# Basic properties of the Liouville quantum gravity metric for $\gamma \in (0, 2)$

#### Ewain Gwynne Based on 4 joint papers with J. Miller, 1 joint paper with J. Dubédat, H. Falconet, J. Pfeffer, and X. Sun, and 1 joint paper with J. Pfeffer

University of Cambridge

#### Outline



#### 2 KPZ formula







• Let  $\gamma \in (0,2)$  and let  $U \subset \mathbb{C}$ .



- Let  $\gamma \in (0,2)$  and let  $U \subset \mathbb{C}$ .
- A  $\gamma$ -Liouville quantum gravity (LQG) surface is the surface with Riemannian metric " $e^{\gamma h(z)} (dx^2 + dy^2)$ ", where h is a variant of the Gaussian free field on U.





- Let  $\gamma \in (0,2)$  and let  $U \subset \mathbb{C}$ .
- A  $\gamma$ -Liouville quantum gravity (LQG) surface is the surface with Riemannian metric " $e^{\gamma h(z)} (dx^2 + dy^2)$ ", where h is a variant of the Gaussian free field on U.
- Definition does not make literal sense since *h* is a distribution, not a function.



- Let  $\gamma \in (0,2)$  and let  $U \subset \mathbb{C}$ .
- A  $\gamma$ -Liouville quantum gravity (LQG) surface is the surface with Riemannian metric " $e^{\gamma h(z)} (dx^2 + dy^2)$ ", where h is a variant of the Gaussian free field on U.
- Definition does not make literal sense since *h* is a distribution, not a function.
- "Canonical random Riemannian metric".



- Let  $\gamma \in (0,2)$  and let  $U \subset \mathbb{C}$ .
- A  $\gamma$ -Liouville quantum gravity (LQG) surface is the surface with Riemannian metric " $e^{\gamma h(z)} (dx^2 + dy^2)$ ", where h is a variant of the Gaussian free field on U.
- Definition does not make literal sense since *h* is a distribution, not a function.
- "Canonical random Riemannian metric".
- Scaling limit of random planar maps.



- Let  $\gamma \in (0,2)$  and let  $U \subset \mathbb{C}$ .
- A  $\gamma$ -Liouville quantum gravity (LQG) surface is the surface with Riemannian metric " $e^{\gamma h(z)} (dx^2 + dy^2)$ ", where h is a variant of the Gaussian free field on U.
- Definition does not make literal sense since *h* is a distribution, not a function.
- "Canonical random Riemannian metric".
- Scaling limit of random planar maps.
  - $\gamma = \sqrt{8/3}$  for uniform random planar maps ("pure gravity").



- Let  $\gamma \in (0,2)$  and let  $U \subset \mathbb{C}$ .
- A  $\gamma$ -Liouville quantum gravity (LQG) surface is the surface with Riemannian metric " $e^{\gamma h(z)} (dx^2 + dy^2)$ ", where h is a variant of the Gaussian free field on U.
- Definition does not make literal sense since *h* is a distribution, not a function.
- "Canonical random Riemannian metric".
- Scaling limit of random planar maps.
  - $\gamma = \sqrt{8/3}$  for uniform random planar maps ("pure gravity").
  - $\gamma = \sqrt{2}$  for spanning-tree weighted maps.



- Let  $\gamma \in (0,2)$  and let  $U \subset \mathbb{C}$ .
- A  $\gamma$ -Liouville quantum gravity (LQG) surface is the surface with Riemannian metric " $e^{\gamma h(z)} (dx^2 + dy^2)$ ", where h is a variant of the Gaussian free field on U.
- Definition does not make literal sense since *h* is a distribution, not a function.
- "Canonical random Riemannian metric".
- Scaling limit of random planar maps.
  - $\gamma = \sqrt{8/3}$  for uniform random planar maps ("pure gravity").
  - $\gamma = \sqrt{2}$  for spanning-tree weighted maps.
  - $\gamma = \sqrt{3}$  for Ising-weighted maps.



Simulation for  $\gamma = 0.6$  (Miller)



Simulation for  $\gamma = 0.6$  (Miller)

• Definition / axiomatic characterization of the  $\gamma$ -LQG metric.



Simulation for  $\gamma = 0.6$  (Miller)

- Definition / axiomatic characterization of the  $\gamma\text{-LQG}$  metric.
- Basic properties of the metric.



Simulation for  $\gamma = 0.6$  (Miller)

- Definition / axiomatic characterization of the  $\gamma\text{-LQG}$  metric.
- Basic properties of the metric.
  - Conformal coordinate change.



Simulation for  $\gamma = 0.6$  (Miller)

- Definition / axiomatic characterization of the  $\gamma\text{-LQG}$  metric.
- Basic properties of the metric.
  - Conformal coordinate change.
  - KPZ formula.



Simulation for  $\gamma = 0.6$  (Miller)

- Definition / axiomatic characterization of the  $\gamma$ -LQG metric.
- Basic properties of the metric.
  - Conformal coordinate change.
  - KPZ formula.
  - Confluence of geodesics.



