Ten Big Problems

Scott Sheffield

Massachusetts Institute of Technology

July 5, 2019

- At least 4 fundamentally distinct (but equivalent) ways to construct pure-LQG/Brownian sphere *directly in the continuum*. Only 2 work for general γ. Are there any more?
 - a. Brownian snake: Glue Brownian-snake-excursion trees. (Chassaing, Schaeffer, Markert and Mokkadem, Le Gall, Miermont)
 - b. **CSBP:** Branching continuous state branching process. (Angel, Curien, Miller and S., etc.)
 - c. **GFF/GMC:** Fix boundary conditions and take large area limit, or use Bessel excursions to build "quantum sphere" or use LCFT. (Duplantier, Miller, S., David, Rhodes, Vargas, Kupiannen, Aru, Huang, Sun, etc.)
 - d. **Correlated CRT:** Mating trees derived from Brownian quadrant excursion. (Duplantier, Miller, S., etc.)

- At least 4 fundamentally distinct (but equivalent) ways to construct pure-LQG/Brownian sphere *directly in the continuum*. Only 2 work for general γ. Are there any more?
 - a. Brownian snake: Glue Brownian-snake-excursion trees. (Chassaing, Schaeffer, Markert and Mokkadem, Le Gall, Miermont)
 - b. **CSBP:** Branching continuous state branching process. (Angel, Curien, Miller and S., etc.)
 - c. **GFF/GMC:** Fix boundary conditions and take large area limit, or use Bessel excursions to build "quantum sphere" or use LCFT. (Duplantier, Miller, S., David, Rhodes, Vargas, Kupiannen, Aru, Huang, Sun, etc.)
 - d. Correlated CRT: Mating trees derived from Brownian quadrant excursion. (Duplantier, Miller, S., etc.)
- Spectacular recent work on metrics (Ding, Dubédat, Dunlap, Falconet, Gwynne, Miller, etc.) raises many questions.

- At least 4 fundamentally distinct (but equivalent) ways to construct pure-LQG/Brownian sphere *directly in the continuum*. Only 2 work for general γ. Are there any more?
 - a. Brownian snake: Glue Brownian-snake-excursion trees. (Chassaing, Schaeffer, Markert and Mokkadem, Le Gall, Miermont)
 - b. **CSBP:** Branching continuous state branching process. (Angel, Curien, Miller and S., etc.)
 - c. **GFF/GMC:** Fix boundary conditions and take large area limit, or use Bessel excursions to build "quantum sphere" or use LCFT. (Duplantier, Miller, S., David, Rhodes, Vargas, Kupiannen, Aru, Huang, Sun, etc.)
 - d. Correlated CRT: Mating trees derived from Brownian quadrant excursion. (Duplantier, Miller, S., etc.)
- Spectacular recent work on metrics (Ding, Dubédat, Dunlap, Falconet, Gwynne, Miller, etc.) raises many questions.
- Can we understand geodesic tree *decorated* by statistical physics data, such as loop data or lsing spin values?

- At least 4 fundamentally distinct (but equivalent) ways to construct pure-LQG/Brownian sphere *directly in the continuum*. Only 2 work for general *γ*. Are there any more?
 - a. Brownian snake: Glue Brownian-snake-excursion trees. (Chassaing, Schaeffer, Markert and Mokkadem, Le Gall, Miermont)
 - b. **CSBP:** Branching continuous state branching process. (Angel, Curien, Miller and S., etc.)
 - c. **GFF/GMC:** Fix boundary conditions and take large area limit, or use Bessel excursions to build "quantum sphere" or use LCFT. (Duplantier, Miller, S., David, Rhodes, Vargas, Kupiannen, Aru, Huang, Sun, etc.)
 - d. Correlated CRT: Mating trees derived from Brownian quadrant excursion. (Duplantier, Miller, S., etc.)
- Spectacular recent work on metrics (Ding, Dubédat, Dunlap, Falconet, Gwynne, Miller, etc.) raises many questions.
- Can we understand geodesic tree *decorated* by statistical physics data, such as loop data or lsing spin values?
- CRT mating *decorated* by metric information (one sided or two sided)?

