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**Renormalization of the Numerical Diffusion
for an Upwind Finite Volume Method.
Application to the Simulation of
Kelvin-Helmholtz Instability**

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1 Introduction

One of the most challenging problem in CFD is the building of a numerical scheme which behaves well for both compressible and (almost) incompressible flows.

Compressible flows : FINITE VOLUMES ($M < \infty$)

Incompressible flows : FINITE ELEMENTS ($M = 0$)

$$M \rightarrow 0$$

2 Kelvin-Helmholtz Instability

2D Euler's equations :

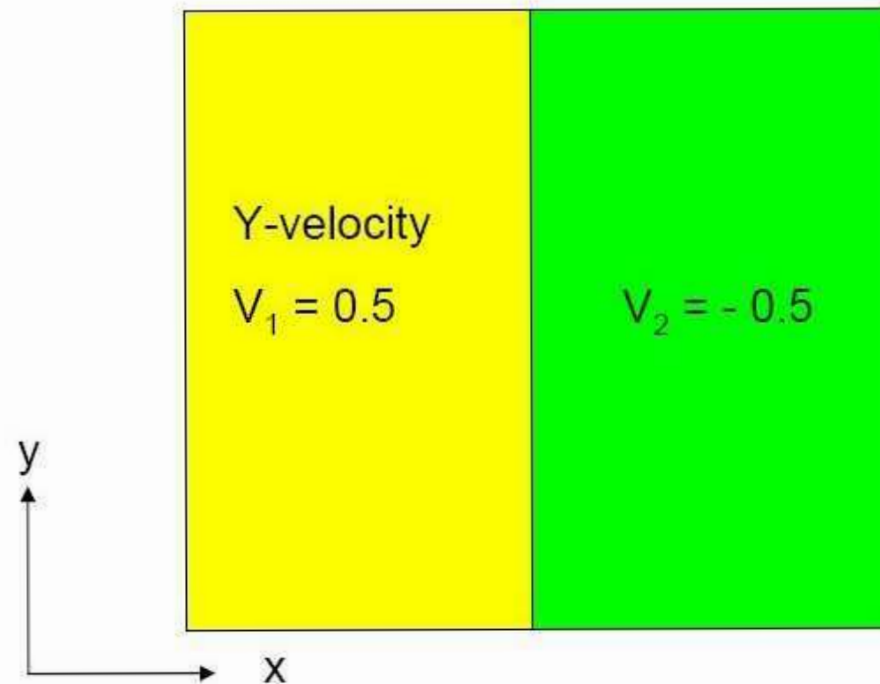
$$\frac{\partial \rho}{\partial t} + \operatorname{div}(\rho u) = 0, \quad (1)$$

$$\frac{\partial(\rho u)}{\partial t} + \operatorname{div}(\rho u \otimes u + p1) = 0, \quad (2)$$

$$\frac{\partial(\rho E)}{\partial t} + \operatorname{div}(\rho E u + p u) = 0, \quad (3)$$

where ρ denotes the density, e the internal energy, $p = p(\rho, e)$ the pressure, $u = (u_x, u_y)$ the velocity, $E = e + \frac{1}{2}|u|^2$ the total specific energy.

KELVIN-HELMHOLTZ TEST PROBLEM – used for dissipation studies



Size of box 1 x 1

Computational cells

64 times 64

$$\Delta V = |V_1 - V_2| = 1$$

Density = 1, Perfect gas equation of state ($\gamma = 5/3$)

Pressure = 15, Sound speed = 5 (default case)

Mach no. $\Delta V / c = 0.2$ (default case)

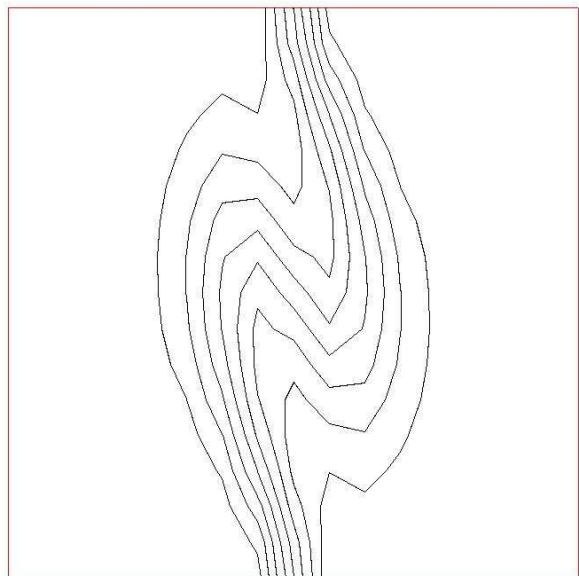
3 Finite Volume Approximation

$$|K| \frac{dv_K}{dt} + \sum_{L \in \mathcal{N}(K)} |K \cap L| \Phi(v_K, v_L; K, L) = 0,$$

