

Aim...

Overview on the references

Problem to be solved and discretization parameters

Finite element discretization

Formulation of a fully implicit discretization scheme

Some known error estimates

New error estimate

Proof of the new $W^{1,\infty}(L^2)$ -error estimate

Conclusion and perspectives

A new error estimate for a fully finite element discretization scheme for parabolic equations using Crank–Nicolson method

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Aim

We consider a conforming finite element method in which the discretization in time is performed using the Crank-Nicolson method for the **non stationary heat equation** (as a model for parabolic equations). We provide an error estimate in $W^{1,\infty}(L^2)$ -norm. This error estimate seems not to be present in the existing literature.

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Overview : References based on...

- Bradji, A.: An analysis of a second-order time accurate scheme for a finite volume method for parabolic equations on general nonconforming multidimensional spatial meshes. Appl. Math. Comput. **219**/11 (2013), 6354–6371.
- Burman, E.: Crank–Nicolson finite element methods using symmetric stabilization with application to optimal control problems subject to transient advection–diffusion equations. Commun. Math. Sci. **9**/1 (2011), 319–329.

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Overview : References based on...(Suite)

- Quarteroni, A. and Valli, A.: Numerical Approximation of Partial Differential Equations. Springer Series in Computational Mathematics 23. Berlin: Springer. (2008)
- Chatzipantelidis, P., Lazarov, R.D., and Thomée, V.: Some error estimates for the lumped mass finite element method for a parabolic problem. Math. Comput. **81**/277 (2012), 1–20.
- Raviart, P. A. and Thomas, J. M.: Introduction à l'Analyse Numérique des Equations aux Dérivées Partielles. Masson, Paris. (1983).

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Overview : References based on...(Suite)

- Thomée, V.: Galerkin Finite Element Methods for Parabolic Problems. Springer-Verlag, Second Edition, Berlin (2006).

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Plan of the presentation

- 1 Problem to be discretized
- 2 Heat equation and a fully implicit finite element discretization using Crank-Nicolson method as discretization in time.
- 3 Some known error estimates e
- 4 Statement of an error estimate (main result)
- 5 Some applications
- 6 Conclusion and some perspectives

Non stationary Heat equation

$$u_t(x, t) - \Delta u(x, t) = f(x, t), \quad (x, t) \in \Omega \times (0, T), \quad (1)$$

where, Ω is an open bounded polyhedral subset in \mathbb{R}^d , with $d \in \mathbb{N}^*$, $T > 0$, and f is a given function.

An initial condition is given by:

$$u(x, 0) = u^0(x), \quad x \in \Omega. \quad (2)$$

A Dirichlet boundary condition is defined by

$$u(x, t) = 0, \quad (x, t) \in \partial\Omega \times (0, T), \quad (3)$$

where, we denote by $\partial\Omega = \overline{\Omega} \setminus \Omega$ the boundary of Ω .

About Heat equation?

- 1 (some physics): Heat equation $u_t - \Delta u$ is typically used in different applications, such as *fluid mechanics*, *heat and mass transfer*,...
- 2 (existence and uniqueness): existence and uniqueness of a **weak** solution of heat equation, with (2) (*initial condition*) and (3) (*Dirichlet boundary condition*) can be formulated using **Bochner spaces**; see for instance **Evans book of partial differential equation**

Finite element discretization

- 1 Let $\{\mathcal{T}_h; h > 0\}$ be a family of shape regular and quasi-uniform triangulations of the domain Ω . The elements of \mathcal{T}_h will be denoted by K . For each triangulation \mathcal{T}_h , the subscript h refers to the level of refinement of the triangulation, which is defined by $h = \max_{K \in \mathcal{T}_h} h_K$, where h_K denotes the diameter of the element K .

Let \mathcal{V}_0^h be the finite element space of continuous, piecewise polynomial functions of degree $k \geq 1$, and vanish on the boundary $\partial\Omega$

- 2 Uniform mesh on $(0, T)$ with constant step $\tau = T/(N + 1)$.

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Discretization of initial conditions

The discretization of the initial condition (2), $u(x, 0) = u^0(x)$ is performed using the **orthogonal projection** : Find $u_h^0 \in \mathcal{V}_0^h$ such that

$$\mathbf{a}(u_h^0, v) = - \left(\Delta u^0, v \right)_{\mathbb{L}^2(\Omega)} = \mathbf{a}(u^0, v), \quad \forall v \in \mathcal{V}_0^h. \quad (4)$$

where $\mathbf{a}(\cdot, \cdot)$ denotes the bilinear form defined for all $(u, v) \in H^1(\Omega) \times H^1(\Omega)$ by

