	Lebesgue measure ($\gamma=$ 0)
Results	General case ($gamma \ge 0$)
ldea of the proof	Long range dependence

Remark: If $d < \frac{\alpha}{\beta}$, $\alpha = 2$, then there is a.s. local extinction, i.e.

 $\lim_{t\to\infty}N_t(A)=0$

a.s. for any bounded $A \in \mathcal{B}(\mathbb{R}^d)$ Hypothesis: true also for $\alpha < 2$.

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Remark: If $d < \frac{\alpha}{\beta}$, $\alpha = 2$, then there is a.s. local extinction, i.e.

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a.s. for any bounded $A \in \mathcal{B}(\mathbb{R}^d)$ Hypothesis: true also for $\alpha < 2$. High density

Theorem 2(i')

Assume that $d \leq \frac{\alpha}{\beta}$, $\nu_T = H_T \lambda$, H_T satisfies

$$\lim_{T\to\infty}H_T^{-\beta}T^{1-\frac{d\beta}{\alpha}}=0,$$

 $F_T^{1+\beta} = H_T T^{(2+\beta-\frac{d}{\alpha}\beta)}$, then $X_T \Rightarrow K_{\lambda}\xi$, where ξ as in Theorem 2 (i).

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Remark. If $\beta = 1$, then ξ is a centered Gaussian process with covariance function

• for $d < \frac{\alpha}{\beta}$

$$-s^{h}-t^{h}+rac{1}{2}[(s+t)^{h}+|s-t|^{h}],$$

where $h = 3 - d/\alpha$ (2 < h < 2 + $\frac{1}{2}$). (This function is positive definite for (2 < h < 4)).

• for
$$d = \frac{\alpha}{\beta}$$

$$\frac{1}{2} \left[(s+t)^2 \log(s+t) + (s-t)^2 \log|s-t| \right] - s^2 \log s - t^2 \log t$$

Introducing high density in Theorem 2 does not change the results, only F_T is different.

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Model Leb Results Ger Idea of the proof Lon

Lebesgue measure ($\gamma = 0$) General case (gamma ≥ 0) Long range dependence

General case ($\gamma \ge 0$)

initial intensity measure: $H_T \frac{dx}{1+|x|^{\gamma}}$, $H_T \to \infty$ sufficiently quickly, in some cases one can take $H_T = 1$ limits: $(1 + \beta)$ -stable processes X dimension *d* spatial structure temporal structure "low" simple complicated $(X = K\lambda \xi)$ (long range dependence) "critical" simple simple $(X = K\lambda \xi)$ (independent increments)

"large"complicatedsimple(space inhomog.)(independent increments)(ξ varies in different cases)

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Model	Lebesgue measure ($\gamma=$ 0)
Results	General case ($gamma \ge 0$)
Idea of the proof	

for low dimensions ξ has the form: $\xi_t = \int_0^t \int_{\mathbb{R}^d} f(t, r, x) M(drdx)$

Critical dimension:

- Lebesgue measure case ($\gamma = 0$): $d = \frac{1+\beta}{\beta}\alpha$
- finite measure case (with high density) ($\gamma > d$): $d = \frac{2+\beta}{1+\beta}$

•
$$0 \le \gamma \le d$$
: $d = \frac{2+\beta}{\beta}\alpha - \frac{\gamma \vee \alpha}{\beta}$

