

Finite Volume Scheme for a Corrosion Model

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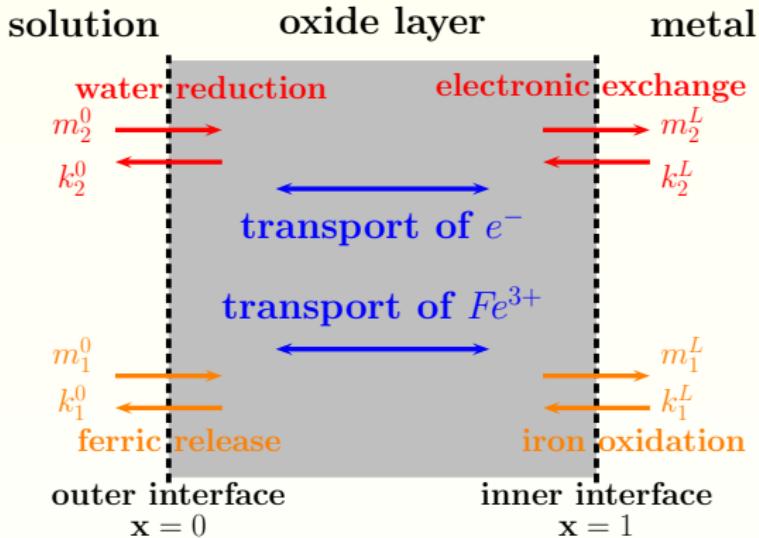
Collaboration with Christian Bataillon, CEA
Work supported by ANDRA

FVCA 5, 2008/06/11

Outline of the talk

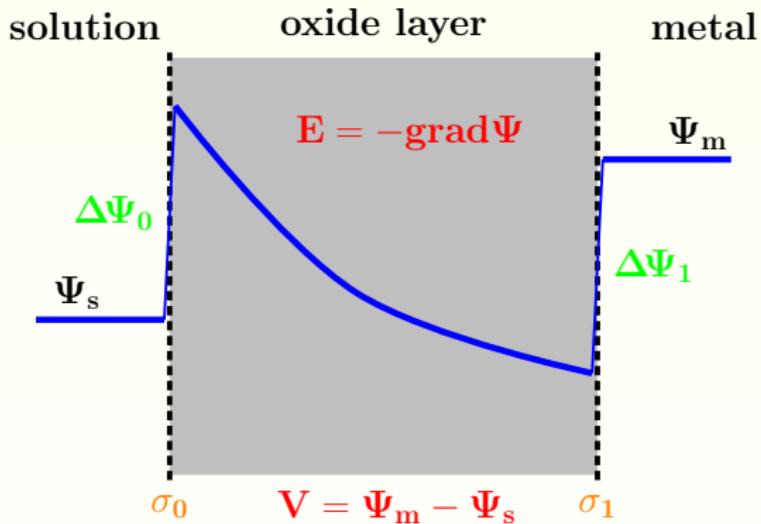
- ① Introduction to corrosion modelling
- ② Presentation of the scheme
- ③ Numerical analysis of the scheme
- ④ Numerical experiments

Introduction to corrosion modelling



Convection-diffusion equations on $\begin{cases} N : \text{density of electrons} \\ P : \text{density of cations} \end{cases}$

Introduction to corrosion modelling



Poisson equation on Ψ (electrostatic potential)

Introduction to corrosion modelling : the system

Equation and boundary conditions for P

$$\partial_t P + \partial_x J_P = 0, \quad J_P = -\partial_x P - 3P\partial_x \Psi$$

$$J_P(x=1) = m_1^L P e^{-3b_1^L(V-\Psi)} - (P_m - P) k_1^L e^{3a_1^L(V-\Psi)}$$

$$J_P(x=0) = -P k_1^0 e^{3a_1^0 \Psi}.$$

Equation and boundary conditions for N

$$\mathbf{D}_1/\mathbf{D}_2 \partial_t N + \partial_x J_N = 0, \quad J_N = -\partial_x N + N\partial_x \Psi$$

$$J_N(x=1) = m_2^L N e^{b_2^L(V-\Psi)} - (N_m - N) k_2^L e^{-a_2^L(V-\Psi)}$$

$$J_N(x=0) = -N k_2^0 e^{-a_2^0 \Psi}.$$

Equation and boundary conditions for Ψ

$$-\lambda^2 \partial_{xx}^2 \Psi = -N + 3P - 5$$

$$\Psi - \alpha_0 \partial_x \Psi = \Delta \Psi_0^{pzc} \quad \text{on } x=0,$$

$$\Psi + \alpha_1 \partial_x \Psi = V - \Delta \Psi_1^{pzc} \quad \text{on } x=1.$$