- At least 4 fundamentally distinct (but equivalent) ways to construct pure-LQG/Brownian sphere *directly in the continuum*. Only 2 work for general *γ*. Are there any more?
 - a. Brownian snake: Glue Brownian-snake-excursion trees. (Chassaing, Schaeffer, Markert and Mokkadem, Le Gall, Miermont)
 - b. **CSBP:** Branching continuous state branching process. (Angel, Curien, Miller and S., etc.)
 - c. **GFF/GMC:** Fix boundary conditions and take large area limit, or use Bessel excursions to build "quantum sphere" or use LCFT. (Duplantier, Miller, S., David, Rhodes, Vargas, Kupiannen, Aru, Huang, Sun, etc.)
 - d. Correlated CRT: Mating trees derived from Brownian quadrant excursion. (Duplantier, Miller, S., etc.)
- Spectacular recent work on metrics (Ding, Dubédat, Dunlap, Falconet, Gwynne, Miller, etc.) raises many questions.
- Can we understand geodesic tree *decorated* by statistical physics data, such as loop data or lsing spin values?
- CRT mating *decorated* by metric information (one sided or two sided)?
- Can we at least show relevant processes are Markovian in continuum?

- At least 4 fundamentally distinct (but equivalent) ways to construct pure-LQG/Brownian sphere *directly in the continuum*. Only 2 work for general *γ*. Are there any more?
 - a. Brownian snake: Glue Brownian-snake-excursion trees. (Chassaing, Schaeffer, Markert and Mokkadem, Le Gall, Miermont)
 - b. **CSBP:** Branching continuous state branching process. (Angel, Curien, Miller and S., etc.)
 - c. **GFF/GMC:** Fix boundary conditions and take large area limit, or use Bessel excursions to build "quantum sphere" or use LCFT. (Duplantier, Miller, S., David, Rhodes, Vargas, Kupiannen, Aru, Huang, Sun, etc.)
 - d. Correlated CRT: Mating trees derived from Brownian quadrant excursion. (Duplantier, Miller, S., etc.)
- Spectacular recent work on metrics (Ding, Dubédat, Dunlap, Falconet, Gwynne, Miller, etc.) raises many questions.
- Can we understand geodesic tree *decorated* by statistical physics data, such as loop data or Ising spin values?
- CRT mating *decorated* by metric information (one sided or two sided)?
- Can we at least show relevant processes are Markovian in continuum?
- Any approach that allows one to compute things more explicitly?

- At least 4 fundamentally distinct (but equivalent) ways to construct pure-LQG/Brownian sphere *directly in the continuum*. Only 2 work for general γ. Are there any more?
 - a. Brownian snake: Glue Brownian-snake-excursion trees. (Chassaing, Schaeffer, Markert and Mokkadem, Le Gall, Miermont)
 - b. **CSBP:** Branching continuous state branching process. (Angel, Curien, Miller and S., etc.)
 - c. **GFF/GMC:** Fix boundary conditions and take large area limit, or use Bessel excursions to build "quantum sphere" or use LCFT. (Duplantier, Miller, S., David, Rhodes, Vargas, Kupiannen, Aru, Huang, Sun, etc.)
 - d. **Correlated CRT:** Mating trees derived from Brownian quadrant excursion. (Duplantier, Miller, S., etc.)
- Spectacular recent work on metrics (Ding, Dubédat, Dunlap, Falconet, Gwynne, Miller, etc.) raises many questions.
- Can we understand geodesic tree *decorated* by statistical physics data, such as loop data or lsing spin values?
- CRT mating *decorated* by metric information (one sided or two sided)?
- Can we at least show relevant processes are Markovian in continuum?
- Any approach that allows one to compute things more explicitly?
- Classic problem: find tractable description of Perron-Frobenius eigenvector.

What else is out there?



Imaginary Geometry

C-Yt house

X_t man t

 $C - Y_t$



Identify points on the graph of X if they are connected by a horizontal line which is below the graph; yields a continuum random tree (CRT)



- Identify points on the graph of X if they are connected by a horizontal line which is below the graph; yields a continuum random tree (CRT)
- Same for $C Y_t$ yields an independent CRT



- Identify points on the graph of X if they are connected by a horizontal line which is below the graph; yields a continuum random tree (CRT)
- Same for $C Y_t$ yields an independent CRT
- Glue the CRTs together by declaring points on the vertical lines to be equivalent



- Identify points on the graph of X if they are connected by a horizontal line which is below the graph; yields a continuum random tree (CRT)
- Same for $C Y_t$ yields an independent CRT
- Glue the CRTs together by declaring points on the vertical lines to be equivalent
- **Q:** What is the resulting structure?