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	Mode Result Idea of the proc	$\begin{array}{llllllllllllllllllllllllllllllllllll$	
of the form $K\lambda\xi$ $X(t) = X(1), t > 0$ Ird \mathbb{I}			
$\gamma < lpha$ $\gamma < \mathbf{d}$	$ \mathbf{f} d < \frac{1+\beta}{\beta} \alpha $ $ \mathbf{F}_{T}^{1+\beta} = \mathbf{H}_{T} T^{2+\beta-\frac{d}{\alpha}\beta-\frac{\gamma}{\alpha}} $	$\mathbf{f}_{T}^{1+\beta} d = \frac{1+\beta}{\beta} \alpha$ $\mathbf{f}_{T}^{1+\beta} = H_{T} T^{1-\frac{\gamma}{\alpha}} \log T$	$\int d > \frac{1+\beta}{\beta} \alpha$ $F_T^{1+\beta} = H_T T^{1-\frac{\gamma}{\alpha}}$
$egin{array}{l} \gamma = lpha \ \gamma < oldsymbol{d} \end{array}$	one can take $H_T = 1$ if $\frac{\alpha}{2} + \gamma < d$	• 4 $d = \frac{1+\beta}{\beta}\alpha$ $F_{\tau}^{1+\beta} = H_{\tau}(\log T)^2$	••• $d > \frac{1+\beta}{\beta}\alpha$ $F_T^{1+\beta} = H_T \log T$
$lpha < \gamma$ $\gamma < oldsymbol{d}$	F ¹ _T	$d = \frac{2+\beta}{\beta}\alpha - \frac{\gamma}{\beta}$ $^{+\beta} = H_T \log T$	$F_T^{1+\beta} = H_T$
$\gamma = d H_1$	$T^{2+\beta-\frac{d}{\alpha}(1+\beta)}\log T F^{1+\beta}_{T}$	$=\frac{2+\beta}{1+\beta}\alpha\\=H_T(\log T)^2$	
$\gamma > d$ F_T^{1+}	$ \stackrel{\bullet}{\stackrel{\circ}{\stackrel{\circ}{\stackrel{\circ}{\stackrel{\circ}{\stackrel{\circ}{\stackrel{\circ}{\stackrel{\circ}{$	$U = \frac{2+\beta}{1+\beta}\alpha$ = $H_T \log T$	

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Occupation time fluctuations

Model	Lebesgue measure ($\gamma = 0$)	
Results	General case ($gamma \ge 0$)	
Idea of the proof		

Case 1: Let
$$\gamma < d$$
, $d < \frac{2+\beta}{\beta} - \frac{\gamma \lor \alpha}{\beta}$, $F_T^{1+\beta} = H_T T^{2+\beta-(d\beta+\gamma)/\alpha}$

$$\lim_{T \to \infty} T^{1-(d-\gamma)\beta/\alpha} H_T^{-\beta} = 0 \quad \text{(only needed if } d < \frac{\alpha}{\beta} + \gamma\text{)}$$

Then $X_T \underset{c}{\Rightarrow} K\lambda \xi$ where ξ

$$\begin{aligned} \xi_t &= \int_{\mathbb{R}^{d+1}} \left(\mathbf{1}_{[0,t]}(r) \left(\int_{\mathbb{R}^d} p_r(x-y) |y|^{-\gamma} dy \right)^{1/(1+\beta)} \\ &\int_r^t p_{u-r}(x) du \right) M(drdx), \quad t \ge 0, \end{aligned}$$

back to the general scheme

Lebesgue measure ($\gamma = 0$) General case (gamma ≥ 0) Long range dependence

Local extinction

Theorem

If $\alpha = 2$ and $d < \frac{\alpha}{\beta} + \gamma$ and the initial intensity measure is μ_{γ} , then the particle system suffers a.s. local extinction, i.e. for each bounded Borel set A

$$P(\lim_{t\to\infty}N_t(A)=0)=1.$$

Hypothesis: the same is true for α < 2

Theorem

Let $0 < \alpha \le 2$, initial intensity measure $= \mu_{\gamma}$, $\gamma < \alpha$ and $d \ge \frac{\alpha}{\beta} + \gamma$. Then there is no almost sure local extinction.

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Model	Lebesgue measure ($\gamma=$ 0)
Results	General case ($gamma \ge 0$)
Idea of the proof	

Note: there is always extinction in probability for $\gamma > 0$ since

$$E \langle N_T, \varphi \rangle \sim \begin{cases} T^{-\frac{\gamma}{\alpha}} & \text{if } \gamma < d \\ T^{-\frac{d}{\alpha}} \log T & \gamma = d \\ T^{-\frac{d}{\alpha}} & \gamma > d \end{cases}$$

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Case 3: Assume
$$\gamma < d, \gamma < \alpha, d > \alpha \frac{1+\beta}{\beta}$$
, and

$$F_T^{1+\beta} = H_T T^{1-\gamma/\alpha},$$

with H_T ≥ 1. Then X_T ⇒_f X, where X is an S'(ℝ^d)-valued (1 + β)-stable process with independent increments
if β < 1 :

$$Ee^{i\langle X(t)-X(s),\varphi\rangle} = \exp\left\{\left(t^{1-\frac{\gamma}{\alpha}} - s^{1-\frac{\gamma}{\alpha}}\right)\right\}$$
$$\int_{\mathbb{R}^d} \frac{V}{2} |G\varphi(x)|^{1+\beta} \left(1 - i(\operatorname{sgn} G\varphi(x)) \tan \frac{\pi}{2}(1+\beta)\right) dx\right\}$$