$$\mathbf{a}(u, v) = \int_{\Omega} \nabla u(x) \cdot \nabla v(x) dx.$$



Discretization of the heat equation

The discretization of the heat equation (1), $u_t - \Delta u = f$ is performed as : for any $n \in \llbracket 0, M \rrbracket$, find $u_h^n \in \mathcal{V}_0^h$ such that, for all $v \in \mathcal{V}_0^h$

$$\left(\partial^1 u_h^{n+1}, v \right)_{\mathbb{L}^2(\Omega)} + \mathbf{a}(u_h^{n+\frac{1}{2}}, v) = \left(\frac{1}{\tau} \int_{t_n}^{t_{n+1}} f(t) dt, v \right)_{\mathbb{L}^2(\Omega)}, \quad (5)$$

where

$$\partial^1 v^{n+1} = \frac{v^{n+1} - v^n}{k} \quad \text{and} \quad v^{n-\frac{1}{2}} = \frac{v^n + v^{n-1}}{2}. \quad (6)$$

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Some known error estimates: $\mathbb{L}^\infty(\mathbb{L}^2)$ -error estimate

Throughout this talk, the notation C stands for a positive constant which is independent of the parameters of the discretization.

$\mathbb{L}^\infty(\mathbb{L}^2)$ -error estimate. Under some regularity assumption on the data and on the exact solution, the following $\mathbb{L}^\infty(\mathbb{L}^2)$ -error estimate holds, see the books of Quarteroni and Valli (2008) and Raviart and Thomas (1983), for all $n \in \llbracket 0, M + 1 \rrbracket$:

$$\|u_h^n - u(t_n)\|_{\mathbb{L}^2(\Omega)} \leq C(h^{k+1} + \tau^2). \quad (7)$$

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Some known error estimates: $\mathbb{L}^2(H^1)$ -error estimate

$\mathbb{L}^2(H^1)$ -error estimate. Under some regularity assumption on the data and on the exact solution, the following $\mathbb{L}^2(H^1)$ -error estimate holds, see the article of Burman (2011), for all $n \in \llbracket 0, M + 1 \rrbracket$:

$$\left(\sum_{n=0}^M \tau \| e_h^{n+\frac{1}{2}} \|_{H^1(\Omega)}^2 \right)^{\frac{1}{2}} \leq C(h^k + \tau^2). \quad (8)$$

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New error estimate: $W^{1,\infty}(\mathbb{L}^2)$ -error estimate

$W^{1,\infty}(\mathbb{L}^2)$ -error estimate. Under some regularity assumption on the data and on the exact solution, the following $W^{1,\infty}(\mathbb{L}^2)$ -error estimate holds, for all $n \in \llbracket 0, M \rrbracket$:

$$\|\partial^1 (u_h^{n+1} - u(t_{n+1}))\|_{\mathbb{L}^2(\Omega)} \leq C(h^{k+1} + \tau^2), \quad (9)$$

where ∂^1 is the discrete time derivative

$$\partial^1 v^{n+1} = \frac{v^{n+1} - v^n}{k}. \quad (10)$$

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An application of the new error estimate: approximation of the first time derivative

$\partial^1 u_h^{n+1}$ approximates the time derivative of u at $t_{\frac{n+1}{2}}$, i.e. $u_t(t_{\frac{n+1}{2}})$, by order $h^{k+1} + \tau^2$ in $\mathbb{L}^\infty(\mathbb{L}^2)$ -norm, where $t_{\frac{n+1}{2}} = (t_{n+1} + t_n) / 2$.

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An a priori estimate

Lemma

Assume that there exists $(\eta_h^n)_{n=0}^{M+1} \in (\mathcal{V}_0^h)^{M+2}$ such that $\eta_h^0 = 0$ and for all $n \in \llbracket 0, M \rrbracket$

$$\left(\partial^1 \eta_h^{n+1}, \mathbf{v} \right)_{\mathbb{L}^2(\Omega)} + \mathbf{a}(\eta_h^{n+\frac{1}{2}}, \mathbf{v}) = (\gamma^n, \mathbf{v})_{\mathbb{L}^2(\Omega)}, \quad \forall \mathbf{v} \in \mathcal{V}_0^h. \quad (11)$$

Then the following estimate holds:

$$\| \partial^1 \eta_h^n \|_{\mathbb{L}^2(\Omega)} \leq C(\gamma + \bar{\gamma}), \quad \forall n \in \llbracket 1, M+1 \rrbracket, \quad (12)$$

where

$$\gamma = \max_{n=0}^M \| \gamma^n \|_{\mathbb{L}^2(\Omega)} \quad \text{and} \quad \bar{\gamma} = \max_{n=1}^M \| \partial^1 \gamma^n \|_{\mathbb{L}^2(\Omega)}. \quad (13)$$

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Idea on the proof of Lemma 1