- Identify points on the graph of X if they are connected by a horizontal line which is below the graph; yields a continuum random tree (CRT)
- Same for $C Y_t$ yields an independent CRT
- Glue the CRTs together by declaring points on the vertical lines to be equivalent
- **Q**: What is the resulting structure? **A**: Sphere with a space-filling path.

X, Y independent Brownian excursions on [0,1]. Pick C > 0 large so that the graphs of X and C - Y are disjoint.



- Identify points on the graph of X if they are connected by a horizontal line which is below the graph; yields a continuum random tree (CRT)
- Same for $C Y_t$ yields an independent CRT
- Glue the CRTs together by declaring points on the vertical lines to be equivalent

Q: What is the resulting structure? **A**: Sphere with a space-filling path. A peanosphere.



- 1. The dancing snake has a scaling limit called the Brownian snake.
- 2. The x and y coordinates of the Brownian snake's head are two functions.
- 3. Each of these describes a tree (via the same construction we used to make CRT from Brownian motion).
- 4. Gluing these two trees together gives a random surface called the **Brownian map**.

Take four basic constructions of Brownian sphere (from Brownian snake, from branching CSBP, from GFF, from Brownian quadrant excursion) and give a torus version of each one of them.

- Take four basic constructions of Brownian sphere (from Brownian snake, from branching CSBP, from GFF, from Brownian quadrant excursion) and give a torus version of each one of them.
- The work has begun. See papers by Rhodes and Vargas plus Guillarmou and/or David and/or Kupiannen. See forthcoming work by Bettinelli, Miermont.

- Take four basic constructions of Brownian sphere (from Brownian snake, from branching CSBP, from GFF, from Brownian quadrant excursion) and give a torus version of each one of them.
- The work has begun. See papers by Rhodes and Vargas plus Guillarmou and/or David and/or Kupiannen. See forthcoming work by Bettinelli, Miermont.
- Understand higher genus analogs of the 4 basic constructions.

- Take four basic constructions of Brownian sphere (from Brownian snake, from branching CSBP, from GFF, from Brownian quadrant excursion) and give a torus version of each one of them.
- The work has begun. See papers by Rhodes and Vargas plus Guillarmou and/or David and/or Kupiannen. See forthcoming work by Bettinelli, Miermont.
- Understand higher genus analogs of the 4 basic constructions.
- Prove their equivalence. In particular, law of conformal modulus should be same in each approach.

Big Problem 3: Understand Growth and QLE

Including isotropic Euclidean DLA.

Big Problem 3: Understand Growth and QLE

- Including isotropic Euclidean DLA.
- ▶ Prove anything. Existence of subsequential limit satisfying QLE definition, bounds on dimensions, exact dimensions, basic properties, etc. Check that metric growth for $\gamma \neq \sqrt{8/3}$ is a QLE. Formulate general case using random matrix variants? SLE variants?



Eden model on $\sqrt{8/3}\text{-}\mathsf{LQG}$



DLA on a $\sqrt{2}$ -LQG

QLE(8/3,0) is a member of a two-parameter family of processes called $QLE(\gamma^2,\eta)$

- $\blacktriangleright~\gamma$ is the type of LQG surface on which the process grows
- $\blacktriangleright \eta$ determines the manner in which it grows

QLE(8/3,0) is a member of a two-parameter family of processes called $QLE(\gamma^2,\eta)$

- $\blacktriangleright~\gamma$ is the type of LQG surface on which the process grows
- η determines the manner in which it grows

Let μ_{HARM} (resp. μ_{LEN}) be harmonic (resp. length) measure on a γ -LQG surface. The rate of growth (i.e., rate at which microscopic particles are added) is proportional to

$$\left(rac{d\mu_{ ext{HARM}}}{d\mu_{ ext{LEN}}}
ight)^\eta d\mu_{ ext{LEN}}.$$

QLE(8/3,0) is a member of a two-parameter family of processes called $QLE(\gamma^2,\eta)$