• if $\beta = 1$, additional term with $\varphi(x)G\varphi(x)$.

back to the general scheme

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Model	Lebesgue measure ($\gamma=$ 0)
Results	General case ($gamma \ge 0$)
ldea of the proof	Long range dependence

Case 2: Assume
$$\gamma < d, \gamma < \alpha, d = \alpha \frac{1+\beta}{\beta}$$
 and

$$F_T^{1+\beta} = H_T T^{1-\gamma/\alpha} \log T,$$

with $H_T \ge 1$. Then $X_T \Rightarrow_f K\lambda\eta$, where η is a $(1 + \beta)$ -stable process with independent, non-stationary increments (for $\gamma > 0$) whose laws are determined by

$$Ee^{iz(\eta_t - \eta_s)} = \exp\left\{\left(t^{1 - \frac{\gamma}{\alpha}} - s^{1 - \frac{\gamma}{\alpha}}\right)|z|^{1 + \beta} \\ \left(1 - i(\operatorname{sgn} z)\tan\frac{\pi}{2}(1 + \beta)\right)\right\}$$

 $z \in \mathbb{R}, t \ge s \ge 0.$

back to the general scheme

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Lebesgue measure ($\gamma = 0$) General case (*gamma* ≥ 0) Long range dependence

Long range dependence

Case 1: $\gamma < d$ and $d < \frac{\alpha(2+\beta)}{\beta} - \frac{\gamma \lor \alpha}{\beta}$. • $\beta = 1$ Gaussian case. Let $0 \le u < v < s < t$. Then

$$Cov(\xi_v - \xi_u, \xi_{t+T} - \xi_{t+T}) \sim CT^{-rac{d}{lpha}}$$

• $\beta < 1$ We introduce dependence exponent: For $0 \le u < v < s < t$, T > 0, $z_1, z_2 \in \mathbb{R}$ let

$$D_{T}(z_{1}, z_{2}; u, v, s, t) = |\log Ee^{i(z_{1}(\xi_{v} - \xi_{u}) + z_{2}(\xi_{T+t} - \xi_{T+s}))} - \log Ee^{iz_{1}(\xi_{v} - \xi_{u})} - \log Ee^{iz_{2}(\xi_{T+t} - \xi_{T+s})}|.$$

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Model	Lebesgue measure ($\gamma=$ 0)
Results	General case ($gamma \ge 0$
dea of the proof	Long range dependence

dependence exponent κ is defined as

$$\kappa = \inf_{\substack{z_1, z_2 \in \mathbb{R} \ 0 \le u < v < s < t \\}} \inf_{sup\{\gamma > 0 : D_T(z_1, z_2; u, v, s, t) = o(T^{-\gamma}) \text{ as } T \to \infty\}.$$

For the process ξ in 1

$$\xi_t = \int_0^t \int_{\mathbb{R}^d} \left(\int_{\mathbb{R}^d} p_r(x-y) |y|^{-\gamma} dy \right)^{\frac{1}{1+\beta}} \int_r^t p_{u-r}(x) du \mathcal{M}(drdx),$$

we have

$$\kappa = \begin{cases} \frac{d}{\alpha} & \text{if} \quad \alpha = 2 \text{ or } \beta > \frac{d - \gamma}{d + \alpha}, \\ \frac{d}{\alpha} \left(1 + \beta - \frac{d - \gamma}{\alpha + d} \right) & \text{if} \quad \alpha < 2 \text{ and } \beta \le \frac{d - \gamma}{d + \alpha}. \end{cases}$$

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Model	Lebesgue measure ($\gamma=$ 0)
Results	General case ($gamma \ge 0$)
Idea of the proof	Long range dependence

Cases 7 and 9: $\gamma \ge d$, $d < \alpha \frac{2+\beta}{1+\beta}$ general scheme The limit process is of the form $K\lambda\zeta$, where

$$\zeta_t = \int_{\mathbb{R}^{d+1}} \left(\texttt{1}_{[0,t]}(r) p_r^{1/(1+\beta)}(x) \int_r^t p_{u-r}(x) du \right) M(drdx), \quad t \ge 0.$$