Simulation for  $\gamma = 0.6$  (Miller)

- Definition / axiomatic characterization of the  $\gamma$ -LQG metric.
- Basic properties of the metric.
  - Conformal coordinate change.
  - KPZ formula.
  - Confluence of geodesics.
- Open problems.



Simulation for  $\gamma = 0.6$  (Miller)

- Definition / axiomatic characterization of the  $\gamma$ -LQG metric.
- Basic properties of the metric.
  - Conformal coordinate change.
  - KPZ formula.
  - Confluence of geodesics.
- Open problems.
- Proofs are elementary: use only basic properties of the GFF.

• **Ding-Gwynne:**  $\exists d_{\gamma} > 2$  (the "Hausdorff dimension of  $\gamma$ -LQG").

- **Ding-Gwynne:**  $\exists d_{\gamma} > 2$  (the "Hausdorff dimension of  $\gamma$ -LQG").
  - e.g., ball volume exponent for random planar maps.

- **Ding-Gwynne:**  $\exists d_{\gamma} > 2$  (the "Hausdorff dimension of  $\gamma$ -LQG").
  - e.g., ball volume exponent for random planar maps.
  - Not known explicitly except that  $d_{\sqrt{8/3}} = 4$ .

• **Ding-Gwynne:**  $\exists d_{\gamma} > 2$  (the "Hausdorff dimension of  $\gamma$ -LQG").

- e.g., ball volume exponent for random planar maps.
- Not known explicitly except that  $d_{\sqrt{8/3}} = 4$ .

• Let  $\xi := \gamma/d_{\gamma}$ . Let  $h_{\varepsilon} = h *$  (heat kernel).  $\varepsilon$ -LFPP metric:

$$D_h^{\varepsilon}(z,w) = \inf_{P:z \to w} \int_0^1 e^{\xi h_{\varepsilon}(P(t))} |P'(t)| dt,$$

where the inf is over piecewise  $C^1$  paths from z to w.

• **Ding-Gwynne:**  $\exists d_{\gamma} > 2$  (the "Hausdorff dimension of  $\gamma$ -LQG").

- e.g., ball volume exponent for random planar maps.
- Not known explicitly except that  $d_{\sqrt{8/3}} = 4$ .

• Let  $\xi := \gamma/d_{\gamma}$ . Let  $h_{\varepsilon} = h *$  (heat kernel).  $\varepsilon$ -LFPP metric:

$$D_h^{\varepsilon}(z,w) = \inf_{P:z \to w} \int_0^1 e^{\xi h_{\varepsilon}(P(t))} |P'(t)| dt,$$

where the inf is over piecewise  $C^1$  paths from z to w.

ξ = γ/d<sub>γ</sub> because scaling areas by C (adding γ<sup>-1</sup> log C to h) corresponds to scaling distances by C<sup>1/d<sub>γ</sub></sup>.

• **Ding-Gwynne:**  $\exists d_{\gamma} > 2$  (the "Hausdorff dimension of  $\gamma$ -LQG").

- e.g., ball volume exponent for random planar maps.
- Not known explicitly except that  $d_{\sqrt{8/3}} = 4$ .

• Let  $\xi := \gamma/d_{\gamma}$ . Let  $h_{\varepsilon} = h *$  (heat kernel).  $\varepsilon$ -LFPP metric:

$$D_h^{\varepsilon}(z,w) = \inf_{P:z \to w} \int_0^1 e^{\xi h_{\varepsilon}(P(t))} |P'(t)| dt,$$

where the inf is over piecewise  $C^1$  paths from z to w.

ξ = γ/d<sub>γ</sub> because scaling areas by C (adding γ<sup>-1</sup> log C to h) corresponds to scaling distances by C<sup>1/d<sub>γ</sub></sup>.

**Theorem (Ding-Dubédat-Dunlap-Falconet, 2019):** The  $\varepsilon$ -LFPP metrics, re-scaled appropriately, are tight w.r.t. local uniform topology on  $\mathbb{C} \times \mathbb{C}$ . Every subsequential limit induces Euclidean topology.

• **Ding-Gwynne:**  $\exists d_{\gamma} > 2$  (the "Hausdorff dimension of  $\gamma$ -LQG").

- e.g., ball volume exponent for random planar maps.
- Not known explicitly except that  $d_{\sqrt{8/3}} = 4$ .

• Let  $\xi := \gamma/d_{\gamma}$ . Let  $h_{\varepsilon} = h *$  (heat kernel).  $\varepsilon$ -LFPP metric:

$$D_h^{\varepsilon}(z,w) = \inf_{P:z \to w} \int_0^1 e^{\xi h_{\varepsilon}(P(t))} |P'(t)| dt,$$

where the inf is over piecewise  $C^1$  paths from z to w.

ξ = γ/d<sub>γ</sub> because scaling areas by C (adding γ<sup>-1</sup> log C to h) corresponds to scaling distances by C<sup>1/d<sub>γ</sub></sup>.