- $\blacktriangleright \ \gamma$ is the type of LQG surface on which the process grows
- η determines the manner in which it grows

Let μ_{HARM} (resp. μ_{LEN}) be harmonic (resp. length) measure on a γ -LQG surface. The rate of growth (i.e., rate at which microscopic particles are added) is proportional to

$$\left(rac{d\mu_{ ext{HARM}}}{d\mu_{ ext{LEN}}}
ight)^\eta d\mu_{ ext{LEN}}.$$

First passage percolation: $\eta = 0$

QLE(8/3,0) is a member of a two-parameter family of processes called $QLE(\gamma^2,\eta)$

- $\blacktriangleright~\gamma$ is the type of LQG surface on which the process grows
- η determines the manner in which it grows

Let $\mu_{\rm HARM}$ (resp. $\mu_{\rm LEN}$) be harmonic (resp. length) measure on a γ -LQG surface. The rate of growth (i.e., rate at which microscopic particles are added) is proportional to

$$\left(rac{d\mu_{ ext{HARM}}}{d\mu_{ ext{LEN}}}
ight)^\eta d\mu_{ ext{LEN}}.$$

- First passage percolation: $\eta = 0$
- Diffusion limited aggregation: $\eta = 1$

QLE(8/3,0) is a member of a two-parameter family of processes called $QLE(\gamma^2,\eta)$

- \blacktriangleright γ is the type of LQG surface on which the process grows
- η determines the manner in which it grows

Let μ_{HARM} (resp. μ_{LEN}) be harmonic (resp. length) measure on a γ -LQG surface. The rate of growth (i.e., rate at which microscopic particles are added) is proportional to

$$\left(rac{d\mu_{ ext{HARM}}}{d\mu_{ ext{LEN}}}
ight)^\eta d\mu_{ ext{LEN}}.$$

- First passage percolation: $\eta = 0$
- Diffusion limited aggregation: $\eta = 1$
- η -dieletric breakdown model: general values of η



Discrete approximation of $\mathrm{QLE}(8/3,0).$ Metric ball on a $\sqrt{8/3}\text{-}\mathsf{LQG}$



Discrete approximation of $\mathrm{QLE}(2,1).$ DLA on a $\sqrt{2}\text{-LQG}$



Each of the $QLE(\gamma^2, \eta)$ processes with (γ^2, η) on the orange curves is built from an SLE_{κ} process using tip re-randomization.



Euclidean Diffusion Limited Aggregation (DLA) introduced by Witten-Sander 1981.



DLA in nature: "A DLA cluster grown from a copper sulfate solution in an electrodeposition cell" (from Wikipedia)



DLA in nature: Magnese oxide patterns on the surface of a rock. (Halsey, Physics Today 2000)


DLA in nature: Magnese oxide patterns on the surface of a rock.



DLA in art: "High-voltage dielectric breakdown within a block of plexiglas" (from Wikipedia)

▶ Recall that in some sense c and γ are related by $Q = 2/\gamma + \gamma/2$ and $c = 25 - 6Q^2$.

- Recall that in some sense c and γ are related by Q = 2/γ + γ/2 and c = 25 - 6Q².
- ▶ Recent work with Ang, Park, Pfeffer gives way to "weight" by loop soup partition function in order to change γ of a "smooth planar map."

- Recall that in some sense c and γ are related by Q = 2/γ + γ/2 and c = 25 - 6Q².
- Recent work with Ang, Park, Pfeffer gives way to "weight" by loop soup partition function in order to change γ of a "smooth planar map."
- Check that higher genus smooth planar maps weighted by loop soups behave as they should. Really understand ghosts and all that.

- Recall that in some sense c and γ are related by Q = 2/γ + γ/2 and c = 25 - 6Q².
- Recent work with Ang, Park, Pfeffer gives way to "weight" by loop soup partition function in order to change γ of a "smooth planar map."
- Check that higher genus smooth planar maps weighted by loop soups behave as they should. Really understand ghosts and all that.
- Another approach. Start with Brownian map, say. Show weighting by number of loop soup loops longer than δ (in Liouville Brownian motion sense) changes γ in expected way. (Can formulate with eigenvalues and heat function also.)