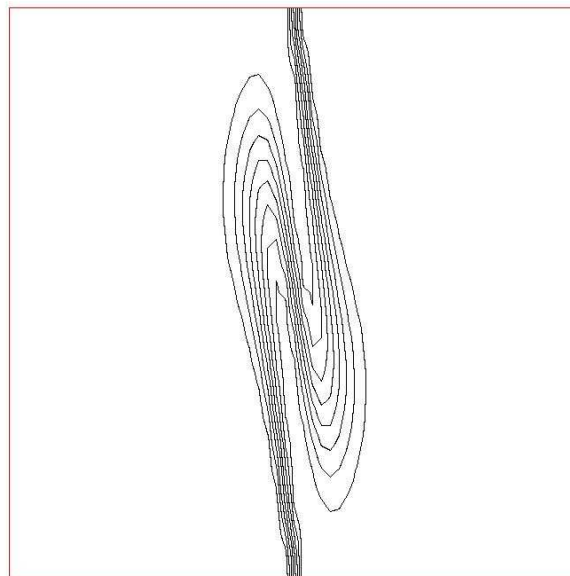
$$\Phi(v, w; K, L) = \frac{F(v) + F(w)}{2} \cdot \nu_{K,L} - \mathcal{U}(v, w; K, L) \frac{F(w) - F(v)}{2} \cdot \nu_{K,L}.$$

$$\mathcal{U}(v, w; K, L) = \text{sgn}(A_{\nu_{K,L}}(\mu_{K,L}))$$

$$p_{j+\frac{1}{2}} = \frac{p_j + p_{j+1}}{2} - \frac{\text{sgn}(u)}{2} (M \Delta p + \frac{\rho u}{\textcolor{red}{M}} \Delta u) + o(\Delta^2 + M^2)$$



(a) $M=0.2$.



(b) $M=0.02$.

Figure 1: Interface at $t = 1$.

4 Renormalization

$$\Phi(v, w; K, L) = \frac{F(v) + F(w)}{2} \cdot \nu_{K,L} - \mathcal{U}(v, w; K, L) \frac{F(w) - F(v)}{2} \cdot \nu_{K,L} .$$

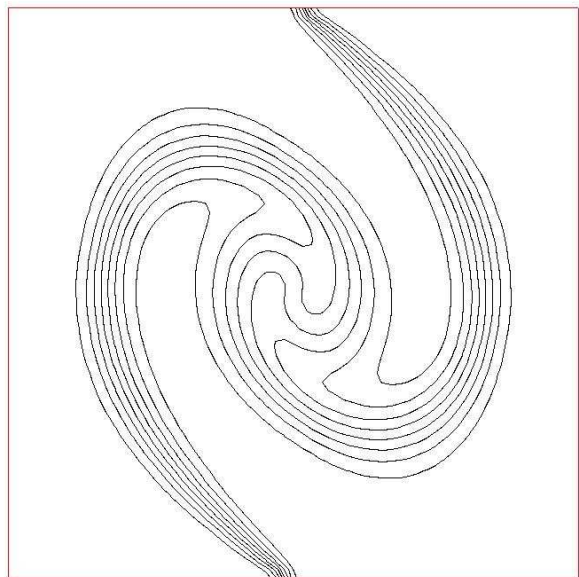
$$\mathcal{U}(v, w; K, L; \beta) = \text{sgn} \left(A_{\nu_{K,L}}(\mu_{K,L}) P(\mu_{K,L}, \beta) \right) ,$$

where $P(\mu_{K,L}, \beta)$ is the preconditioning matrix proposed E. Turkel *J. Comp. Phys.* in 1987 with β a parameter of the order of the Mach number. [Here: time consistent scheme.](#)

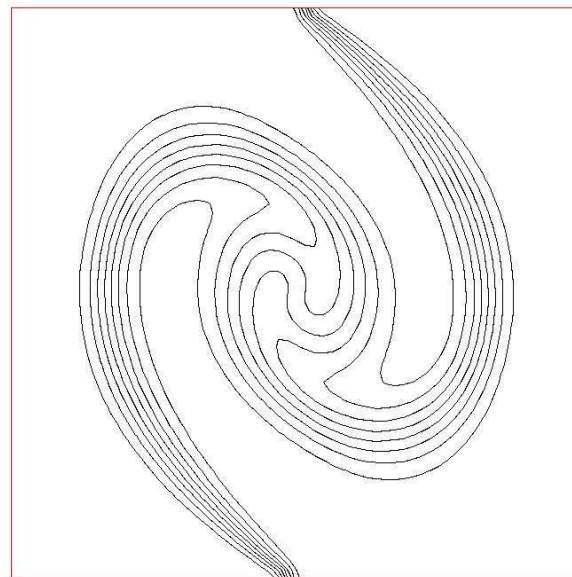
Proposition 1 *For the renormalized FVCF scheme ($\beta = M$), we have*