**Theorem (Ding-Dubédat-Dunlap-Falconet, 2019):** The  $\varepsilon$ -LFPP metrics, re-scaled appropriately, are tight w.r.t. local uniform topology on  $\mathbb{C} \times \mathbb{C}$ . Every subsequential limit induces Euclidean topology.

**Theorem (Gwynne Miller, 2019):**  $\varepsilon$ -LFPP converges in probability to a conformally covariant metric  $D_h$ , the  $\gamma$ -LQG metric.

E. Gwynne (Cambridge)

A (strong)  $\gamma$ -LQG metric is a function  $h \mapsto D_h$  from distributions to metrics on  $\mathbb{C}$  which induce Eucl. topology satisfying the following.

A (strong)  $\gamma$ -LQG metric is a function  $h \mapsto D_h$  from distributions to metrics on  $\mathbb{C}$  which induce Eucl. topology satisfying the following.

• Length space. The  $D_h$ -distance between two points is the infimum of the  $D_h$ -lengths of continuous paths between them.

A (strong)  $\gamma$ -LQG metric is a function  $h \mapsto D_h$  from distributions to metrics on  $\mathbb{C}$  which induce Eucl. topology satisfying the following.

- Length space. The D<sub>h</sub>-distance between two points is the infimum of the D<sub>h</sub>-lengths of continuous paths between them.
- ② Locality. Let U ⊂ C be open. The D<sub>h</sub>-internal metric on U determined by h|<sub>U</sub>.

A (strong)  $\gamma$ -LQG metric is a function  $h \mapsto D_h$  from distributions to metrics on  $\mathbb{C}$  which induce Eucl. topology satisfying the following.

- Length space. The D<sub>h</sub>-distance between two points is the infimum of the D<sub>h</sub>-lengths of continuous paths between them.
- ② Locality. Let U ⊂ C be open. The D<sub>h</sub>-internal metric on U determined by h|<sub>U</sub>.
- **(3) Weyl scaling.** Let  $\xi = \gamma/d_{\gamma}$ . A.s.,  $\forall$  continuous  $f : \mathbb{C} \to \mathbb{R}$ ,

$$D_{h+f}(z,w) = \inf_{P:z\to w} \int_0^{\operatorname{len}(P;D_h)} e^{\xi f(P(t))} dt,$$

where the inf is over continuous paths parametrized by  $D_h$ -length.

A (strong)  $\gamma$ -LQG metric is a function  $h \mapsto D_h$  from distributions to metrics on  $\mathbb{C}$  which induce Eucl. topology satisfying the following.

- Length space. The D<sub>h</sub>-distance between two points is the infimum of the D<sub>h</sub>-lengths of continuous paths between them.
- ② Locality. Let U ⊂ C be open. The D<sub>h</sub>-internal metric on U determined by h|<sub>U</sub>.
- **(3) Weyl scaling.** Let  $\xi = \gamma/d_{\gamma}$ . A.s.,  $\forall$  continuous  $f : \mathbb{C} \to \mathbb{R}$ ,

$$D_{h+f}(z,w) = \inf_{P:z\to w} \int_0^{\operatorname{len}(P;D_h)} e^{\xi f(P(t))} dt,$$

where the inf is over continuous paths parametrized by  $D_h$ -length.

 Goordinate change for complex affine maps. Let Q = 2/γ + γ/2. For each fixed a ∈ C \ {0} and b ∈ C, a.s.

$$D_h(a \cdot +b, a \cdot +b) = D_{h(a \cdot +b)+Q \log |a|}(\cdot, \cdot).$$

#### Characterization of the $\gamma$ -LQG metric

# **Theorem (Gwynne-Miller, 2019):** The limit of the $\varepsilon$ -LFPP metrics is a $\gamma$ -LQG metric.



Simulation for  $\gamma = 0.9$  (Miller)

#### Characterization of the $\gamma$ -LQG metric

**Theorem (Gwynne-Miller, 2019):** The limit of the  $\varepsilon$ -LFPP metrics is a  $\gamma$ -LQG metric. Moreover, if  $D_h$  and  $\widetilde{D}_h$  are two  $\gamma$ -LQG metrics, then there is a deterministic C > 0 such that a.s.  $\widetilde{D}_h = CD_h$ .



Simulation for  $\gamma = 0.9$  (Miller)

Simulation for  $\gamma = 0.9$  (Miller)

**Theorem (Gwynne-Miller, 2019):** The limit of the  $\varepsilon$ -LFPP metrics is a  $\gamma$ -LQG metric. Moreover, if  $D_h$  and  $\widetilde{D}_h$  are two  $\gamma$ -LQG metrics, then there is a deterministic C > 0 such that a.s.  $\widetilde{D}_h = CD_h$ .

• Hence, we can refer to the  $\gamma$ -LQG metric.



Simulation for  $\gamma = 0.9$  (Miller)

**Theorem (Gwynne-Miller, 2019):** The limit of the  $\varepsilon$ -LFPP metrics is a  $\gamma$ -LQG metric. Moreover, if  $D_h$  and  $\widetilde{D}_h$  are two  $\gamma$ -LQG metrics, then there is a deterministic C > 0 such that a.s.  $\widetilde{D}_h = CD_h$ .