Theorem for compact surfaces without boundary

THEOREM (Adapted by Ang, Park, Pfeffer, S.): Let (M, g) be a compact orientable surface. Then for δ > 0 small and C > 0 large we have, with γ the Euler-Mascheroni constant,

$$\mu_{\mathcal{M},g}^{\mathrm{loop}}(\mathcal{L}(\mathcal{M},g,\delta)ackslash\mathcal{L}(\mathcal{M},g,\mathcal{C})) = rac{\mathrm{Vol}_g(\mathcal{M})}{2\pi\delta} - rac{\chi(\mathcal{M})}{6}\log\delta + \log\mathcal{C} - \log\det'\Delta_g + O(\delta^{1/2}) + O(e^{-lpha\mathcal{C}}),$$

where $\alpha > 0$ depends on the manifold (M, g).

Theorem for compact surfaces without boundary

THEOREM (Adapted by Ang, Park, Pfeffer, S.): Let (M, g) be a compact orientable surface. Then for δ > 0 small and C > 0 large we have, with γ the Euler-Mascheroni constant,

$$\mu_{M,g}^{\mathrm{loop}}(\mathcal{L}(M,g,\delta) \setminus \mathcal{L}(M,g,C)) = \frac{\mathrm{Vol}_g(M)}{2\pi\delta} - \frac{\chi(M)}{6}\log\delta + \log C - \log \det'\Delta_g + O(\delta^{1/2}) + O(e^{-lpha C}),$$

where $\alpha > 0$ depends on the manifold (M, g).

COROLLARY: Let (S², g) be a sphere and η a simple smooth closed curve on the sphere. Then the mass of loops hitting γ of size between δ and C is given by

$$\frac{\operatorname{Len}_{g}(\eta)}{\sqrt{2\pi\delta}} + \log C - \log \operatorname{Vol}_{g}(S^{2}) - \frac{1}{12}I_{L}(\eta) -\mathcal{H}(S^{1},g) - \gamma - \log 2 + O(\delta^{1/2}) + O(e^{-\alpha C}),$$

where $I_L(\eta)$ is the Loewner energy of the curve η , and γ is the Euler-Mascheroni constant.

Theorem for compact surfaces without boundary

THEOREM (Adapted by Ang, Park, Pfeffer, S.): Let (M, g) be a compact orientable surface. Then for δ > 0 small and C > 0 large we have, with γ the Euler-Mascheroni constant,

$$\mu_{M,g}^{\mathrm{loop}}(\mathcal{L}(M,g,\delta) \setminus \mathcal{L}(M,g,C)) = rac{\mathrm{Vol}_g(M)}{2\pi\delta} - rac{\chi(M)}{6}\log\delta + \log C - \log \det'\Delta_g + O(\delta^{1/2}) + O(e^{-lpha C}),$$

where $\alpha > 0$ depends on the manifold (M, g).

COROLLARY: Let (S², g) be a sphere and η a simple smooth closed curve on the sphere. Then the mass of loops hitting γ of size between δ and C is given by

$$\begin{split} & \frac{\operatorname{Len}_g(\eta)}{\sqrt{2\pi\delta}} + \log \mathcal{C} - \log \operatorname{Vol}_g(\mathcal{S}^2) - \frac{1}{12} I_L(\eta) \\ & -\mathcal{H}(\mathcal{S}^1, g) - \gamma - \log 2 + O(\delta^{1/2}) + O(e^{-\alpha \mathcal{C}}), \end{split}$$

where $I_L(\eta)$ is the Loewner energy of the curve η , and γ is the Euler-Mascheroni constant.

Take loop mass on sphere, subtract loop mass in each half of $S^2 \setminus \eta$ applying Yilin-Wang (the quantity $\mathcal{H}(S^1, g)$ there is a nonexplicit constant).