Introduction to corrosion modelling : the system

Equation and boundary conditions for P

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$$\Psi + \alpha_1 \partial_x \Psi = V - \Delta \Psi_1^{pzc} \quad \text{on } x=1.$$

Presentation of the scheme

Scheme for P : $\mathcal{F}_{i+\frac{1}{2}} - \mathcal{F}_{i-\frac{1}{2}} = 0, 1 \leq i \leq I$

$$\mathcal{F}_{\frac{1}{2}} = m_1^0 e^{-3b_1^0 \Psi_0} P_m - \left(k_1^0 e^{3a_1^0 \Psi_0} + m_1^0 e^{-3b_1^0 \Psi_0} \right) P_0,$$

$$\mathcal{F}_{I+\frac{1}{2}} = \left(k_1^L e^{3a_1^L (V - \Psi_{I+1})} + m_1^L e^{-3b_1^L (V - \Psi_{I+1})} \right) P_{I+1} - k_1^L e^{3a_1^L (V - \Psi_{I+1})} P_m.$$

Scheme for N : $\mathcal{G}_{i+\frac{1}{2}} - \mathcal{G}_{i-\frac{1}{2}} = 0, 1 \leq i \leq I$

$$\mathcal{G}_{\frac{1}{2}} = m_2^0 e^{b_2^0 \Psi_0} N_m - \left(m_2^0 e^{b_2^0 \Psi_0} + k_2^0 e^{-a_2^0 \Psi_0} \right) N_0,$$

$$\mathcal{G}_{I+\frac{1}{2}} = \left(m_2^L e^{b_2^L (V - \Psi_{I+1})} + k_2^L e^{-a_2^L (V - \Psi_{I+1})} \right) N_{I+1} - k_2^L e^{-a_2^L (V - \Psi_{I+1})} N_m.$$

Scheme for Ψ : $-\lambda^2 (d\Psi_{i+\frac{1}{2}} - d\Psi_{i-\frac{1}{2}}) = h_i (3P_i - N_i - 5), 1 \leq i \leq I$

$$\begin{aligned} \Psi_0 - \alpha_0 d\Psi_{\frac{1}{2}} &= U_0, & d\Psi_{i+\frac{1}{2}} &= \frac{\Psi_{i+1} - \Psi_i}{h_{i+\frac{1}{2}}} \\ \Psi_{I+1} + \alpha_1 d\Psi_{I+\frac{1}{2}} &= U_1. \end{aligned}$$

Presentation of the scheme : the fluxes

$$J_P = -\partial_x P - 3P\partial_x \Psi, \quad J_N = -\partial_x N + N\partial_x \Psi$$

$$\mathcal{F}_{i+\frac{1}{2}} = \frac{B(3(\Psi_{i+1} - \Psi_i))P_i - B(-3(\Psi_{i+1} - \Psi_i))P_{i+1}}{h_{i+\frac{1}{2}}} \quad 0 \leq i \leq I$$

$$\mathcal{G}_{i+\frac{1}{2}} = \frac{B(-(\Psi_{i+1} - \Psi_i))N_i - B(\Psi_{i+1} - \Psi_i)N_{i+1}}{h_{i+\frac{1}{2}}} \quad 0 \leq i \leq I$$

Choices for B

- Classical upwind scheme

$$B(x) = 1 + (-x)^+$$

- Scharfetter-Gummel scheme

$$B(x) = \frac{x}{e^x - 1} \text{ for } x \neq 0, \quad B(0) = 1.$$

Presentation of the scheme : the Slotboom variables

Change the variables

$$u_i = e^{3\Psi_i} P_i, \quad v_i = e^{-\Psi_i} N_i$$



$$\mathcal{F}_{i+\frac{1}{2}} = \frac{1}{\mathbf{f}(\Psi_i, \Psi_{i+1})} \frac{u_i - u_{i+1}}{h_{i+\frac{1}{2}}}, \quad \mathcal{G}_{i+\frac{1}{2}} = \frac{1}{\mathbf{g}(\Psi_i, \Psi_{i+1})} \frac{v_i - v_{i+1}}{h_{i+\frac{1}{2}}}$$