- Hence, we can refer to the  $\gamma\text{-}\mathsf{LQG}$  metric.
- Recall: Miller and Sheffield (2016) constructed a  $\sqrt{8/3}$ -LQG metric using QLE.



Simulation for  $\gamma = 0.9$  (Miller)

**Theorem (Gwynne-Miller, 2019):** The limit of the  $\varepsilon$ -LFPP metrics is a  $\gamma$ -LQG metric. Moreover, if  $D_h$  and  $\widetilde{D}_h$  are two  $\gamma$ -LQG metrics, then there is a deterministic C > 0 such that a.s.  $\widetilde{D}_h = CD_h$ .

- Hence, we can refer to the  $\gamma\text{-}\mathsf{LQG}$  metric.
- Recall: Miller and Sheffield (2016) constructed a  $\sqrt{8/3}$ -LQG metric using QLE.
- $\sqrt{8/3}$ -LQG surface = Brownian map = scaling limit of uniform random planar maps w.r.t. Gromov-Hausdorff distance.
# Characterization of the $\gamma\text{-}\text{LQG}$ metric

Simulation for  $\gamma = 0.9$  (Miller)

**Theorem (Gwynne-Miller, 2019):** The limit of the  $\varepsilon$ -LFPP metrics is a  $\gamma$ -LQG metric. Moreover, if  $D_h$  and  $\widetilde{D}_h$  are two  $\gamma$ -LQG metrics, then there is a deterministic C > 0 such that a.s.  $\widetilde{D}_h = CD_h$ .

- Hence, we can refer to the  $\gamma\text{-}\mathsf{LQG}$  metric.
- Recall: Miller and Sheffield (2016) constructed a  $\sqrt{8/3}$ -LQG metric using QLE.
- $\sqrt{8/3}$ -LQG surface = Brownian map = scaling limit of uniform random planar maps w.r.t. Gromov-Hausdorff distance.

**Corollary:** The Miller-Sheffield  $\sqrt{8/3}$ -LQG metric agrees with the limit of  $\sqrt{8/3}$ -LFPP.

# Characterization of the $\gamma$ -LQG metric

Simulation for  $\gamma = 0.9$  (Miller)

**Theorem (Gwynne-Miller, 2019):** The limit of the  $\varepsilon$ -LFPP metrics is a  $\gamma$ -LQG metric. Moreover, if  $D_h$  and  $\widetilde{D}_h$  are two  $\gamma$ -LQG metrics, then there is a deterministic C > 0 such that a.s.  $\widetilde{D}_h = CD_h$ .

- Hence, we can refer to the  $\gamma$ -LQG metric.
- Recall: Miller and Sheffield (2016) constructed a  $\sqrt{8/3}$ -LQG metric using QLE.
- $\sqrt{8/3}$ -LQG surface = Brownian map = scaling limit of uniform random planar maps w.r.t. Gromov-Hausdorff distance.

**Corollary:** The Miller-Sheffield  $\sqrt{8/3}$ -LQG metric agrees with the limit of  $\sqrt{8/3}$ -LFPP. **Conjecture:** For general  $\gamma \in (0, 2)$ , the  $\gamma$ -LQG metric is the scaling limit of weighted random planar maps w.r.t. the Gromov-Hausdorff topology.

#### Metrics on other domains

If U ⊂ C and h is a GFF on U, can define D<sub>h</sub> by local absolute continuity.

#### Metrics on other domains

If U ⊂ C and h is a GFF on U, can define D<sub>h</sub> by local absolute continuity.

Coordinate change (Gwynne-Miller, 2019): If  $\phi: \widetilde{U} \to U$  is a conformal map and

$$\widetilde{h} = h \circ \phi + Q \log |\phi'| \quad ext{for} \quad Q = rac{2}{\gamma} + rac{\gamma}{2},$$

then a.s.  $D_{\widetilde{h}}(z,w) = D_h(\phi(z),\phi(w)).$ 



#### Outline

#### 1) The $\gamma$ -LQG metric







• For  $X \subset \mathbb{C}$ , let dim<sup>0</sup><sub>H</sub> X and dim<sup> $\gamma$ </sup><sub>H</sub> X be the Hausdorff dimension of X w.r.t. the Euclidean and  $\gamma$ -LQG metrics, resp.



Simulation for  $\gamma = 2$  (Miller)

 For X ⊂ C, let dim<sup>0</sup><sub>H</sub> X and dim<sup>γ</sup><sub>H</sub> X be the Hausdorff dimension of X w.r.t. the Euclidean and γ-LQG metrics, resp.