Liouville quantum gravity: e^{γh(z)}dz where h is a GFF and γ ∈ [0, 2)





- Liouville quantum gravity: e^{γh(z)}dz where h is a GFF and γ ∈ [0, 2)
- Random surface model: Polyakov, 1980. Motivated by string theory.
- Rigorous construction of measure: Høegh-Krohn, 1971, $\gamma \in [0, \sqrt{2})$. Kahane, 1985, $\gamma \in [0, 2)$.

$$\gamma = 0.5$$



- Liouville quantum gravity: e^{γh(z)}dz where h is a GFF and γ ∈ [0, 2)
- Random surface model: Polyakov, 1980. Motivated by string theory.
- Rigorous construction of measure: Høegh-Krohn, 1971, $\gamma \in [0, \sqrt{2})$. Kahane, 1985, $\gamma \in [0, 2)$.
- Does not make literal sense since h takes values in the space of distributions.

$$\gamma = 0.5$$



- Liouville quantum gravity: e^{γh(z)}dz where h is a GFF and γ ∈ [0, 2)
- Random surface model: Polyakov, 1980. Motivated by string theory.
- Rigorous construction of measure: Høegh-Krohn, 1971, $\gamma \in [0, \sqrt{2})$. Kahane, 1985, $\gamma \in [0, 2)$.
- Does not make literal sense since h takes values in the space of distributions.
- Can make sense of random area measure using a regularization procedure.

$$\gamma = 0.5$$



- Liouville quantum gravity: e^{γh(z)}dz where h is a GFF and γ ∈ [0, 2)
- Random surface model: Polyakov, 1980. Motivated by string theory.
- Rigorous construction of measure: Høegh-Krohn, 1971, $\gamma \in [0, \sqrt{2})$. Kahane, 1985, $\gamma \in [0, 2)$.
- Does not make literal sense since h takes values in the space of distributions.
- Can make sense of random area measure using a regularization procedure.
- Areas of regions and lengths of curves are well defined.

$$\gamma = 1.0$$



- Liouville quantum gravity: e^{γh(z)}dz where h is a GFF and γ ∈ [0, 2)
- Random surface model: Polyakov, 1980. Motivated by string theory.
- Rigorous construction of measure: Høegh-Krohn, 1971, $\gamma \in [0, \sqrt{2})$. Kahane, 1985, $\gamma \in [0, 2)$.
- Does not make literal sense since h takes values in the space of distributions.
- Can make sense of random area measure using a regularization procedure.
- Areas of regions and lengths of curves are well defined.

$$\gamma = 1.5$$



- Liouville quantum gravity: e^{γh(z)}dz where h is a GFF and γ ∈ [0, 2)
- Random surface model: Polyakov, 1980. Motivated by string theory.
- Rigorous construction of measure: Høegh-Krohn, 1971, $\gamma \in [0, \sqrt{2})$. Kahane, 1985, $\gamma \in [0, 2)$.
- Does not make literal sense since h takes values in the space of distributions.
- Can make sense of random area measure using a regularization procedure.
- Areas of regions and lengths of curves are well defined.

$$\gamma = 2.0$$



Recall: add loops to SLE to get restriction measure.

- Recall: add loops to SLE to get restriction measure.
- Leads to some interested "mismatched" κ couplings. Mismatch makes them different from most obvious imaginary geometry flow line coupings.

- Recall: add loops to SLE to get restriction measure.
- Leads to some interested "mismatched" κ couplings. Mismatch makes them different from most obvious imaginary geometry flow line coupings.
- Is there a more general theory of such things?

- Recall: add loops to SLE to get restriction measure.
- Leads to some interested "mismatched" κ couplings. Mismatch makes them different from most obvious imaginary geometry flow line coupings.
- Is there a more general theory of such things?
- What if we move to random surface and throw in mismatched γ values?

- Recall: add loops to SLE to get restriction measure.
- Leads to some interested "mismatched" κ couplings. Mismatch makes them different from most obvious imaginary geometry flow line coupings.
- Is there a more general theory of such things?
- What if we move to random surface and throw in mismatched γ values?
- What Markov properties can we derive?

Big Problem 6: Ultra-High-Genus Maps

Very interesting new paper by Budzinski and Louf solves conjecture of Benjamini and Curien. Explains what happens when we take infinite volume limit of random surface in which the genus is proportional to the area, in terms of PSHT.

Big Problem 6: Ultra-High-Genus Maps

- Very interesting new paper by Budzinski and Louf solves conjecture of Benjamini and Curien. Explains what happens when we take infinite volume limit of random surface in which the genus is proportional to the area, in terms of PSHT.
- At first glance somewhat surprising that small handles not seen in infinite volume limit. But one can think about (have at least n^{2g} choices) and it starts to make sense.