First step: Estimate on $\partial^1 \eta_h^{j+1}$, $j \in \llbracket 1, M \rrbracket$. Acting the discrete operator ∂^1 on (11) to get, for all $n \in \llbracket 1, M \rrbracket$

$$\left(\partial^2 \eta_h^{n+1}, v \right)_{\mathbb{L}^2(\Omega)} + \mathbf{a}(\partial^1 \eta_h^{n+\frac{1}{2}}, v) = \left(\partial^1 \gamma^n, v \right)_{\mathbb{L}^2(\Omega)}. \quad (14)$$

Taking $v = \partial^1 \eta_h^{n+1} + \partial^1 \eta_h^n$ in (14) to get

$$\begin{aligned} \|\partial^1 \eta_h^{n+1}\|_{\mathbb{L}^2(\Omega)}^2 &- \|\partial^1 \eta_h^n\|_{\mathbb{L}^2(\Omega)}^2 + \frac{\tau}{2} \left| \partial^1 \left(\eta_h^{n+1} + \eta_h^n \right) \right|_{1,\Omega}^2 \\ &= \tau \left(\partial^1 \gamma^n, \partial^1 \left(\eta_h^{n+1} + \eta_h^n \right) \right)_{\mathbb{L}^2(\Omega)}. \end{aligned}$$

Summing the previous equality over $n \in \llbracket 1, j \rrbracket$, where $j \in \llbracket 1, M \rrbracket$ and using some technical steps with a discrete Poincaré inequality yields

$$\|\partial^1 \eta_h^{j+1}\|_{\mathbb{L}^2(\Omega)}^2 \leq \|\partial^1 \eta_h^1\|_{\mathbb{L}^2(\Omega)}^2 + 4T (C_p)^2 (\bar{\gamma})^2. \quad (15)$$

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Idea on the proof of Lemma 1 (Suite)

Second step: Estimate on $\partial^1 \eta_h^j$. Taking $n = 0$ in (11) to get (note that $\eta_h^0 = 0$)

$$\left(\partial^1 \eta_h^1, v \right)_{\mathbb{L}^2(\Omega)} + \frac{1}{2} \mathbf{a}(\eta_h^1, v) = \left(\gamma^0, v \right)_{\mathbb{L}^2(\Omega)}.$$

Taking $v = \partial^1 \eta_h^1$ in the previous equality and using some technical steps leads to

$$\| \partial^1 \eta_h^1 \|_{\mathbb{L}^2(\Omega)} \leq \gamma. \quad (16)$$

This with (15) implies the desired estimate of Lemma 1. ■

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Idea on the proof of the new $W^{1,\infty}(\mathbb{L}^2)$ -error estimate

The proof of the new $W^{1,\infty}(\mathbb{L}^2)$ -error estimate is based essentially on the comparison with the following finite element scheme:

For each $n \in \llbracket 0, M + 1 \rrbracket$, we compute $\bar{u}_h^n \in \mathcal{V}_0^h$ such that

$$\mathbf{a}(\bar{u}_h^n, v) = -(\Delta u(t_n), v)_{\mathbb{L}^2(\Omega)} = \mathbf{a}(u(t_n), v), \quad \forall v \in \mathcal{V}_0^h. \quad (17)$$

The following convergence result can be shown using the classical error estimates in finite element methods

$$|\bar{u}_h^n - u(t_n)|_{1,\Omega} \leq Ch^k, \quad (18)$$

and, for all $j \in \{1, 2\}$

$$\|\partial^j \bar{u}_h^n - \partial^j u(t_n)\|_{\mathbb{L}^2(\Omega)} \leq Ch^{k+1}. \quad (19)$$

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Idea on the proof of the new $W^{1,\infty}(\mathbb{L}^2)$ -error estimate (Suite)

We write the error between the exact solution u and its finite element approximate solution as

$$u(t_n) - u_h^n = (u(t_n) - \bar{u}_h^n) + (\bar{u}_h^n - u_h^n). \quad (20)$$

The error $(u(t_n) - \bar{u}_h^n)$ is already estimated in (17)–(18). It remains now to estimate $\eta_h^n =: \bar{u}_h^n - u_h^n$. Using the schemes satisfied by \bar{u}_h^n and u_h^n , we get

$$\left(\partial^1 \eta_h^{n+1}, \mathbf{v} \right)_{\mathbb{L}^2(\Omega)} + \mathbf{a}(\eta_h^{n+\frac{1}{2}}, \mathbf{v}) = \left(\mathbb{K}^{n,1} - \mathbb{K}^{n,2}, \mathbf{v} \right)_{\mathbb{L}^2(\Omega)}, \quad (21)$$

where

$$\mathbb{K}^{n,1} = -\partial^1 \left(u(t_{n+1}) - \bar{u}_h^{n+1} \right) \quad \text{and} \quad \mathbb{K}^{n,2} = -\frac{1}{\tau} \int_{t_n}^{t_{n+1}} \Delta u(t) dt + \frac{\Delta u(t_{n+1}) + \Delta u(t_n)}{2}.$$