Simulation for  $\gamma = 2$  (Miller)

**KPZ formula (Gwynne-Pfeffer, 2019):** If X is a random Borel set independent from h, then a.s.

$$\dim^0_{\mathcal{H}} X = rac{\gamma}{d_\gamma} Q \dim^\gamma_{\mathcal{H}} X - rac{\gamma^2}{2d_\gamma^2} (\dim^\gamma_{\mathcal{H}} X)^2.$$

 For X ⊂ C, let dim<sup>0</sup><sub>H</sub> X and dim<sup>γ</sup><sub>H</sub> X be the Hausdorff dimension of X w.r.t. the Euclidean and γ-LQG metrics, resp.



Simulation for  $\gamma = 2$  (Miller)

**KPZ formula (Gwynne-Pfeffer, 2019):** If X is a random Borel set independent from h, then a.s.

$$\dim_{\mathcal{H}}^{0} X = \frac{\gamma}{d_{\gamma}} Q \dim_{\mathcal{H}}^{\gamma} X - \frac{\gamma^{2}}{2d_{\gamma}^{2}} (\dim_{\mathcal{H}}^{\gamma} X)^{2}.$$

• Agrees with other versions of KPZ (e.g., Duplantier-Sheffield) if we use the re-scaled dimension  $\frac{1}{d_{\gamma}}(\dim_{\mathcal{H}}^{\gamma}X)$ .

 For X ⊂ C, let dim<sup>0</sup><sub>H</sub> X and dim<sup>γ</sup><sub>H</sub> X be the Hausdorff dimension of X w.r.t. the Euclidean and γ-LQG metrics, resp.



Simulation for  $\gamma = 2$  (Miller)

**KPZ formula (Gwynne-Pfeffer, 2019):** If X is a random Borel set independent from h, then a.s.

$$\dim_{\mathcal{H}}^{0} X = \frac{\gamma}{d_{\gamma}} Q \dim_{\mathcal{H}}^{\gamma} X - \frac{\gamma^{2}}{2d_{\gamma}^{2}} (\dim_{\mathcal{H}}^{\gamma} X)^{2}.$$

• Agrees with other versions of KPZ (e.g., Duplantier-Sheffield) if we use the re-scaled dimension  $\frac{1}{d_{\gamma}}(\dim_{\mathcal{H}}^{\gamma}X)$ .

**Corollary:** dim $_{\mathcal{H}}^{\gamma} \mathbb{C} = d_{\gamma}$ .

$$\dim_{\mathcal{H}}^{\gamma} X \leq \begin{cases} \frac{\dim_{\mathcal{H}}^{0} X}{\frac{\gamma}{d_{\gamma}} \left( Q - \sqrt{4 - 2 \dim_{\mathcal{H}}^{0} X} \right)}, & \text{if } \dim_{\mathcal{H}}^{0} X < 2 - \frac{\gamma}{2} \\ d_{\gamma}, & \text{if } \dim_{\mathcal{H}}^{0} X \geq 2 - \frac{\gamma}{2} \end{cases}$$

$$\dim_{\mathcal{H}}^{\gamma} X \leq \begin{cases} \frac{\dim_{\mathcal{H}}^{0} X}{\frac{\gamma}{d_{\gamma}} \left( Q - \sqrt{4 - 2 \dim_{\mathcal{H}}^{0} X} \right)}, & \text{ if } \dim_{\mathcal{H}}^{0} X < 2 - \frac{\gamma^{2}}{2} \\ d_{\gamma}, & \text{ if } \dim_{\mathcal{H}}^{0} X \geq 2 - \frac{\gamma^{2}}{2} \end{cases}$$

$$\begin{split} \dim_{\mathcal{H}}^{0} X \\ &\leq \begin{cases} \frac{\gamma}{d_{\gamma}} \dim_{\mathcal{H}}^{\gamma} X \Big( Q - \frac{\gamma}{d_{\gamma}} \dim_{\mathcal{H}}^{\gamma} X + \sqrt{4 - \frac{2\gamma Q}{d_{\gamma}} \dim_{\mathcal{H}}^{\gamma} X + \frac{\gamma^{2}}{d_{\gamma}^{2}} (\dim_{\mathcal{H}}^{\gamma} X)^{2}} \Big), & \text{if } \dim_{\mathcal{H}}^{\gamma} X < \frac{2d_{\gamma}}{\gamma Q} \\ 2, & \text{if } \dim_{\mathcal{H}}^{\gamma} X \geq \frac{2d_{\gamma}}{\gamma Q} \end{cases} \end{split}$$

$$\dim_{\mathcal{H}}^{\gamma} X \leq \begin{cases} \frac{\dim_{\mathcal{H}}^{0} X}{\frac{\gamma}{d_{\gamma}} \left( Q - \sqrt{4 - 2 \dim_{\mathcal{H}}^{0} X} \right)}, & \text{ if } \dim_{\mathcal{H}}^{0} X < 2 - \frac{\gamma^{2}}{2} \\ d_{\gamma}, & \text{ if } \dim_{\mathcal{H}}^{0} X \geq 2 - \frac{\gamma^{2}}{2} \end{cases}$$



$$\dim_{\mathcal{H}}^{\gamma} X \leq \begin{cases} \frac{\dim_{\mathcal{H}}^{0} X}{\frac{\gamma}{d_{\gamma}} \left( Q - \sqrt{4 - 2 \dim_{\mathcal{H}}^{0} X} \right)}, & \text{ if } \dim_{\mathcal{H}}^{0} X < 2 - \frac{\gamma^{2}}{2} \\ d_{\gamma}, & \text{ if } \dim_{\mathcal{H}}^{0} X \geq 2 - \frac{\gamma^{2}}{2} \end{cases}$$







• Idea of proof: assume that the  $2^{-n} \times 2^{-n}$  squares which intersect X are the ones with the largest  $D_h$ -diameters. Do the reverse for squares with "quantum size"  $2^{-m}$ .