Dependence exponent of ζ is

$$\kappa = \frac{d}{\alpha}$$

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Lebesgue measure ($\gamma = 0$) General case (*gamma* ≥ 0) Long range dependence

Other comments

- The same results hold for superprocesses (some differences if $\beta = 1$ in "large" dimensions)
- Questions
 - probabilistic interpretation of the two long range dependence regimes for γ < d in "low" dimensions
 - interpretation of the "critical" dimension

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Idea of the proof of convergence

• Let $(\tilde{X}(t))_{t \in \mathbb{R}_+}$ be an $\mathcal{S}'(\mathbb{R}^d)$ -valued process. Define $\mathcal{S}'(\mathbb{R}^{d+1})$ -valued random variable \tilde{X} by:

$$\left\langle \tilde{X},\Phi\right\rangle =\int_{0}^{1}\left\langle X(t),\Phi(\cdot,t)
ight
angle \,dt,\qquad\Phi\in\mathcal{S}(\mathbb{R}^{d+1})$$

In $C([0, 1], \mathcal{S}'(\mathbb{R}^d))$: convergence $\langle \tilde{X}_T, \Phi \rangle \Rightarrow \langle \tilde{X}, \Phi \rangle$, in law $\forall \Phi \in \mathcal{S}(\mathbb{R}^{d+1})$ + tightness of $\langle X_T, \varphi \rangle_{T \ge 2}$ in $C([0, 1], \mathbb{R}), \forall \varphi \in \mathcal{S}(\mathbb{R}^d)$ \implies convergence of $X_T \Rightarrow X$ in law in w $C([0, 1], \mathcal{S}'(\mathbb{R}^d))$. (Remark: in general convergence of finite dimensional disributions of $X_T \Leftrightarrow$ convergence in law of \tilde{X}_T)

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- 2 Convergence of $\left< \tilde{X}, \Phi \right>$ is reduced to te case $\Phi \ge 0$.
- The limit process is totally skewed to the right ⇒ one can use Laplace transform. We show

$$Ee^{-\langle \tilde{X}_T, \Phi
angle} = \exp\left\{-\int_{\mathbb{R}^d} v_T(x, T) \nu_T(dx) + \frac{1}{F_T} E \int_0^T \langle N_s, \Phi
angle ds
ight\},$$

where ν_T – initial intensity measure, v_T satisfies a certain integral equation.

(Remark: a similar formula holds for finite dimensional distributions: $(\langle X_T(t_1), \varphi_1 \rangle, \dots, \langle X_T(t_n), \varphi_n \rangle)$ - by approximating $\sum_{i=1}^n \delta_{t_i} \varphi_i$ by $\Phi_m \in \mathcal{S}(\mathbb{R}^{d+1})$).

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$${\it Ee}^{-\left< ilde{X}_{{\it T}}, \Phi
ight>}
ightarrow {\it Ee}^{-\left< ilde{X}, \Phi
ight>}, \qquad \Phi \geq 0.$$

S Tightness (if the limit process is continuous)

• $\beta = 1$ one can study

$$E |\langle X_T(t_2) - X_T(t_1), \varphi \rangle|^k$$

• $\beta \leq 1$. We show: there exist $h, \varepsilon, r > 0$ such that

$$egin{aligned} & \mathcal{P}\left(|\langle X_{T}(t_{2})-X_{T}(t_{1}),arphi
angle|\geq\delta
ight) \ &\leq C\delta\int_{0}^{rac{1}{\delta}}\left(1-\operatorname{\mathsf{Re}}\mathcal{E}e^{-i heta\langle X_{T}(t_{2})-X_{T}(t_{1}),arphi
angle}
ight)d heta \ &\leq rac{C(arphi)}{\delta^{r}}\left(t_{2}^{h}-t_{1}^{h}
ight)^{1+arepsilon} \end{aligned}$$

for all $0 \leq t_1 \leq t_2$, $\delta > 0$.