Idea on the proof of the new $W^{1,\infty}(\mathbb{L}^2)$ -error estimate (Suite)

Recall that

$$\left(\partial^1 \eta_h^{n+1}, \mathbf{v} \right)_{\mathbb{L}^2(\Omega)} + \mathbf{a}(\eta_h^{n+\frac{1}{2}}, \mathbf{v}) = \left(\mathbb{K}^{n,1} - \mathbb{K}^{n,2}, \mathbf{v} \right)_{\mathbb{L}^2(\Omega)}. \quad (22)$$

Using Lemma 1 to obtain

$$\| \partial^1 \eta_h^n \|_{\mathbb{L}^2(\Omega)} \leq C(\gamma + \bar{\gamma}), \quad \forall n \in \llbracket 1, M+1 \rrbracket, \quad (23)$$

where

$$\gamma = \max_{n=0}^M \| \mathbb{K}^{n,1} - \mathbb{K}^{n,2} \|_{\mathbb{L}^2(\Omega)} \quad \text{and} \quad \bar{\gamma} = \max_{n=1}^M \| \partial^1 (\mathbb{K}^{n,1} - \mathbb{K}^{n,2}) \|_{\mathbb{L}^2(\Omega)}. \quad (24)$$

Error estimates (17)–(18) imply that, for $j \in \{0, 1\}$

$$\| \partial^j \mathbb{K}^{n,1} \|_{\mathbb{L}^2(\Omega)} \leq Ch^{k+1}. \quad (25)$$

Idea on the proof of the new $W^{1,\infty}(\mathbb{L}^2)$ -error estimate (Suite)

Recall that

$$\mathbb{K}^{n,2} = -\frac{1}{\tau} \int_{t_n}^{t_{n+1}} \Delta u(t) dt + \frac{\Delta u(t_{n+1}) + \Delta u(t_n)}{2}.$$

We use

$$\mathbb{K}^{n,2} = -\frac{1}{\tau} \int_0^\tau \left(\frac{(t - \frac{\tau}{2})^2}{2} - \frac{\tau^2}{8} \right) \Delta u_{tt}(t + t_n) dt. \quad (26)$$

We can easily check that $\frac{(t - \frac{\tau}{2})^2}{2} - \frac{\tau^2}{8}$ is non-positive for $t \in [0, \tau]$ and by some elementary calculations, we get

$$\int_0^\tau \left(\frac{(t - \frac{\tau}{2})^2}{2} - \frac{\tau^2}{8} \right) dt = -\frac{\tau^3}{12}.$$

This with allows us to get, for $j \in \{0, 1\}$

$$\|\partial^j \mathbb{K}^{n,2}\|_{\mathbb{L}^2(\Omega)} \leq C\tau^2. \quad (27)$$

Idea on the proof of the new $W^{1,\infty}(L^2)$ -error estimate (Suite)

Recall that

$$\gamma = \max_{n=0}^M \|\mathbb{K}^{n,1} - \mathbb{K}^{n,2}\|_{\mathbb{L}^2(\Omega)} \quad \text{and} \quad \bar{\gamma} = \max_{n=1}^M \|\partial^1(\mathbb{K}^{n,1} - \mathbb{K}^{n,2})\|_{\mathbb{L}^2(\Omega)}, \quad (28)$$

and the following estimates have been obtained

$$\|\partial^j \mathbb{K}^{n,1}\|_{\mathbb{L}^2(\Omega)} \leq Ch^{k+1} \quad \text{and} \quad \|\partial^j \mathbb{K}^{n,2}\|_{\mathbb{L}^2(\Omega)} \leq C\tau^2. \quad (29)$$

Consequently, with (28)

$$\|\gamma\|_{\mathbb{L}^2(\Omega)} \leq Ch^{k+1} \quad \text{and} \quad \|\bar{\gamma}\|_{\mathbb{L}^2(\Omega)} \leq C\tau^2. \quad (30)$$

This with (23) implies that

$$\|\partial^1 \eta_h^n\|_{\mathbb{L}^2(\Omega)} \leq C(h^{k+1} + \tau^2), \quad \forall n \in \llbracket 1, M+1 \rrbracket. \quad (31)$$

This completes the proof of new $W^{1,\infty}(L^2)$ -error estimate. ■

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Conclusion

We considered the heat equation (as model), with initial and homogeneous boundary conditions in any space dimension. A new error error estimate in $W^{1,\infty}(L^2)$ -norm is derived for an implicit Crank-Nicolson finite element scheme.

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Perspectives

- Is it possible to prove a convergence in $W^{1,\infty}(L^2)$ -norm under a weak-regularity ?
- Is it possible to extend the obtained error estimate to other complex equations (or systems), e.g. time dependent incompressible Navier-Stokes equations, in which heat equation is involved?