• Idea of proof: assume that the  $2^{-n} \times 2^{-n}$  squares which intersect X are the ones with the largest  $D_h$ -diameters. Do the reverse for squares with "quantum size"  $2^{-m}$ .

• Upper bound for dim $^{0}_{\mathcal{H}}X$  is closest to optimal when *h* is "small" on *X*.



• Idea of proof: assume that the  $2^{-n} \times 2^{-n}$  squares which intersect X are the ones with the largest  $D_h$ -diameters. Do the reverse for squares with "quantum size"  $2^{-m}$ .

• Upper bound for dim $_{\mathcal{H}}^0 X$  is closest to optimal when *h* is "small" on *X*.

• A  $D_h$ -geodesic has  $D_h$ -dimension 1.



• Idea of proof: assume that the  $2^{-n} \times 2^{-n}$  squares which intersect X are the ones with the largest  $D_h$ -diameters. Do the reverse for squares with "quantum size"  $2^{-m}$ .

• Upper bound for dim $_{\mathcal{H}}^0 X$  is closest to optimal when *h* is "small" on *X*.

• A  $D_h$ -geodesic has  $D_h$ -dimension 1. **Corollary:** A.s., every  $D_h$ -geodesic has Eucl. dimension at most

$$rac{\gamma}{d_\gamma}igg(Q-rac{\gamma}{d_\gamma}+\sqrt{4-rac{2\gamma Q}{d_\gamma}+rac{\gamma^2}{d_\gamma^2}}igg)$$



• Idea of proof: assume that the  $2^{-n} \times 2^{-n}$  squares which intersect X are the ones with the largest  $D_h$ -diameters. Do the reverse for squares with "quantum size"  $2^{-m}$ .

• Upper bound for dim $_{\mathcal{H}}^0 X$  is closest to optimal when *h* is "small" on *X*.

• A  $D_h$ -geodesic has  $D_h$ -dimension 1. **Corollary:** A.s., every  $D_h$ -geodesic has Eucl. dimension at most

$$rac{\gamma}{d_{\gamma}}\left( \mathcal{Q} - rac{\gamma}{d_{\gamma}} + \sqrt{4 - rac{2\gamma\mathcal{Q}}{d_{\gamma}} + rac{\gamma^2}{d_{\gamma}^2}} 
ight)$$

• Equals  $\frac{4+\sqrt{15}}{6} \approx 1.312$  for  $\gamma = \sqrt{8/3}$ .

#### Outline

1) The  $\gamma$ -LQG metric







# **Theorem (Gwynne-Miller, 2019):** For every s > 0, $\exists t \in (0, s)$ such that all $D_h$ -geodesics from 0 to points outside $\mathcal{B}_s(z; D_h)$ coincide on [0, t].



 $\mathcal{B}^{\bullet}_{s}(z; D_{h}$ 



 $\tilde{B}_{1}(z; D_{h})$ 

**Theorem (Gwynne-Miller, 2019):** For every s > 0,  $\exists t \in (0, s)$  such that all  $D_h$ -geodesics from 0 to points outside  $\mathcal{B}_s(z; D_h)$  coincide on [0, t].

• Very different from behavior of geodesics for a smooth Riemannian metric.



**Theorem (Gwynne-Miller, 2019):** For every s > 0,  $\exists t \in (0, s)$  such that all  $D_h$ -geodesics from 0 to points outside  $\mathcal{B}_s(z; D_h)$  coincide on [0, t].

- Very different from behavior of geodesics for a smooth Riemannian metric.
- Proven by Le Gall (2010) for the Brownian map (using very different methods).





**Theorem (Gwynne-Miller, 2019):** For every s > 0,  $\exists t \in (0, s)$  such that all  $D_h$ -geodesics from 0 to points outside  $\mathcal{B}_s(z; D_h)$  coincide on [0, t].

- Very different from behavior of geodesics for a smooth Riemannian metric.
- Proven by Le Gall (2010) for the Brownian map (using very different methods).
- Key tool in the proof of uniqueness of the metric.



**Theorem (Gwynne-Miller, 2019):** For every s > 0,  $\exists t \in (0, s)$  such that all  $D_h$ -geodesics from 0 to points outside  $\mathcal{B}_s(z; D_h)$  coincide on [0, t].