$$p_{j+\frac{1}{2}} = \frac{p_j + p_{j+1}}{2} - \frac{\text{sgn}(u)}{2\sqrt{5}} (\Delta p + 3\rho u \Delta u) + o(\Delta^2 + M^2) .$$

$$\left(p_{j+\frac{1}{2}} = \frac{p_j + p_{j+1}}{2} - \frac{\text{sgn}(u)}{2} (M \Delta p + \frac{\rho u}{M} \Delta u) + o(\Delta^2 + M^2) \right)$$

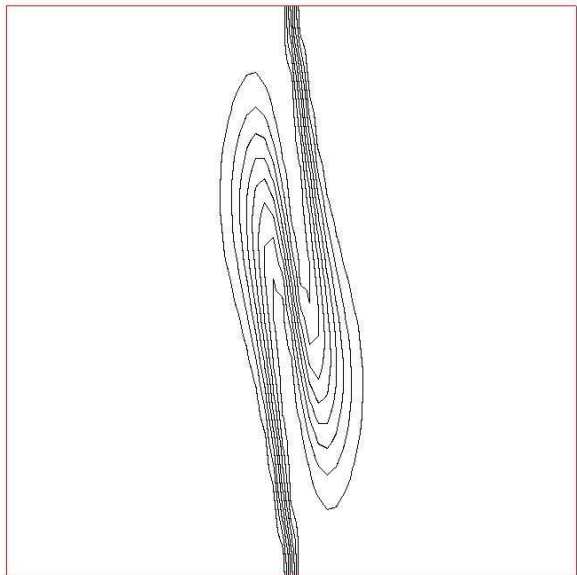


(a) $M=0.2, \beta=M$.

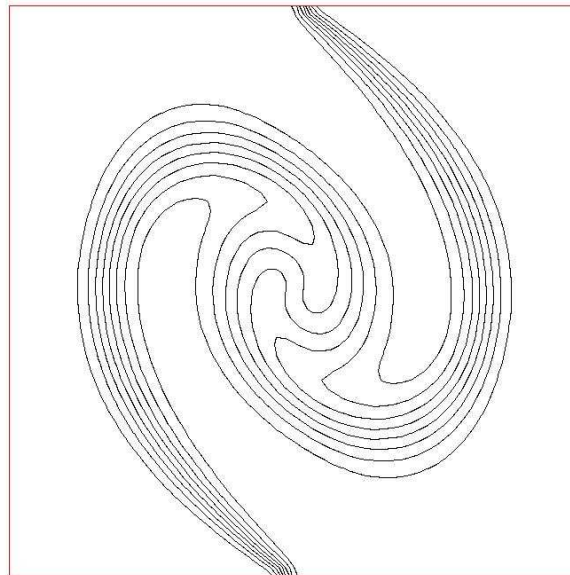


(b) $M=0.02, \beta=M$.

Figure 2: Interface at $t = 1$ with renormalization.



(a) $M=0.02, \beta=1$.



(b) $M=0.02, \beta=M$.

Figure 3: Interface at $t = 1$ without and with renormalization.

5 Explicit versus Implicit schemes

Euler explicit scheme reads:

$$\frac{|K|}{\Delta t^n} (v_K^{n+1} - v_K^n) + \sum_{L \in \mathcal{N}(K)} |K \cap L| \Phi(v_K^n, v_L^n; K, L; \beta) = 0. \quad (4)$$

Bypass the C.F.L. condition: implicit Euler discretization:

$$\frac{|K|}{\Delta t^n} (v_K^{n+1} - v_K^n) + \sum_{L \in \mathcal{N}(K)} |K \cap L| \Phi(v_K^{n+1}, v_L^{n+1}; K, L; \beta) = 0. \quad (5)$$

Stability

Tool for the study of stability: Von Neuman analysis (amplification matrix).

	<i>without renormalization</i>	<i>with renormalization</i>
Explicit	$u \frac{\Delta t}{\Delta x} = \mathcal{O}(M)$	$u \frac{\Delta t}{\Delta x} = \mathcal{O}(M^2)$
Implicit	unconditionally	unconditionally

Table 1: Stability (time discretization)

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