$\implies \mathbf{u}$ and \mathbf{v} expressed as a function of Ψ

Nonlinear system of equations on Ψ

$$-\lambda^2 \left(\frac{\Psi_{i+1} - \Psi_i}{h_{i+\frac{1}{2}}} - \frac{\Psi_i - \Psi_{i-1}}{h_{i-\frac{1}{2}}} \right) = h_i (3u_i e^{-3\Psi_i} - v_i e^{\Psi_i} - 5),$$

$$\Psi_0 - \alpha_0 d\Psi_{\frac{1}{2}} = U_0,$$

$$\Psi_{I+1} + \alpha_1 d\Psi_{I+\frac{1}{2}} = U_1.$$

Numerical analysis

Theorem 1

Under hypothesis : $a_i^0 + b_i^0 = a_i^L + b_i^L = 1$, $i = 1, 2$, we have :

- **Existence of a solution** to the scheme :

$$\mathbf{u} = (u_i)_{0 \leq i \leq I+1}, \mathbf{v} = (v_i)_{0 \leq i \leq I+1}, \Psi = (\Psi_i)_{0 \leq i \leq I+1},$$

- with **L^∞ and discrete H^1 -estimates** :

$$0 \leq u_i \leq \frac{k_1^L}{m_1^L} e^{3V} P_m, \quad 0 \leq v_i \leq \frac{k_2^L}{m_2^L} e^{-V} N_m$$

$$|\Psi_i| \leq M \text{ and } \sum_{i=0}^I \frac{(\Psi_{i+1} - \Psi_i)^2}{h_{i+\frac{1}{2}}} + \Psi_0^2 + \Psi_{I+1}^2 \leq M$$

$$\sum_{i=0}^I \frac{(u_i - u_{i+1})^2}{h_{i+\frac{1}{2}}} \leq M, \quad \sum_{i=0}^I \frac{(v_i - v_{i+1})^2}{h_{i+\frac{1}{2}}} \leq M.$$

Numerical analysis

Theorem 2

- $a_i^0 + b_i^0 = a_i^L + b_i^L = 1, i = 1, 2$
- $(\mathbf{u}^h, \mathbf{v}^h, \Psi^h)_{h>0}$ sequence of approximate solutions

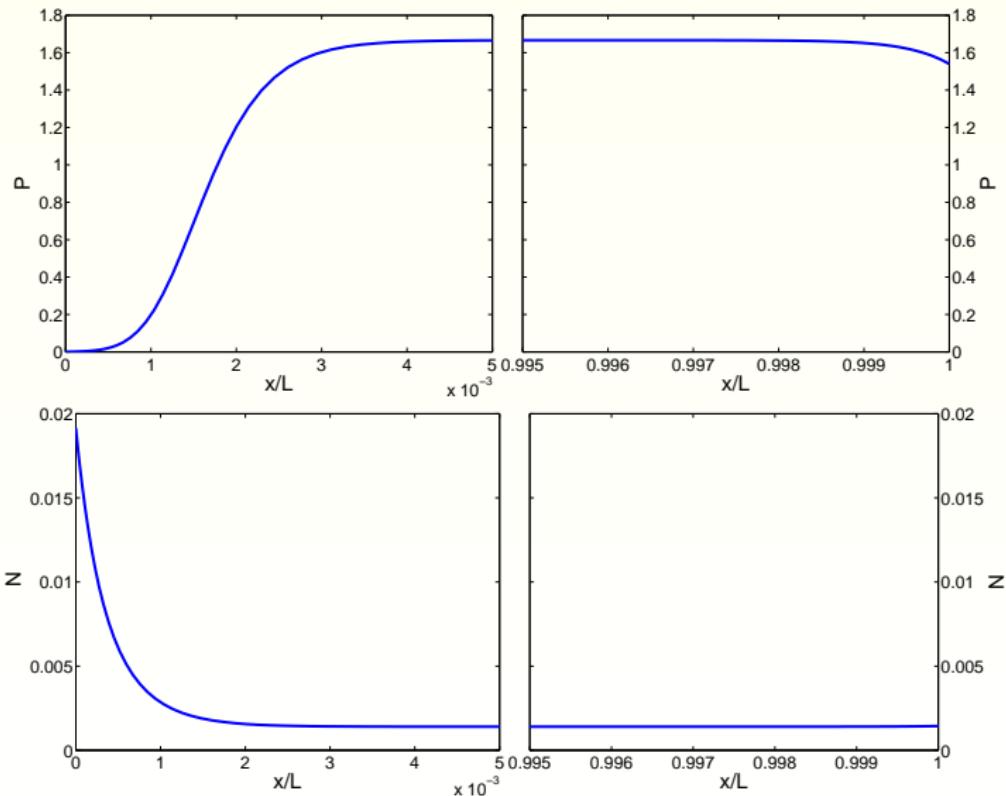
Then, up to a subsequence,

$$\mathbf{u}^h \rightarrow u, \quad \mathbf{v}^h \rightarrow v, \quad \Psi^h \rightarrow \Psi \text{ as } h \rightarrow 0 \text{ in } C([0, 1]).$$