- Very different from behavior of geodesics for a smooth Riemannian metric.
- Proven by Le Gall (2010) for the Brownian map (using very different methods).
- Key tool in the proof of uniqueness of the metric.
- Let  $\mathcal{B}_{s}^{\bullet}(z; D_{h})$  be the union of  $\mathcal{B}_{s}(z; D_{h})$  and the "holes" which it disconnects from  $\infty$ .



**Theorem (Gwynne-Miller, 2019):** For every s > 0,  $\exists t \in (0, s)$  such that all  $D_h$ -geodesics from 0 to points outside  $\mathcal{B}_s(z; D_h)$  coincide on [0, t].

- Very different from behavior of geodesics for a smooth Riemannian metric.
- Proven by Le Gall (2010) for the Brownian map (using very different methods).
- Key tool in the proof of uniqueness of the metric.
- Let B<sup>●</sup><sub>s</sub>(z; D<sub>h</sub>) be the union of B<sub>s</sub>(z; D<sub>h</sub>) and the "holes" which it disconnects from ∞.

**Theorem (Gwynne-Miller, 2019):** For any  $0 < t < s < \infty$ , there are only finitely many possibilities for  $P|_{[0,t]}$ , where P is a leftmost  $D_{h}$ -geodesic from 0 to  $\partial \mathcal{B}_{s}^{\bullet}(z; D_{h})$ .

# Step 1: Partition $\partial B_t$ into small arcs



#### Step 2: each arc is "killed off" with positive probability



#### Step 2: each arc is "killed off" with positive probability



#### Step 2: each arc is "killed off" with positive probability



#### Step 3: increasing radius "kills off" positive fraction of arcs



# Step 4: iterate logarithmically many times



# Confluence implies independence along a geodesic



• A  $D_h$ -geodesic P is not a local set for h.

#### Confluence implies independence along a geodesic



- A  $D_h$ -geodesic P is not a local set for h.
- Should still have long-range independence for behavior of *h* near different points of *P*.

#### Confluence implies independence along a geodesic



- A  $D_h$ -geodesic P is not a local set for h.
- Should still have long-range independence for behavior of *h* near different points of *P*.
- For a typical time s ∈ [0, D<sub>h</sub>(z, w)], all D<sub>h</sub>-geodesics from ∂B<sub>ε</sub>(P(t)) to z agree on [0, s − ε<sup>β</sup>].
#### Confluence implies independence along a geodesic



- A  $D_h$ -geodesic P is not a local set for h.
- Should still have long-range independence for behavior of *h* near different points of *P*.
- For a typical time s ∈ [0, D<sub>h</sub>(z, w)], all D<sub>h</sub>-geodesics from ∂B<sub>ε</sub>(P(t)) to z agree on [0, s − ε<sup>β</sup>].
- If we make a small change to  $h|_{B_{\varepsilon}(P(s))}$ , we do not change  $P|_{[0,s-\varepsilon^{\beta}]}$ .

#### Confluence implies independence along a geodesic



- A  $D_h$ -geodesic P is not a local set for h.
- Should still have long-range independence for behavior of *h* near different points of *P*.
- For a typical time s ∈ [0, D<sub>h</sub>(z, w)], all D<sub>h</sub>-geodesics from ∂B<sub>ε</sub>(P(t)) to z agree on [0, s − ε<sup>β</sup>].
- If we make a small change to h|<sub>Bε</sub>(P(s)), we do not change P|<sub>[0,s-εβ]</sub>.
- Hence, conditional law of h|<sub>B<sub>ε</sub>(P(s))</sub> given P|<sub>[0,s-ε<sup>β</sup>]</sub> is not much different from its marginal law.

#### Outline

1) The  $\gamma$ -LQG metric







$$D_h^arepsilon(z,w) = \min_{P:z o w} \int_0^1 e^{\xi h_arepsilon(P(t))} |P'(t)| \, dt, \quad \xi = rac{\gamma}{d_\gamma}.$$

γ-LQG metric for each γ ∈ (0, 2) constructed as the limit (in probability) of LFPP:

$$D_h^arepsilon(z,w) = \min_{P:z
ightarrow w} \int_0^1 e^{\xi h_arepsilon(P(t))} |P'(t)| \, dt, \quad \xi = rac{\gamma}{d_\gamma}.$$

• Uniquely characterized by

$$D_h^arepsilon(z,w) = \min_{P:z o w} \int_0^1 e^{\xi h_arepsilon(P(t))} |P'(t)| \, dt, \quad \xi = rac{\gamma}{d_\gamma}.$$

- Uniquely characterized by
  - Length space.

$$D_h^arepsilon(z,w) = \min_{P:z o w} \int_0^1 e^{\xi h_arepsilon(P(t))} |P'(t)| \, dt, \quad \xi = rac{\gamma}{d_\gamma}.$$

- Uniquely characterized by
  - Length space.
  - Locality.

$$D_h^{\varepsilon}(z,w) = \min_{P:z \to w} \int_0^1 e^{\xi h_{\varepsilon}(P(t))} |P'(t)| dt, \quad \xi = \frac{\gamma}{d_{\gamma}}.$$