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Let

$$\Psi_T(x,s) = \frac{1}{F_T} \int_{\frac{s}{T}}^1 \Phi(x,t) dt$$

 $(\eta_s)_{s\geq 0}$ – standard α -stable process

g – probability generating function of the branching mechanism N^x – empirical process for a tree starting from one particle at site x

Since N_0 is a Poisson random measure with intensity ν_T , we get

$$\begin{split} \mathsf{E} e^{-\langle \tilde{X}_{T}, \Phi \rangle} &= \exp\left\{ \int_{\mathbb{R}^{d}} \left(\mathsf{E} e^{-\int_{0}^{T} \langle N_{s}^{x}, \Psi_{T}(\cdot, s) \rangle ds} - 1 \right) \nu_{T}(dx) \right\} \\ & \qquad \times \exp\left\{ \int_{\mathbb{R}^{d}} \int_{0}^{T} \mathcal{T}_{s} \Psi_{T}(\cdot, s)(x) ds \nu_{T}(dx) \right\} \tag{L} \end{split}$$

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Let

$$w_T(x, r, t) := Ee^{-\int_0^T \langle N_s^x, \Psi_T(\cdot, r+s) \rangle ds}.$$

Conditioning with respect to the first branching

$$w_{T}(x, r, t) = e^{-Vt} \underbrace{Ee^{-\int_{0}^{t} \Psi_{T}(\eta_{s}^{x}, r+s)ds}}_{=: h_{T}(x, r, t)} + V \int_{0}^{t} e^{-V(t-s)} \underbrace{Ee^{-\int_{0}^{t-s} \Psi_{T}(\eta_{u}^{x}, r+u)du}}_{=: k_{T}(x, r, s, t-s)} g(w_{T}(\eta_{t-s}^{x}, r+t-s, s)) ds$$
(1)

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By the Feynman Kac Formula:

$$\begin{cases} \frac{\partial}{\partial t} h_T(x,r,t) = \left(\Delta_{\alpha} + \frac{\partial}{\partial r} - \Psi_T(x,r)\right) h_T(x,r,t) \\ h_T(x,r,0) = 1 \\ \begin{cases} \frac{\partial}{\partial \sigma} k_T(x,r,s,\sigma) = \left(\Delta_{\alpha} + \frac{\partial}{\partial r} - \Psi_T(x,r)\right) k_T(x,r,t) \\ k_T(x,r,s,0) = g(w_T(x,r,s)) \end{cases} \end{cases}$$

This and (1) gives

$$\begin{cases} \frac{\partial}{\partial t} w_{T}(x,r,t) = \left(\Delta_{\alpha} + \frac{\partial}{\partial r} - \Psi_{T}(x,r)\right) w_{T}(x,r,t) \\ + \frac{V}{1+\beta} \left(1 - w_{T}(x,r,t)\right)^{1+\beta} \\ w_{T}(x,r,0) = 1 \end{cases}$$

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Writing the equation in the mild form and substituting $v_T(x, t) = 1 - w_T(x, T - t, t)$, $(0 \le t \le T)$ we obtain

$$egin{aligned} & v_{T}(x,t) = \int_{0}^{t} \mathcal{T}_{t-s} \left(\Psi_{T}(\cdot,T-t)(1-v_{T}(\cdot,s))
ight. \ & \left. - rac{V}{1+eta} (v_{T}(\cdot,t))^{1+eta}
ight) (x) ds, \end{aligned}$$

From (L) we obtain

$$Ee^{-\langle \tilde{X}_{T}, \Phi \rangle}$$

$$= \exp\left\{\int_{\mathbb{R}^{d}} \left[-v_{T}(x, T) + \int_{0}^{T} \mathcal{T}_{T-s} \Psi_{T}(\cdot, T-s)(x) ds\right] \nu_{T}(dx)\right\}$$

$$= \exp\left\{\int_{\mathbb{R}^{d}} \int_{0}^{T} \mathcal{T}_{T-s} \left[\Psi_{T}(\cdot, T-s) v_{T}(\cdot, s) + (v_{T}(\cdot, s))^{1+\beta}\right] (x) ds \nu_{T}(dx)\right\}$$

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Usually (always for $\beta <$ 1, and if $\beta =$ 1 – in case of low dimensions)

$$\lim_{T \to \infty} E e^{-\langle \tilde{X}_{T}, \Phi \rangle} = \lim_{T \to \infty} \exp\left\{ \frac{V}{1+\beta} \int_{0}^{T} \mathcal{T}_{T-s} \left(\int_{0}^{s} \mathcal{T}_{s-u} \Psi_{T}(\cdot, T-u) \right)^{1+\beta}(x) \nu_{T}(dx) \right\}$$

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