Big Problem 6: Ultra-High-Genus Maps

- Very interesting new paper by Budzinski and Louf solves conjecture of Benjamini and Curien. Explains what happens when we take infinite volume limit of random surface in which the genus is proportional to the area, in terms of PSHT.
- At first glance somewhat surprising that small handles not seen in infinite volume limit. But one can think about (have at least n^{2g} choices) and it starts to make sense.
- Problem is to give theory that incudes weighting by statistical physics models. New notion of entropy: exponential part that persists after removing factorial part?

▶ Understand discrete analog of LQG for $c \in (1, 25)$. Applications?

- ▶ Understand discrete analog of LQG for $c \in (1, 25)$. Applications?
- Nina Holden will talk about recent work on this (Gwynne, Holden, Pfeffer, Remy). Other work in progress by Ang, Park, Pfeffer, S.

- Understand discrete analog of LQG for $c \in (1, 25)$. Applications?
- Nina Holden will talk about recent work on this (Gwynne, Holden, Pfeffer, Remy). Other work in progress by Ang, Park, Pfeffer, S.
- Embedding surface in high dimension (and fixing area) seems to give tree. But there are interesting infinite volume surfaces.

▶ Understand discrete analog of LQG for $c \in (1, 25)$. Applications?

- Nina Holden will talk about recent work on this (Gwynne, Holden, Pfeffer, Remy). Other work in progress by Ang, Park, Pfeffer, S.
- Embedding surface in high dimension (and fixing area) seems to give tree. But there are interesting infinite volume surfaces.
- What interesting random surfaces have something like this as infinite volume limit?

▶ Understand discrete analog of LQG for $c \in (1, 25)$. Applications?

- Nina Holden will talk about recent work on this (Gwynne, Holden, Pfeffer, Remy). Other work in progress by Ang, Park, Pfeffer, S.
- Embedding surface in high dimension (and fixing area) seems to give tree. But there are interesting infinite volume surfaces.
- What interesting random surfaces have something like this as infinite volume limit?
- High genus but forced to have lots of small disjoint homotopically different loops?

Big Problem 8: scaling limits, scaling limits, scaling limits

PLANAR MAP MODEL

- a. Laplacian-determinant power weighted (loop soup, tree, GFF)
- b. Ising/FK/Potts weighted
- c. Other decoration (bipolar orientation, Schnyder wood, various things)
- d. Continuum-derived: mated-CRT map, Poisson-Voronoi, square subdivision

Big Problem 8: scaling limits, scaling limits, scaling limits

PLANAR MAP MODEL

- a. Laplacian-determinant power weighted (loop soup, tree, GFF)
- b. Ising/FK/Potts weighted
- c. Other decoration (bipolar orientation, Schnyder wood, various things)
- d. Continuum-derived: mated-CRT map, Poisson-Voronoi, square subdivision

TOPOLOGY OF CONVERGENCE

- A. Path/measure converence under Tutte/Cardy/circle-packing/square tiling
- B. Gromov-Hausdorff
- C. Peanosphere (law of pair of trees)
- D. Topology encoding loop lengths and adjacency relationships, etc.

Lots of truly spectacular work (Le Gall, Miermont, Curien, Holden-Sun, Gwynne, etc.) But most questions are open.

Big Problem 9: gauge theory, gauge theory, gauge theory

Don't have to start with d = 4 Yang-Mills with compact gauge group. Just give any satisfying LQG/gauge-theory link.

Big Problem 9: gauge theory, gauge theory, gauge theory

- Don't have to start with d = 4 Yang-Mills with compact gauge group. Just give any satisfying LQG/gauge-theory link.
- ► Think about moving simple case (c =≤ 1, genus finite, Gaussian matrix measure) across one or more of the three barriers.

Considered by many the most challenging open problem in the field.

- Considered by many the most challenging open problem in the field.
- So challenging that the people who solve it get...

- Considered by many the most challenging open problem in the field.
- So challenging that the people who solve it get...
- ► No recognition at all.

- Considered by many the most challenging open problem in the field.
- So challenging that the people who solve it get...
- ► No recognition at all.
- Proceedings of London Math Society (PLMS)... like JEMS, JAMS