- Uniquely characterized by
  - Length space.
  - Locality.
  - Weyl scaling (behavior when adding a function to *h*).

$$D_h^{\varepsilon}(z,w) = \min_{P:z \to w} \int_0^1 e^{\xi h_{\varepsilon}(P(t))} |P'(t)| dt, \quad \xi = \frac{\gamma}{d_{\gamma}}.$$

- Uniquely characterized by
  - Length space.
  - Locality.
  - Weyl scaling (behavior when adding a function to *h*).
  - Coordinate change (behavior when scaling/translating space).

$$D_h^{\varepsilon}(z,w) = \min_{P:z \to w} \int_0^1 e^{\xi h_{\varepsilon}(P(t))} |P'(t)| dt, \quad \xi = \frac{\gamma}{d_{\gamma}}.$$

- Uniquely characterized by
  - Length space.
  - Locality.
  - Weyl scaling (behavior when adding a function to *h*).
  - Coordinate change (behavior when scaling/translating space).
- Properties:

$$D_h^{\varepsilon}(z,w) = \min_{P:z \to w} \int_0^1 e^{\xi h_{\varepsilon}(P(t))} |P'(t)| dt, \quad \xi = \frac{\gamma}{d_{\gamma}}.$$

- Uniquely characterized by
  - Length space.
  - Locality.
  - Weyl scaling (behavior when adding a function to h).
  - Coordinate change (behavior when scaling/translating space).
- Properties:
  - Conformal coordinate change.

$$D_h^{\varepsilon}(z,w) = \min_{P:z \to w} \int_0^1 e^{\xi h_{\varepsilon}(P(t))} |P'(t)| dt, \quad \xi = \frac{\gamma}{d_{\gamma}}.$$

- Uniquely characterized by
  - Length space.
  - Locality.
  - Weyl scaling (behavior when adding a function to h).
  - Coordinate change (behavior when scaling/translating space).
- Properties:
  - Conformal coordinate change.
  - KPZ formula.

γ-LQG metric for each γ ∈ (0, 2) constructed as the limit (in probability) of LFPP:

$$D_h^{\varepsilon}(z,w) = \min_{P:z \to w} \int_0^1 e^{\xi h_{\varepsilon}(P(t))} |P'(t)| dt, \quad \xi = \frac{\gamma}{d_{\gamma}}.$$

- Uniquely characterized by
  - Length space.
  - Locality.
  - Weyl scaling (behavior when adding a function to h).
  - Coordinate change (behavior when scaling/translating space).
- Properties:
  - Conformal coordinate change.
  - KPZ formula.
  - Confluence of geodesics.

#### The LQG metric



• Compute the Hausdorff dimension  $d_{\gamma}$  for  $\gamma \neq \sqrt{8/3}$ .



- Compute the Hausdorff dimension  $d_{\gamma}$  for  $\gamma \neq \sqrt{8/3}$ .
- Compute the Euclidean dimension of LQG geodesics, outer boundaries of LQG metric balls, etc.



- Compute the Hausdorff dimension  $d_{\gamma}$  for  $\gamma \neq \sqrt{8/3}$ .
- Compute the Euclidean dimension of LQG geodesics, outer boundaries of LQG metric balls, etc.
- Show that the γ-LQG metric is the scaling limit of random planar maps for general γ ∈ (0, 2) (w.r.t. Gromov-Hausdorff).

• Construct the metric in the critical case  $\gamma = 2$ , or more generally for  $\xi > 2/d_2$  ("central charge > 1"; see Gwynne-Holden-Pfeffer-Remy, 2019).

- Construct the metric in the critical case  $\gamma = 2$ , or more generally for  $\xi > 2/d_2$  ("central charge > 1"; see Gwynne-Holden-Pfeffer-Remy, 2019).
- Can we see that  $\gamma = \sqrt{8/3}$  is special just from the axioms?

- Construct the metric in the critical case  $\gamma = 2$ , or more generally for  $\xi > 2/d_2$  ("central charge > 1"; see Gwynne-Holden-Pfeffer-Remy, 2019).
- Can we see that  $\gamma = \sqrt{8/3}$  is special just from the axioms?
- Are there any exact formulas for LQG distances, analogous to exact formulas for correlation functions / LQG areas by Kupiannen-Rhodes-Vargas, Remy, et. al. or exact formulas for Brownian surfaces?

- Construct the metric in the critical case  $\gamma = 2$ , or more generally for  $\xi > 2/d_2$  ("central charge > 1"; see Gwynne-Holden-Pfeffer-Remy, 2019).
- Can we see that  $\gamma = \sqrt{8/3}$  is special just from the axioms?
- Are there any exact formulas for LQG distances, analogous to exact formulas for correlation functions / LQG areas by Kupiannen-Rhodes-Vargas, Remy, et. al. or exact formulas for Brownian surfaces?
- Is there a general theory of random metrics obtained by exponentiating the GFF analogous to Gaussian multiplicative chaos?