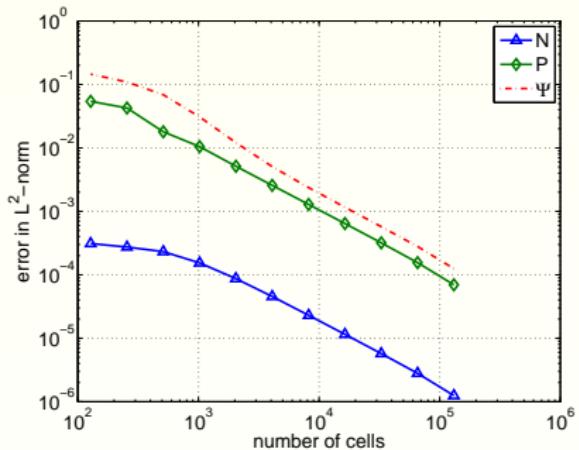
and

$(N = e^\Psi v, P = e^{-3\Psi} u, \Psi)$ is a weak solution
to the corrosion model.

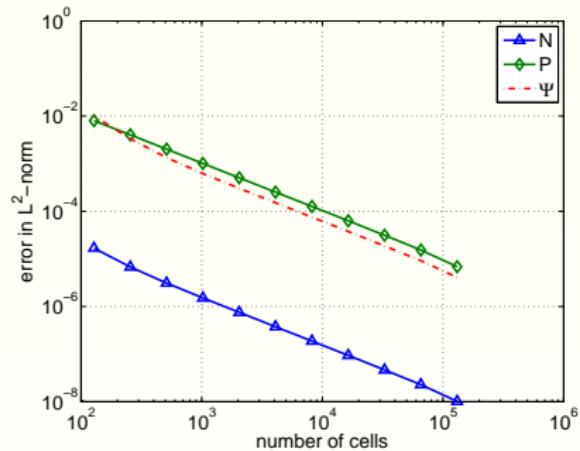
Numerical experiments : profiles



Numerical experiments : convergence



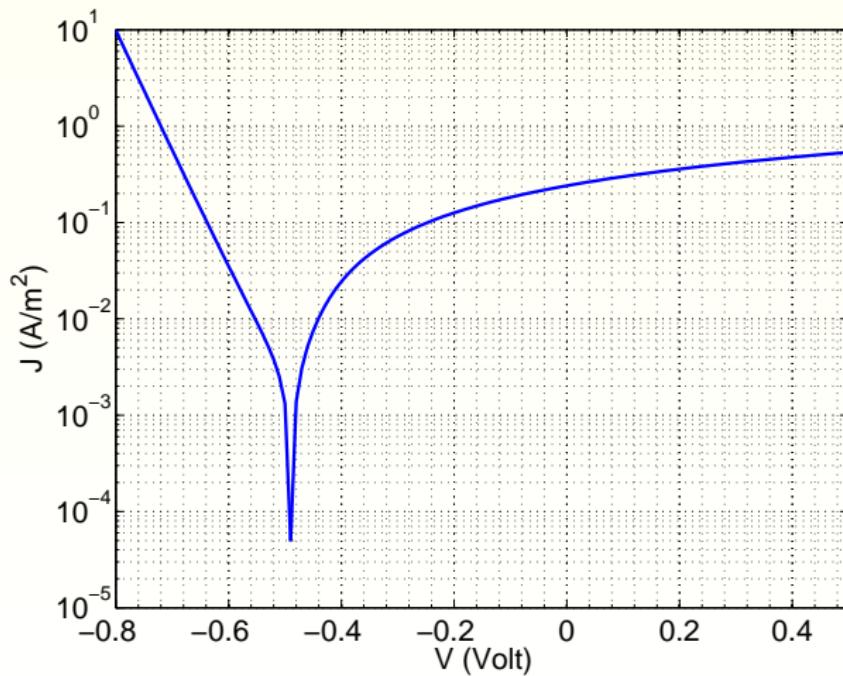
Uniform mesh



Tchebychev mesh

Numerical experiments

Current-voltage curve



Prospects

- More numerical simulations and comparison to the experiments :
 ⇒ Help for the choice of all parameters
- Addition of a third charge carrier : the oxygen vacancies
- Evolution in time of the oxide layer :
 - Scheme for the non stationary system of equations
 - Free boundary problem
- Computation of the corrosion speed